

QUANTUM LAX PAIRS VIA DUNKL AND CHEREDNIK OPERATORS

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ABSTRACT. We establish a direct link between Dunkl operators and quantum Lax matrices \mathcal{L} for the Calogero–Moser systems associated to an arbitrary Weyl group W (or an arbitrary finite reflection group in the rational case). This interpretation also provides a companion matrix \mathcal{A} so that \mathcal{L}, \mathcal{A} form a quantum Lax pair. Moreover, such an \mathcal{A} can be associated to any of the higher commuting quantum Hamiltonians of the system, so we obtain a family of quantum Lax pairs. These Lax pairs can be of various sizes, matching the sizes of orbits in the reflection representation of W , and in the elliptic case they contain a spectral parameter. This way we reproduce universal classical Lax pairs by D’Hoker–Phong and Bordner–Corrigan–Sasaki, and complement them with quantum Lax pairs in all cases (including the elliptic case, where they were not previously known). The same method, with the Dunkl operators replaced by the Cherednik operators, produces quantum Lax pairs for the generalised Ruijsenaars systems for arbitrary root systems. As one of the main applications, we calculate a Lax matrix for the elliptic BC_n case with nine coupling constants (van Diejen system), thus providing an answer to a long-standing open problem.

1. INTRODUCTION

The notion of a Lax pair has for a long time been instrumental in both finite- and infinite-dimensional integrable systems, with the earliest examples given by P. Lax and H. Flaschka [L, Fl]. Another famous example is the Lax pair found by J. Moser [Mo] for the classical rational Calogero–Moser system [Ca1], which is a system of n interacting particles on the line with coordinates x_1, \dots, x_n , described by the Hamiltonian

$$H = \frac{1}{2} \sum_{k=1}^n p_k^2 + g^2 \sum_{k<l}^n (x_k - x_l)^{-2}. \quad (1.1)$$

Here g is a coupling constant. The Lax presentation for this system involves two matrices L and A of size n of the following form:

$$L_{kl} = \begin{cases} ig(x_k - x_l)^{-1} & \text{for } k \neq l \\ p_k & \text{for } k = l, \end{cases} \quad A_{kl} = \begin{cases} ig(x_k - x_l)^{-2} & \text{for } k \neq l \\ -ig \sum_{j \neq k}^n (x_j - x_k)^{-2} & \text{for } k = l. \end{cases} \quad (1.2)$$

This allows presenting the equations of motion in the form

$$\frac{dL}{dt} = [A, L]. \quad (1.3)$$

One immediate corollary is that $H_k = \text{tr } L^k$, $k \in \mathbb{N}$ are conserved quantities (integrals of motion).

The quantum Calogero–Moser system is described by a Schrödinger operator

$$\widehat{H} = \frac{1}{2} \sum_{k=1}^n \widehat{p}_k^2 + g(g - \hbar) \sum_{k<l}^n (x_k - x_l)^{-2}, \quad \widehat{p}_k = -i\hbar \frac{\partial}{\partial x_k}. \quad (1.4)$$

Quantum analogues \mathcal{L}, \mathcal{A} of the above L, A were introduced in [UHW] (see also [BGHP, SS]). They are obtained by replacing p_k by \widehat{p}_k and multiplying A by $i\hbar$:

$$\mathcal{L}_{kl} = \begin{cases} ig(x_k - x_l)^{-1} & \text{for } k \neq l \\ \widehat{p}_k & \text{for } k = l, \end{cases} \quad \mathcal{A} = i\hbar A. \quad (1.5)$$

These are matrices of size n whose entries are partial differential operators. To write down the quantum Lax equation, introduce a diagonal matrix $\mathcal{H} = \widehat{H} \mathbb{1}_n$. The following can then be confirmed by a direct calculation:

$$[\mathcal{L}, \mathcal{H}] = [\mathcal{A}, \mathcal{L}]. \quad (1.6)$$

To see why this is indeed analogous to (1.3), note that (1.6) can be written as

$$(i\hbar)^{-1}(\mathcal{L}_{kl}\widehat{H} - \widehat{H}\mathcal{L}_{kl}) = [A, \mathcal{L}]_{kl}, \quad k, l = 1, \dots, n. \quad (1.7)$$

In the classical limit $\hbar \rightarrow 0$, \mathcal{L} and \widehat{H} reduce to L and H , respectively, and the left-hand side reduces to the Poisson bracket $\{L_{kl}, H\}$. Thus,

$$\frac{dL_{kl}}{dt} = \{L_{kl}, H\} = [A, L]_{kl}, \quad (1.8)$$

which is (1.3).

Similarly to the classical case, the above quantum Lax matrix can be used to produce first integrals for the Hamiltonian (1.4). Namely, following [UHW, SS, BGHP], consider a pair \mathbf{v}, \mathbf{w} of n -component column and row vectors

$$\mathbf{v} = (1, \dots, 1)^T, \quad \mathbf{w} = (1, \dots, 1). \quad (1.9)$$

Then we have the following easily verified properties of the matrix \mathcal{A} in (1.2), (1.5):

$$\mathbf{w}\mathcal{A} = 0, \quad \mathcal{A}\mathbf{v} = 0. \quad (1.10)$$

Now define

$$\widehat{H}_k = \mathbf{w}\mathcal{L}^k\mathbf{v}, \quad k \in \mathbb{N}. \quad (1.11)$$

(Note that \widehat{H}_k is the sum of all entries of \mathcal{L}^k .) Then

$$[\widehat{H}, \widehat{H}_k] = \mathbf{w}[\widehat{H}\mathbb{1}_n, \mathcal{L}^k]\mathbf{v} = \mathbf{w}[\mathcal{H}, \mathcal{L}^k]\mathbf{v} = \mathbf{w}[\mathcal{L}^k, \mathcal{A}]\mathbf{v} = \mathbf{w}\mathcal{L}\mathcal{A}\mathbf{v} - \mathbf{w}\mathcal{A}\mathcal{L}\mathbf{v} = 0,$$

where we used (1.10) and the relation $[\mathcal{L}^k, \mathcal{H} + \mathcal{A}] = 0$. Therefore, \widehat{H}_k are quantum integrals.

The original papers [Mo, UHW] simply present the above Lax pairs but do not explain how they were found (cf. [Ca2] where this was related to solving certain functional equations, leading to a Lax pair in the elliptic case). More conceptual ways of actually *deriving* Moser's Lax matrix have been subsequently discovered in [KKS] in the framework of symplectic reduction, and in [Kr1] in connection with the KP hierarchy. There is by now a vast literature devoted to various further generalisations and development of those ideas. However, the quantum Lax pairs lacked such an interpretation, although many authors remark on a similarity between Moser's Lax matrix and the Dunkl operators [D] which played a pivotal role in the theory of Calogero–Moser systems since the works [He1, He2, C1, Op]. The present paper fills that gap: as we explain, there is a direct link between the Dunkl operators and quantum Lax pairs. Our approach is inspired by an observation due to Etingof and Ginzburg, who in [EG] derived the classical Moser's Lax matrix from the representation theory of Cherednik algebras. The main difference is that we work at the quantum level (and in a more general situation), so the classical Lax pairs are obtained by letting $\hbar \rightarrow 0$. In the elliptic case we use elliptic Dunkl operators [BFV, EM] together with some important ideas from [EFMV]. In this case the Lax matrices which we construct contain a spectral parameter (intrinsically linked to the theory of Dunkl operators). This way we reproduce, in a much simpler and more conceptual way, the previously known Lax pairs from the papers [DHP, BCS, BMS, KPS], as well as find some new ones (for instance, quantum Lax pairs were not known in the elliptic case). Also, our method allows us to associate a Lax partner to each of the commuting Hamiltonians of the Calogero–Moser problem, so we get a family of compatible Lax pairs. As a corollary, this gives a simple uniform proof of the fact that the classical Lax matrix L remains isospectral under all of the commuting flows, implying that the functions $\text{tr } L^k$ are in involution. Such a property is well-known in type A [Mo, Pe] and it is, of course, to be expected in other cases, but no general proof of that fact was available. Note that it is customary to study an integrable system first at the classical level and to use its classical Lax matrix to get an insight into a possibility of a quantisation. By contrast, we derive non-trivial properties of the classical Lax matrix by studying its quantum counterpart. It seems rather surprising that the Lax pairs are actually easier to understand at the quantum level.

Perhaps more importantly, our construction works for the systems of Ruijsenaars–Schneider type (also referred to as *relativistic Calogero–Moser systems*). The usual Ruijsenaars–Schneider system [RS] corresponds to $R = A_n$; its quantum version was introduced by Ruijsenaars [R1], who also proved its complete integrability. A classical Lax pair for this system is well-known [R1, BCa, KrZ], see also [Ha] where a quantum Lax matrix was introduced. However, for the

models related to other root systems the question remained open for a long time (for instance, it was raised already by Inozemtsev in [I]). The best result in that direction has been obtained recently by Görbe and Pusztai, who constructed in [GoP] a Lax pair for a two-parameter subfamily of the Koornwinder–van Diejen system (see also [Pu] where a Lax matrix was found for the three-parameter rational case). The Koornwinder–van Diejen system [Ko, vD2] is a BC_n version of the trigonometric Ruijsenaars–Schneider system, depending on five coupling parameters. There is also an elliptic version with nine parameters, introduced by van Diejen in [vD1], whose integrability was shown in [KH1]. However, even in the trigonometric case with five parameters a Lax matrix remained unknown, let alone the nine-parameter elliptic version. Thus, it is rather pleasing that within our approach we are able to calculate it explicitly without much effort.

We also give a general construction of Lax pairs for the generalised Ruijsenaars systems related to arbitrary root systems. Note that in the relativistic case instead of the Dunkl operators one needs to use their q -analogues, known as Cherednik operators [C2, C3]. The theory of Cherednik operators is well established in the trigonometric case, where they are intimately related to the theory of double affine Hecke algebras and Macdonald polynomials [C4, M2]. Their elliptic analogues were introduced by Komori and Hikami [KH2] following the ideas of Cherednik [C6], but some features of the trigonometric case seemed missing (or looked puzzlingly different) in the elliptic case. An issue here is that in the elliptic case one has braid (or Yang–Baxter) relations but no quadratic Hecke relations. As a result, elliptic Cherednik operators are only defined up to scaling (which in addition may depend on the dynamical variables), and so a correct way of defining them is not immediately obvious. We observe that a particularly well-behaved choice is the one associated with *unitary* R -matrices. It is this choice which allows us to draw a parallel with the results of [EFMV] and construct Lax pairs for the generalised elliptic Ruijsenaars systems for all root systems. Calculating these Lax pairs explicitly is not easy in general (or even impractical: for instance, the smallest Lax matrix in the E_8 case has size 240). We carry out such a calculation in two important cases: for the standard elliptic Ruijsenaars system and for the elliptic van Diejen system, i.e. for the A_n and BC_n cases of the theory. A crucial realization that such a calculation was possible came to us after seeing a paper of Nazarov and Sklyanin [NS] in which they calculated, rather nicely, a quantum Lax matrix for the trigonometric Ruijsenaars system. Similarly to us, they derive their Lax matrix directly from Cherednik operators, and although they do not consider quantum Lax pairs, some of their considerations are very close to ours. However, our approach is more general, in particular it extends to the BC_n case, including the elliptic version. A nice special feature of the relativistic A_n and BC_n cases is that we can construct a Lax pair for each of the commuting Hamiltonians. In the BC_n case this relies on results of Rains, who recently developed a geometric approach to elliptic DAHAs [Ra]. For other root systems we are able to construct Lax pairs only for the Hamiltonians corresponding to minuscule and quasi-minuscule coweights. These Hamiltonians are the Macdonald operators [M1] and their elliptic analogues [KH2]. Since every root system has a (quasi)-minuscule coweight, we obtain at least one Lax pair for each root system.

Let us remark that there exist various geometric approaches to Calogero–Moser and Ruijsenaars–Schneider systems, see [GNe, Ne, KrZ, FR, HM, Kr2, Kr3, KrS, FeK, BZN, LOSZ, FeM, KPSZ] (where also many further references can be found). It would be interesting to see whether our quantum Lax pairs admit a geometric interpretation within any of those approaches. We also would like to mention that our interest in this problem was triggered by a paper by Sergeev and Veselov [SV3] in which they construct quantum Lax pairs for certain deformed Calogero–Moser systems. We expect that the methods of the present paper can be adapted to give a conceptual approach to quantum Lax pairs for other deformed systems [SV1, SV2, Fe, FeS]. Our results in the BC_n case are also of crucial importance for constructing action-angle coordinates for the classical Koornwinder–van Diejen system (in the A_n case this was done in [R2]). Also, they give valuable hints towards a possible description of the center of the double affine Hecke algebra of type CC_n^V (cf. [Ob] in the A_n case). This is also an important ingredient for describing the moduli space of ideals of the Askey–Wilson algebra. Some of these problems will be a subject of future work.

The structure of the paper is as follows. In Section §2 we present the main construction in the rational case and illustrate it by deriving the quantum Lax pair (1.5). Section §3 generalises this to the trigonometric relativistic case, where the Dunkl operators are replaced by the Cherednik operators. In Section §4 we apply the results of Section §3 for calculating a Lax matrix for the

Koornwinder–van Diejen system. Section §5 deals with the elliptic Calogero–Moser systems. Here we use elliptic Dunkl operators [BFV], and while the main idea remains the same, the construction of a Lax pair is more involved and uses ideas from [EFMV]. Section §6 is devoted to the elliptic difference case, related to elliptic Cherednik and Macdonald–Ruijsenaars operators from [KH2]. Our main effort here is to establish the existence of a Lax pair for any root system (Proposition 6.13 and Theorem 6.15). Subsections §6.6–§6.10 are devoted to the BC_n case; a Lax matrix for the elliptic van Diejen system is calculated in Subsection §6.10.

The structure of the paper reflects how it developed over time: the main constructions and results in Sections §2, §3 go back to 2015, while the calculations in §3.9 and §4.2 were inspired by the work [NS]. The result of Proposition 3.3 is also a later addition, prompted by [NS, Corollary 2.6]. Sections §5, §6 are more recent. Note that some of the results in Sections §2, §3 can be obtained from the elliptic case as a limit. However, we decided to have them derived independently, mainly because some interesting features (for instance, Proposition 3.3) do not seem to have an analogue in the elliptic case.

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2. DUNKL OPERATORS AND QUANTUM LAX PAIRS

2.1. Let us first recall the well-known link between Dunkl operators and rational Calogero–Moser systems [D, He1], cf. [Po, BHV]. For simplicity, we restrict ourselves to the case of real Coxeter groups, but everything applies with minimal changes to any complex reflection group by replacing the Dunkl operators by their complex analogues [DO].

Let W be a finite Coxeter group and $V_{\mathbb{R}}$ be its reflection representation. We will work over \mathbb{C} , so $V := V_{\mathbb{R}} \otimes_{\mathbb{R}} \mathbb{C}$ will be an n -dimensional complex vector space with a fixed W -invariant scalar product $\langle -, - \rangle$. Let $R = R_+ \sqcup -R_+$ be the root system of W (not necessarily crystallographic). For each $\alpha \in R$ we have the orthogonal reflection $s_{\alpha} \in W$ acting on V by the formula $s_{\alpha}(x) = x - 2 \frac{\langle \alpha, x \rangle}{\langle \alpha, \alpha \rangle} \alpha$, and these reflections generate the group W . We assume that the set R is W -invariant. Below we will always identify V with its dual by using the scalar product $\langle -, - \rangle$, and hence equip $V \times V$ with a symplectic form transferred from $V \times V^* = T^*V$.

Denote by $\mathbb{C}(V)$ and $\mathcal{D}(V)$ the rings of meromorphic functions and differential operators on V with meromorphic coefficients, respectively. The group W acts naturally on $\mathbb{C}(V)$ and $\mathcal{D}(V)$, so we form the crossed products $\mathbb{C}(V) * W$ and $\mathcal{D}(V) * W$. As an algebra, $\mathcal{D}(V) * W$ is generated by the elements $w \in W$, $f \in \mathbb{C}(V)$, and derivations ∂_{ξ} , $\xi \in V$, subject to the relations

$$w \partial_{\xi} = \partial_{w\xi} w, \quad w f = f^w w, \quad \text{where } f^w(x) = f(w^{-1}x).$$

Any $a \in \mathcal{D}(V) * W$, admits a unique presentation

$$a = \sum_{w \in W} a_w w \quad \text{with } a_w \in \mathcal{D}(V). \quad (2.1)$$

Let us fix parameters $t \neq 0$ and a W -invariant function $c : R \rightarrow \mathbb{C}$, with $c(\alpha)$ abbreviated to c_{α} . Introduce *Dunkl operators* as the following elements of $\mathcal{D}(V) * W$:

$$y_{\xi} := t \partial_{\xi} + \sum_{\alpha \in R_+} \frac{\langle \alpha, \xi \rangle}{\langle \alpha, x \rangle} c_{\alpha} s_{\alpha}, \quad \xi \in V. \quad (2.2)$$

The two main properties of the Dunkl operators are their commutativity and equivariance: for all $\xi, \eta \in V$ and $w \in W$,

$$y_{\xi} y_{\eta} = y_{\eta} y_{\xi}, \quad w y_{\xi} w^{-1} = y_{w(\xi)}. \quad (2.3)$$

Therefore, the assignment $\xi \mapsto y_{\xi}$ extends to a W -equivariant injective algebra map

$$SV = \bigoplus_{i \geq 0} S^i V \rightarrow \mathcal{D}(V) * W. \quad (2.4)$$

The image of $q \in SV$ under this map is denoted by $q(y)$. Let $\partial_i = \partial_{\xi_i}$ and $y_i = y_{\xi_i}$, where $\{\xi_i \mid i = 1 \dots n\}$ is an orthonormal basis in V . Writing $\langle y, y \rangle := y_1^2 + \dots + y_n^2$, we have by [D]:

$$\langle y, y \rangle = t^2 \Delta_V - \sum_{\alpha \in R_+} \langle \alpha, \alpha \rangle \langle \alpha, x \rangle^{-2} c_\alpha (c_\alpha + t s_\alpha), \quad \Delta_V = \sum_{i=1}^n \partial_i^2. \quad (2.5)$$

Let

$$e = \frac{1}{|W|} \sum_{w \in W} w \quad (2.6)$$

be the symmetrizing idempotent in the group algebra $\mathbb{C}W$. For any W -invariant element $q \in (SV)^W$, we have

$$q(y)e = L_q e, \quad L_q \in \mathcal{D}(V)^W \quad (2.7)$$

for some uniquely defined W -invariant differential operator L_q . Explicitly, if $q(y)$ is presented in the form (2.1), $q(y) = \sum_{w \in W} a_w w$, then $L_q = \sum_{w \in W} a_w$. In particular, for $\xi^2 := \xi_1^2 + \dots + \xi_n^2$ we find from (2.5) that

$$L_{\xi^2} = t^2 \Delta_V - \sum_{\alpha \in R_+} \frac{c_\alpha (c_\alpha + t) \langle \alpha, \alpha \rangle}{\langle \alpha, x \rangle^2}. \quad (2.8)$$

Substituting $t = -i\hbar$ and $c_\alpha = ig_\alpha$ gives the quantum Calogero–Moser Hamiltonian associated to the group W [OP1, OP2]:

$$\widehat{H} = \sum_{k=1}^n \widehat{p}_k^2 + \sum_{\alpha \in R_+} \frac{g_\alpha (g_\alpha - \hbar) \langle \alpha, \alpha \rangle}{\langle \alpha, x \rangle^2}, \quad \widehat{p}_k = -i\hbar \partial_{\xi_k}. \quad (2.9)$$

From the commutativity of the Dunkl operators it follows that the operators L_q , $q \in (SV)^W$ pairwise commute [He1]. Since $(SV)^W$ is a free polynomial algebra on $n = \dim V$ generators, this proves that the quantum Hamiltonian (2.9) is completely integrable.

2.2. Let us now explain how to construct a quantum Lax pair for the Hamiltonian (2.9). For that we will work in a special representation of $\mathcal{D}(V) * W$. Namely, let us view $\mathbb{C}(V)$ as a left $\mathcal{D}(V)$ -module with the usual action by differential operators, and consider the induced module

$$M = \text{Ind}_{\mathcal{D}(V)}^{\mathcal{D}(V) * W} \mathbb{C}(V).$$

We can write elements of M as $f = \sum_{w \in W} w f_w$ with $f_w \in \mathbb{C}(V)$, thus identifying M and $\mathbb{C}W \otimes \mathbb{C}(V)$ (as a vector spaces). The algebra $\text{End}_{\mathbb{C}}(M)$ then is identified with $\text{End}_{\mathbb{C}}(\mathbb{C}W) \otimes \text{End}_{\mathbb{C}}(\mathbb{C}(V))$, i.e. with *operator-valued* matrices of size $|W|$. As a result, the (left) action of $\mathcal{D}(V) * W$ on M gives a faithful representation

$$\mathcal{D}(V) * W \rightarrow \text{Mat}(|W|, \mathcal{D}(V)). \quad (2.10)$$

For a W -invariant $a \in \mathcal{D}(V)$ we have: $a(\sum_{w \in W} w f_w) = \sum_{w \in W} w (a f_w)$. Therefore, in the above representation such a acts as $a\mathbb{1}$.

Now pick a Dunkl operator y_ξ ; obviously, it commutes with $\langle y, y \rangle$. From (2.5) we have:

$$\frac{1}{2} \langle y, y \rangle = \widehat{H} + \widehat{A}, \quad \text{where} \quad \widehat{H} := \frac{1}{2} L_{\xi^2}, \quad \widehat{A} := \frac{t}{2} \sum_{\alpha \in R_+} c_\alpha \langle \alpha, \alpha \rangle \langle \alpha, x \rangle^{-2} (1 - s_\alpha). \quad (2.11)$$

As a result, if we set \mathcal{L} , \mathcal{H} , \mathcal{A} to be the matrices representing under (2.10) the action of y_ξ , \widehat{H} and \widehat{A} , respectively, we obtain

$$[\mathcal{L}, \mathcal{H} + \mathcal{A}] = 0, \quad (2.12)$$

which is (1.6). Since \widehat{H} is W -invariant, the matrix \mathcal{H} is $\widehat{H}\mathbb{1}$. Therefore, we have obtained a quantum Lax pair \mathcal{L}, \mathcal{A} of matrices of size $|W|$.

In fact, using this approach one can associate a suitable \mathcal{A} to any of the commuting quantum Hamiltonians L_q , $q \in (SV)^W$. Indeed, suppose $q(y) = \sum_{w \in W} a_w w$ with $a_w \in \mathcal{D}(V)$. Then we have

$$q(y) = L_q + \widehat{A}, \quad \text{where} \quad L_q = \sum_{w \in W} a_w, \quad \widehat{A} := \sum_{w \in W} a_w (w - 1), \quad (2.13)$$

so the above construction gives a Lax pair with the same \mathcal{L} but with different \mathcal{H}, \mathcal{A} .

Remark 2.1. In the above construction one can replace the ring $\mathcal{D}(V)$ with a smaller ring $\mathcal{D}(V_{\text{reg}})$ of algebraic differential operators on V_{reg} , the complement to the reflection hyperplanes. Furthermore, when constructing the module M , one can induce from any $\mathcal{D}(V_{\text{reg}})$ -module, e.g. space of analytic functions on a small neighbourhood of a point in V_{reg} . Therefore, one can allow elements of M to be multivalued, with branching along the reflection hyperplanes in V .

2.3. The classical limit corresponds to taking $t \rightarrow 0$. More precisely, we set $t = -i\hbar$ and view the Dunkl operators as elements of the algebra

$$A_{\hbar} * W = \mathbb{C}(V)[\hat{p}_1, \dots, \hat{p}_n][[\hbar]] * W,$$

where the quantum momenta $\hat{p}_k = -i\hbar\partial_{\xi_k}$ satisfy the standard relations $[\hat{p}_k, f] = -i\hbar\partial_{\xi_k}f$ for $f \in \mathbb{C}(V)$. We have an algebra isomorphism

$$\eta_0 : A_{\hbar} * W / \hbar A_{\hbar} * W \rightarrow A_0 * W, \quad f \mapsto f, \quad \hat{p}_k \mapsto p_k, \quad w \mapsto w,$$

where $A_0 = \mathbb{C}(V)[p_1, \dots, p_n]$ is the classical version of A_{\hbar} . Therefore, A_{\hbar} (resp. $A_{\hbar} * W$) is a formal deformation of A_0 (resp. $A_0 * W$), see, e.g., [E, 3.1]. Note that A_0 is commutative, with the standard Poisson bracket satisfying $[a, b] = i\hbar\{\eta_0(a), \eta_0(b)\} + o(\hbar)$ for $a, b \in A_{\hbar}$. For any $a \in A_{\hbar} * W$, we call $\eta_0(a)$ the *classical limit* of a . For example, the classical limit of (2.2) is

$$y_{\xi}^c = p_{\xi} + \sum_{\alpha \in R_+} \frac{\langle \alpha, \xi \rangle}{\langle \alpha, x \rangle} c_{\alpha} s_{\alpha},$$

which is called a classical Dunkl operator, see [E, 6.30]. Here p_{ξ} is the classical momentum in direction ξ . The operators y_{ξ}^c are commuting elements of $A_0 * W$, so we have a classical variant of the map (2.4):

$$SV = \bigoplus_{i \geq 0} S^i V \rightarrow A_0 * W, \quad q \mapsto q(y^c).$$

The classical limits of (2.7) and (2.13) can be obtained by replacing Dunkl operators y_{ξ} by their classical counterparts y_{ξ}^c . We will use the following important fact.

Lemma 2.2 (Lemma 2.2, [EFMV]). *For any $q \in (SV)^W$, when writing $q(y^c) = \sum_{w \in W} a_w w$ with $a_w \in A_0$, we have $a_w = 0$ for all $w \neq \text{id}$.*

This tells us that the classical limit of \hat{A} in (2.13) is zero, in other words, the classical limit of $\hbar^{-1}\hat{A}$ is well-defined. As a result, we have well-defined classical limits L, H, A of \mathcal{L}, \hat{H} and $(i\hbar)^{-1}\mathcal{A}$, respectively, after which the classical Lax equation follows in the same way as in (1.7), (1.8).

2.4. In [BCS, BMS], classical and quantum Lax pairs of various sizes were constructed, so let us explain how they arise within our approach. We start again by picking y_{ξ} and $q(y)$ with $\xi \in V$, $q \in (SV)^W$ and writing $q(y) = L_q + \hat{A}$, as in (2.13). To get a Lax pair of a smaller size, we choose ξ with non-trivial stabiliser, writing

$$W' = \{w \in W \mid w\xi = \xi\}, \quad e' = \frac{1}{|W'|} \sum_{w \in W'} w.$$

Obviously, $y_{\xi}e' = e'y_{\xi}$; also, $q(y)e' = e'q(y)$, $L_q e' = e' L_q$ by their W -invariance. As a result, the operators y_{ξ} , $q(y)$, L_q and \hat{A} preserve the subspace

$$M' = e' M \cong e' CW \otimes \mathbb{C}(V).$$

The left W -module $e'CW$ has dimension equal to $|W/W'|$, i.e. to the size of the orbit $W(\xi)$. Therefore, restricting $\mathcal{L}, \mathcal{H}, \mathcal{A}$ onto M' produces a quantum Lax pair of size $|W/W'|$, with the smaller sizes achieved when ξ is a fundamental weight. This agrees with the Lax pairs in [BCS, BMS, KPS].

2.5. We can also explain why (1.10) holds in general. Namely, pick representatives w_i for the cosets in $W' \backslash W$. Elements of M' are linear combinations of $e'w_i f_i$ with $f_i \in \mathbb{C}(V)$. Now, from (2.7) and (2.13) we have $\widehat{A}e = 0$. Similarly, $eq(y) = eL_q$, so $e\widehat{A} = 0$. Multiplication by e acts on M' , and it is easy to see that the associated matrix is $\frac{1}{|W/W'|} \mathbf{v}\mathbf{w}$, where \mathbf{v}, \mathbf{w} are the column/row vectors as in (1.9) (of size $|W/W'|$). The relations $\widehat{A}e = e\widehat{A} = 0$ easily imply (1.10). Therefore, the formula (1.11) always produces quantum integrals; for instance, it works for any complex reflection group W .

An alternative explanation of (1.11) is as follows. Consider the operators $ey_\xi^k e$ with $k \in \mathbb{N}$. The symmetrizer e acts on M' by $\frac{1}{|W/W'|} \mathbf{v}\mathbf{w}$, and y_ξ acts by \mathcal{L} . Therefore, up to a constant factor, $ey_\xi^k e$ acts as $\mathbf{v}\mathbf{w}\mathcal{L}^k \mathbf{v}\mathbf{w} = (\mathbf{w}\mathcal{L}^k \mathbf{v})\mathbf{v}\mathbf{w} = |W/W'|(\mathbf{w}\mathcal{L}^k \mathbf{v})e$. On the other hand, $ey_\xi^k e = \frac{1}{|W|} \left(\sum_{w \in W} y_{w\xi}^k \right) e$, so it commutes with any of the operators $eq(y)e = L_q e$, $q \in (SV)^W$. As a result, $\mathbf{w}\mathcal{L}^k \mathbf{v}$ and L_q commute.

In Proposition 3.3 below we generalise this result to the relativistic case.

2.6. As an illustration, let us derive the quantum Lax pair (1.5). We consider $W = \mathfrak{S}_n$ acting on $V = \mathbb{C}^n$ by permuting the basis vectors; it is generated by permutations \mathfrak{s}_{ij} , $i \neq j$. The ring $\mathbb{C}(V) = \mathbb{C}(x)$ is the ring of functions of n variables x_1, \dots, x_n . We have n commuting Dunkl operators,

$$y_i = t\partial_i + \sum_{j \neq i} c(x_i - x_j)^{-1} \mathfrak{s}_{ij}, \quad i = 1 \dots n.$$

Choose y_1 , so $\xi = (1, 0, \dots, 0)$ and $W' = \mathfrak{S}_{n-1}$ is the permutation group on $\{2, \dots, n\}$, with

$$e' = \frac{1}{(n-1)!} \sum_{w \in \mathfrak{S}_{n-1}} w.$$

We pick \mathfrak{s}_{1j} , $i = 1 \dots n$ as representatives for the cosets in $W' \backslash W$ (with $\mathfrak{s}_{11} := \text{id}$), and write elements of $M' = e'M$ as

$$f = \sum_{i=1}^n e' \mathfrak{s}_{1i} f_i, \quad \text{with } f_i \in \mathbb{C}(x).$$

To find the matrix representing the action of y_1 on M' , it will be useful to work in a greater generality.

Lemma 2.3. *Suppose that we have an element $Z = \sum_{i=1}^n Z_i \mathfrak{s}_{1i}$ with $Z_i \in \mathcal{D}(V)$, whose action preserves $M' = e'M$.*

(1) Z_1 is symmetric in x_2, \dots, x_n and Z_i for $i > 1$ is symmetric in x_j with $j \neq 1, i$. Also, we have $Z_i^{\mathfrak{s}_{ij}} = Z_j$ for any $i, j > 1$.

(2) For each $1 \leq i, j \leq n$, we have

$$e' \mathfrak{s}_{1i} \mathfrak{s}_{1j} = e' \mathfrak{s}_{1k}, \quad \text{with } k = k(i, j) = \begin{cases} i & \text{for } i \neq 1, j \\ j & \text{for } i = 1 \\ 1 & \text{for } i = j. \end{cases}$$

(3) For any i, j and $k = k(i, j)$ as above, the (k, j) -th entry the matrix representing the action of Z on M' is calculated as $(Z_i)^{\mathfrak{s}_{1j} \mathfrak{s}_{1i}}$.

Proof. Parts (1) follows from $\mathfrak{s}_{ij} Z e' = Z e'$ for $i, j > 1$. Part (2) is straightforward. For part (3),

$$Z \sum_{j=1}^n e' \mathfrak{s}_{1j} f_j = \sum_{j=1}^n e' Z \mathfrak{s}_{1j} f_j = \sum_{i,j=1}^n e' \mathfrak{s}_{1i} \mathfrak{s}_{1j} Z_i^{\mathfrak{s}_{1j} \mathfrak{s}_{1i}} f_j = \sum_{i,j=1}^n e' \mathfrak{s}_{1k} Z_i^{\mathfrak{s}_{1j} \mathfrak{s}_{1i}} f_j.$$

□

Applying the lemma to $y_1 = t\partial_1 \mathfrak{s}_{11} + \sum_{i \neq 1} c(x_1 - x_i)^{-1} \mathfrak{s}_{1i}$, we find that it is represented by a matrix \mathcal{L} with

$$\mathcal{L}_{kl} = \begin{cases} c(x_k - x_l)^{-1} & \text{for } k \neq l \\ t\partial_k & \text{for } k = l. \end{cases}$$

To calculate a Lax partner \mathcal{A} , we need to consider the action of \widehat{A} (2.11). Note that $(1 - \mathfrak{s}_{ij})e' = 0$ if $i, j > 1$; as a result, the action of \widehat{A} can be replaced by $ct \sum_{i \neq 1} (x_1 - x_i)^{-2} (1 - \mathfrak{s}_{1i})$. This element is W' -invariant, and its action on M' is again calculated from the above lemma. The result is:

$$\mathcal{A}_{kl} = \begin{cases} -ct(x_k - x_l)^{-2} & \text{for } k \neq l \\ ct \sum_{j \neq k}^n (x_j - x_k)^{-2} & \text{for } k = l. \end{cases}$$

Upon a substitution $t = -i\hbar$, $c = ig$, these \mathcal{L}, \mathcal{A} coincide with the Lax pair (1.5).

Remark 2.4. By analogy with [EG], we can also consider the action of x_1 on M' . The corresponding matrix is $\mathcal{X} = \text{diag}(x_1, \dots, x_n)$. We then have $\mathcal{X}\mathcal{L} - \mathcal{L}\mathcal{X} + i(g - \hbar)\mathbb{1}_n = ig\mathbf{v}\mathbf{w}$. This is a quantum version of the well-known relation from [KKS].

Remark 2.5. In a similar manner one can calculate Lax matrices for the Calogero–Moser systems associated to complex reflection groups, e.g., for the case $W = G(m, p, n)$. Note that classical Lax matrices for $W = G(m, 1, n)$ (generalised symmetric group) were found by a different method in [ChS].

3. CHEREDNIK OPERATORS AND LAX PAIRS FOR RELATIVISTIC CALOGERO–MOSER SYSTEMS

We start by outlining how the relativistic Calogero–Moser systems of Ruijsenaars–Schneider type can be constructed using affine Hecke algebras and Cherednik operators. This construction is due to Cherednik [C2, C3], and in the GL_n case it reproduces the quantum Ruijsenaars system [R1]. For other root systems some of the commuting Hamiltonians are expressed by Macdonald difference operators appearing in the theory of Macdonald polynomials [M1]. In the $C^\vee C_n$ case, the simplest Hamiltonian is the Koornwinder operator [Ko], and higher commuting Hamiltonians were constructed by van Diejen [vD2]. We therefore will refer to this case as Koornwinder–van Diejen system. Its interpretation in the framework of affine Hecke algebras can be found in [No, Sa, St].

We will be largely following Macdonald’s book [M2], see also [C4, Ki]. Our setting is not the most general (it corresponds to the case [M2, (1.4.1)]), but the method is exactly the same in all other cases. The $C^\vee C_n$ case is treated separately in Section §4.

3.1. Let R be a reduced, irreducible root system in a (complexified) Euclidean vector space V with an inner product denoted as $\langle -, - \rangle$, and W be the Weyl group of R , generated by the orthogonal reflections s_α , $\alpha \in R$. We write $R^\vee = \{\alpha^\vee\}$ for the dual system formed by the coroots $\alpha^\vee = 2\alpha/\langle \alpha, \alpha \rangle$. Let a_1, \dots, a_n be a fixed basis of simple roots in R , associated with a decomposition $R = R_+ \sqcup R_-$. We have the coroot and coweight lattices: $Q^\vee = \sum_{i=1}^n \mathbb{Z}a_i^\vee$ and $P^\vee = \sum_{i=1}^n \mathbb{Z}b_i$, where the fundamental coweights $b_i \in V$ are defined by $\langle a_i, b_j \rangle = \delta_{ij}$. We write $P_+^\vee := \sum_{i=1}^n \mathbb{Z}_{\geq 0}b_i$ for the cone of *dominant* coweights, .

The *affine Weyl group* is defined as $W_a = W \rtimes t(Q^\vee)$, where $t(Q^\vee)$ denotes the group of translations $t(\lambda)$, $\lambda \in Q^\vee$ acting on V by

$$t(\lambda)x = x - c\lambda, \quad (3.1)$$

where c is a fixed parameter. The *extended* affine Weyl group is $\widehat{W} := W \rtimes t(P^\vee)$. The group \widehat{W} acts naturally on the ring of meromorphic functions $\mathbb{C}(V)$ by

$$\widehat{w}f(x) = f(\widehat{w}^{-1}x) \quad \text{for } f(x) \in \mathbb{C}(V), \widehat{w} \in \widehat{W}. \quad (3.2)$$

In particular, a translation $t(\lambda)$, $\lambda \in P^\vee$ acts on functions by

$$t(\lambda)f(x) = f(x + c\lambda). \quad (3.3)$$

Writing $q = e^c$, we have $t(\lambda) = q^{\partial\lambda}$. We form a crossed product $\mathbb{C}(V) * \widehat{W}$ which we view as a subalgebra of $\text{End}_{\mathbb{C}}(\mathbb{C}(V))$, with $\mathbb{C}(V)$ acting on itself by multiplication. Inside $\mathbb{C}(V) * \widehat{W}$ we have an algebra \mathcal{D}_q generated by $\mathbb{C}(V)$ and $t(P^\vee)$; this is the algebra of difference operators on V . Clearly, $\mathbb{C}(V) * \widehat{W} \cong \mathcal{D}_q * W$, with every element admitting a unique presentation as

$$a = \sum_{w \in W} a_w w, \quad a_w \in \mathcal{D}_q. \quad (3.4)$$

Let \widehat{V} denote the space of affine-linear functions on V . We identify \widehat{V} with $V \oplus \mathbb{C}\delta$, where vectors in V are considered as linear functionals on V via the scalar product $\langle -, - \rangle$ and where $\delta \equiv c$ on V (so $e^\delta = e^c = q$). Let

$$R_a = \{\tilde{\alpha} = \alpha + k\delta, k \in \mathbb{Z}, \alpha \in R\} \subset \widehat{V} \quad (3.5)$$

be the affine root system associated with R . The action of \widehat{W} on $\widehat{V} \subset \mathbb{C}(V)$ permutes affine roots. For any $\tilde{\alpha} = \alpha + k\delta$ we have the orthogonal reflection with respect to the hyperplane $\tilde{\alpha}(x) = 0$ in V ,

$$s_{\tilde{\alpha}}(x) = x - \tilde{\alpha}(x)\alpha^\vee, \quad x \in V.$$

We extend the set of simple roots a_i to a basis in R_a by adding $a_0 = \delta - \varphi$, where φ is the highest root in R_+ . Then the reflections $s_i = s_{a_i}$, $i = 0, \dots, n$ generate the group W_a , and the *length* $l(w)$ of $w \in W_a$ is defined as the length l of a reduced decomposition

$$w = s_{i_1} \dots s_{i_l}, \quad \text{with } 0 \leq i_k \leq n. \quad (3.6)$$

Let Ω be the subgroup of the elements $\pi \in \widehat{W}$ which map the basis a_0, \dots, a_n to itself. It is known that Ω is an abelian group, isomorphic to P^\vee/Q^\vee , and the extended affine Weyl group is isomorphic to $W_a \rtimes \Omega$. Each $w \in \widehat{W}$ admits a unique presentation as $w = \tilde{w}\pi$ with $\tilde{w} \in W_a$ and $\pi \in \Omega$. We use this to extend the notion of the length from W_a to \widehat{W} by setting $l(\tilde{w}\pi) = l(\tilde{w})$, so $l(\pi) = 0$ for all $\pi \in \Omega$.

The *braid group* \mathfrak{B} of \widehat{W} is the group with generators T_w , $w \in \widehat{W}$, and relations

$$T_v T_w = T_{vw} \quad \text{if } l(v) + l(w) = l(vw).$$

Write $T_i := T_{s_i}$ for $i = 0, \dots, n$. Then for any reduced decomposition $w = s_{i_1} \dots s_{i_l}\pi$ we have $T_w = T_{i_1} \dots T_{i_l} T_\pi$. It follows that \mathfrak{B} is generated by T_i , $i = 0, \dots, n$ and T_π , $\pi \in \Omega$, subject to the following relations [M2, (3.1.6)]:

$$T_i T_j \dots = T_j T_i \dots \quad \text{for } i \neq j, \text{ with } m_{ij} \text{ factors on either side,} \quad (3.7)$$

$$T_\pi T_{\pi'} = T_{\pi\pi'} \quad \text{for } \pi, \pi' \in \Omega, \quad (3.8)$$

$$T_\pi T_i T_\pi^{-1} = T_j \quad \text{if } \pi s_i \pi^{-1} = s_j. \quad (3.9)$$

Here $m_{ij} = 2, 3, 4, 6$ is the order of $s_i s_j \in W_a$.

The braid group \mathfrak{B} contains an abelian subgroup $\{Y^\lambda \mid \lambda \in P^\vee\}$ [M2, 3.2]. Namely, for dominant λ we define $Y^\lambda = T_{i(\lambda)}$ and then extend this definition to all $\lambda \in P^\vee$ by setting $Y^\lambda = Y^\mu (Y^\nu)^{-1}$ whenever $\lambda = \mu - \nu$ with dominant μ, ν .

Choose nonzero parameters τ_i , $i = 0, \dots, n$ such that $\tau_i = \tau_j$ if s_i and s_j are conjugated in \widehat{W} . The (extended) *affine Hecke algebra* $\widehat{\mathfrak{H}}$ is the quotient of the group algebra $\mathbb{C}\mathfrak{B}$ by relations

$$(T_i - \tau_i)(T_i + \tau_i^{-1}) = 0, \quad i = 0, \dots, n. \quad (3.10)$$

The image of T_i (resp. T_w , Y^λ) in $\widehat{\mathfrak{H}}$ will be denoted by the same symbol T_i (resp. T_w , Y^λ). By [M2, (4.1.3)], the elements T_w , $w \in \widehat{W}$ form a \mathbb{C} -basis of $\widehat{\mathfrak{H}}$.

3.2. The algebra $\widehat{\mathfrak{H}}$ can be realized by difference-reflection operators, as was observed by Cherednik. This is an injective algebra map

$$\beta: \widehat{\mathfrak{H}} \rightarrow \mathcal{D}_q * W,$$

called the *basic representation* [M2, (4.3.10)]. To describe it, let us extend the set of parameters τ_i to τ_α , $\alpha \in R_a$ so that $\tau_\alpha = \tau_{w(\alpha)}$ for $w \in \widehat{W}$, and introduce functions \mathbf{c}_α as follows:

$$\mathbf{c}_\alpha = \frac{\tau_\alpha^{-1} - \tau_\alpha e^\alpha}{1 - e^\alpha}, \quad \alpha \in R_a. \quad (3.11)$$

Theorem 3.1 (cf. (4.3.10), (4.3.12) [M2]). *The extended affine Hecke algebra admits a faithful representation $\beta: \widehat{\mathfrak{H}} \rightarrow \mathcal{D}_q * W$ such that*

$$\beta: T_i \mapsto \tau_i + \mathbf{c}_i(x)(s_i - 1), \quad \mathbf{c}_i = \mathbf{c}_{a_i} \quad (i = 0, \dots, n), \quad (3.12)$$

$$\beta: T_\pi \mapsto \pi \quad \text{for all } \pi \in \Omega. \quad (3.13)$$

Let us identify $\widehat{\mathfrak{H}}$ with its image under β , so the affine Hecke algebra from now on will be viewed as a subalgebra of $\mathcal{D}_q * W$. The *Cherednik operators*, by definition, are the images of Y^λ under β . They form a commutative family of difference-reflection operators, and should be viewed as q -analogues of the Dunkl operators. In comparison, they are rather complicated. For example, for a dominant λ , one obtains Y^λ by first finding a reduced decomposition $t(\lambda) = s_{i_1} \dots s_{i_r} \pi$ and then writing the product $Y^\lambda = T_{i_1} \dots T_{i_r} T_\pi$ in the basic representation. Below we will often write Y^λ in terms of the elements $R(\alpha)$ (“ R -matrices”) defined by

$$R(\alpha) = \tau_\alpha s_\alpha + \mathbf{c}_\alpha(x)(1 - s_\alpha), \quad \alpha \in R_a. \quad (3.14)$$

The following property of these elements is important:

$$wR(\alpha)w^{-1} = R(w(\alpha)), \quad \text{for any } w \in W_a.$$

Using this and the fact that $R(a_i) = T_i s_i$ for $i = 0, \dots, n$, it is straightforward to rewrite Y^λ in terms of $R(\alpha)$ instead of T_i .

The commutative subalgebra generated by the Cherednik operators will be denoted as $\mathbb{C}[Y]$, so elements $f(Y) \in \mathbb{C}[Y]$ are arbitrary linear combinations of Y^λ , $\lambda \in P^\vee$. Inside $\mathbb{C}[Y]$ we have a subalgebra $\mathbb{C}[Y]^W$, spanned by the orbitsums $f = \sum_{\mu \in W\lambda} Y^\mu$.

3.3. The (finite) Hecke algebra \mathfrak{H} of W is a subalgebra of $\widehat{\mathfrak{H}}$, generated by T_i , $i = 1, \dots, n$. That is, \mathfrak{H} is generated by T_1, \dots, T_n which satisfy the relations (3.7), (3.10). We have an isomorphism $\widehat{\mathfrak{H}} \cong \mathfrak{H} \otimes \mathbb{C}[Y]$, as vector spaces. The cross-relations between T_i and Y^λ are the so-called Lusztig relations [M2, (4.2.4)]. They imply the following property, cf. [M2, (4.2.8)]:

$$T_i f(Y) = f(Y) T_i \quad \text{for all } f \in \mathbb{C}[Y] \text{ such that } s_i \cdot f = f, \quad (3.15)$$

where the action of s_i on $f \in \mathbb{C}[Y]$ is by $s_i \cdot Y^\lambda = Y^{s_i(\lambda)}$.

For any $w \in W$ with a reduced decomposition $w = s_{i_1} \dots s_{i_r}$, we define $\tau_w = \tau_{i_1} \dots \tau_{i_r}$. The following is a Hecke-algebra analogue of the symmetrizer (2.6):

$$e_\tau = \frac{1}{\sum_{w \in W} \tau_w^2} \sum_{w \in W} \tau_w T_w. \quad (3.16)$$

By [M2, 5.5.17],

$$T_i e_\tau = e_\tau T_i = \tau_i e_\tau, \quad e_\tau^2 = e_\tau. \quad (3.17)$$

Let us also define parabolic symmetrizers. Given $I \subset \{1, \dots, n\}$, we define $W_I \subset W$ as the subgroup generated by s_i with $i \in I$. It is known that W_I is isomorphic to a Weyl group with Dynkin diagram obtained by removing the vertices $j \notin I$ from the diagram of W . Similarly, define a parabolic subalgebra $\mathfrak{H}_I \subset \mathfrak{H}$ as the subalgebra generated by T_i with $i \in I$. By [GeP, 4.4.7], \mathfrak{H}_I is spanned by T_w with $w \in W_I$ and it is isomorphic to the Hecke algebra of W_I with parameters τ_i , $i \in I$. As a result, if we define

$$e_{\tau, I} = \frac{1}{\sum_{w \in W_I} \tau_w^2} \sum_{w \in W_I} \tau_w T_w, \quad (3.18)$$

then

$$T_i e_{\tau, I} = e_{\tau, I} T_i = \tau_i e_{\tau, I} \quad \forall i \in I, \quad e_{\tau, I}^2 = e_{\tau, I}.$$

3.4. The Hamiltonians of the quantum relativistic Calogero–Moser system based on a root system R are given by certain difference operators. The following construction for them was given by Cherednik. We start by picking an arbitrary $f(Y) \in \mathbb{C}[Y]^W$; then there is a unique difference operator $L_f \in \mathcal{D}_q$ such that $f(Y)e = L_f e$, where e is the symmetrizer (2.6). By (3.15), $(T_i - \tau_i)f(Y) = f(Y)(T_i - \tau_i)$. From (3.12) we have $T_i - \tau_i = \mathbf{c}_i(x)(s_i - 1)$, therefore,

$$\mathbf{c}_i(x)(s_i - 1)L_f e = (T_i - \tau_i)f(Y)e = f(Y)(T_i - \tau_i)e = f(Y)\mathbf{c}_i(x)(s_i - 1)e = 0. \quad (3.19)$$

As a result, $(s_i - 1)L_f e = 0$, from which $L_f^{s_i} = L_f$ for all i , i.e. L_f is a W -invariant difference operator, with $eL_f e = eL_f = L_f e = f(Y)e = ef(Y)e$. (Warning: $f(Y)e \neq ef(Y)e$!) These operators pairwise commute, for if $L_f, L_g \in \mathcal{D}_q^W$ are constructed from $f, g \in \mathbb{C}[Y]^W$ then

$$L_f L_g e = L_f e L_g e = f(Y)e L_g e = f(Y)L_g e = f(Y)g(Y)e,$$

which is invariant under exchanging f and g . This gives an algebra map

$$\mathbb{C}[Y]^W \rightarrow \mathcal{D}_q^W, \quad f \mapsto L_f, \quad (3.20)$$

hence a quantum integrable system with commuting Hamiltonians $L_f, f \in \mathbb{C}[Y]^W$. In the case $R = A_{n-1}$, this is the trigonometric Ruijsenaars system [R1]. In general, the difference operators $L_f, f \in \mathbb{C}[Y]^W$ are complicated; some of them are known explicitly [M1, Ko, vD2, vDI, vDE].

3.5. Quantum Lax pairs are now constructed in the same way as before. First, the algebra \mathcal{D}_q of difference operators acts naturally on $\mathbb{C}(V)$, with translations $t(\lambda), \lambda \in P^\vee$ acting by (3.3). Consider the induced module

$$M = \text{Ind}_{\mathcal{D}_q}^{\mathcal{D}_q * W} \mathbb{C}(V). \quad (3.21)$$

We have $M \cong \mathbb{C}W \otimes \mathbb{C}(V)$ as vector spaces, so $\text{End}_{\mathbb{C}}(M)$ can be identified with *operator-valued* matrices of size $|W|$. As a result, the (left) action of $\mathcal{D}_q * W$ on M gives a representation

$$\mathcal{D}_q * W \rightarrow \text{Mat}(|W|, \mathcal{D}_q). \quad (3.22)$$

Again, any $a \in \mathcal{D}_q^W$ is represented by $a\mathbb{1}$.

Now pick a Cherednik operator $Y^\lambda, \lambda \in P^\vee$ and choose an element $f(Y), f \in \mathbb{C}[Y]^W$, writing it as $f(Y) = \sum_{w \in W} a_w w$ with $a_w \in \mathcal{D}_q$. Then

$$f(Y) = L_f + \widehat{A}, \quad \text{where} \quad L_f = \sum_{w \in W} a_w, \quad \widehat{A} := \sum_{w \in W} a_w (w - 1). \quad (3.23)$$

Note that L_f is one of the Hamiltonians (3.20), so $L_f \in \mathcal{D}_q^W$. We now set $\mathcal{L}, \mathcal{H}, \mathcal{A}$ to be the images of Y^λ, L_f and \widehat{A} , respectively, under (3.22). Since Y^λ and $f(Y) = L_f + \widehat{A}$ commute, we obtain (1.3) with $\mathcal{H} = L_f \mathbb{1}$, that is, a quantum Lax pair of size $|W|$.

3.6. The classical limit corresponds to $q = e^c \rightarrow 1$, and the procedure is similar to §2.3. Namely, we set $c = -i\hbar\beta$, with some fixed β , and consider the algebra

$$A_\hbar * W = \mathbb{C}(V)[t_1^{\pm 1}, \dots, t_n^{\pm 1}][[\hbar]] * W, \quad t_k := t(b_k),$$

where b_k are the fundamental coweights. We have

$$[t_k, f] = \sum_{l=1}^{\infty} (-i\hbar\beta)^l \partial_k^l(f) t_k, \quad \forall f \in \mathbb{C}(V),$$

where we denote $\partial_k := \partial_{b_k}$.

Consider a commutative algebra $A_0 = \mathbb{C}(V)[P^\vee]$. Elements of $\mathbb{C}[P^\vee]$ are linear combinations with coefficients in $\mathbb{C}(V)$ of $e^\lambda, \lambda \in P^\vee$. We will view e^λ as a function of the classical momenta $p \in V$ by setting $e^\lambda := e^{\beta p_\lambda}$, where p_λ is the momentum in direction λ . Writing $p_k := p_{b_k}$, we have an algebra isomorphism

$$\eta_0 : A_\hbar * W / (\hbar A_\hbar * W) \rightarrow A_0 * W, \quad f \mapsto f, \quad t_k \mapsto e^{\beta p_k}, \quad w \mapsto w.$$

We may view A_\hbar (resp. $A_\hbar * W$) as a formal deformation of A_0 (resp. $A_0 * W$). The algebra A_0 is commutative, with the Poisson bracket determined by $[a, b] = i\hbar\{\eta_0(a), \eta_0(b)\} + o(\hbar)$ for $a, b \in A_\hbar$. Explicitly, $\{e^{\beta p_k}, f\} = -\beta \partial_k(f) e^{\beta p_k}$. For $a \in A_\hbar * W$, we call $\eta_0(a)$ the *classical limit* of a .

The classical Cherednik operators $Y_c^\lambda := \eta_0(Y^\lambda)$ can be computed by using the classical version of the basic representation, cf. [Ob, Section 5]; let us write $f(Y)_c$ for the classical version of $f(Y) \in \mathbb{C}[Y]$. We then have the following analogue of Lemma 2.2.

Lemma 3.2. *For any $f(Y) \in \mathbb{C}[Y]^W$, when writing $f(Y)_c = \sum_{w \in W} a_w w$ with $a_w \in A_0$, we have $a_w = 0$ for $w \neq \text{id}$.*

Proof. Indeed, by [Ob, Lemma 5.1] the algebra $\mathbb{C}[Y]^W$ belongs to the center of the double affine Hecke algebra at the classical level $q = 1$. Therefore, the same proof as in [EFMV, Lemma 2.2] applies. \square

This tells us that the classical limit of \widehat{A} in (3.23) is zero. Therefore, $\hbar^{-1}\widehat{A}$ has a well-defined classical limit, and so the classical Lax pair is obtained in the same way as in §2.3.

3.7. To get a Lax pair of a smaller size, we take λ to be on the boundary of the Weyl chamber, i.e. $\langle \lambda, a_i \rangle \geq 0$ for $i = 1, \dots, n$, with at least one i such that $\langle \lambda, a_i \rangle = 0$. Set $I = \{i \mid \langle a_i, \lambda \rangle = 0\}$; then the stabiliser $W' = W_I$ of λ is generated by s_i with $i \in I$. Write

$$e' = \frac{1}{|W'|} \sum_{w \in W'} w, \quad M' = e'M \cong e'CW \otimes \mathbb{C}(V). \quad (3.24)$$

By (3.15), we have $T_i Y^\lambda = Y^\lambda T_i$ for all $i \in I$. Then, similarly to (3.19), we have $(s_i - 1)Y^\lambda e' = 0$ for $i \in I$, implying $e'Y^\lambda e' = Y^\lambda e'$. This tells us that the action of Y^λ preserves $M' = e'M$. Similarly, $f(Y)$ for any $f \in \mathbb{C}[Y]^W$ preserves M' . Therefore, the operators Y^λ , $f(Y) = L_f + \hat{A}$, as well as L_f , can be restricted onto M' , resulting in a quantum Lax pair \mathcal{L} , \mathcal{A} of size $|W/W'|$.

3.8. There is a q -analogue of the formula (1.11) for constructing first integrals. It generalises a formula discovered in the GL_n -case by Nazarov and Sklyanin [NS], see § 3.10 below. As before, we choose a dominant $\lambda \in P^\vee$ and write $W' = W_I$ for the stabiliser of λ and $e_{\tau, I}$, e' for the symmetrizers (3.18), (3.24). Denote by $R' \subset R$ the root system of W' , and choose coset representatives w_i for $W' \backslash W$. We write elements of $M' = e'M$ as linear combinations of $e'w_i f_i$ with $f_i \in \mathbb{C}(V)$. Let \mathcal{L} be a matrix of size $|W/W'|$ which represents the action of Y^λ on M' .

Recall the functions $\mathbf{c}_\alpha(x)$ (3.11). Introduce a pair of row/column vectors with $l := |W/W'|$ components as follows:

$$\mathbf{v} = (1, \dots, 1)^T, \quad \mathbf{u} = (\phi_1, \dots, \phi_l), \quad \phi_i(x) := \phi(w_i x), \quad \text{where } \phi(x) = \prod_{\alpha \in R_+ \setminus R'_+} \mathbf{c}_\alpha(-x). \quad (3.25)$$

Proposition 3.3. *The difference operators $\hat{H}_k = \mathbf{u} \mathcal{L}^k \mathbf{v}$, $k \in \mathbb{Z}$ belong to the algebra of quantum Hamiltonians L_f , $f \in \mathbb{C}[Y]^W$.*

Proof. First, we claim that for any $f \in \mathbb{C}[Y]^W$ the elements $e_\tau Y^\lambda e$ and $ef(Y)e = L_f e = eL_f$ commute:

$$[e_\tau Y^\lambda e, ef(Y)e] = 0. \quad (3.26)$$

Indeed, by (3.17) and (3.12) we have $0 = (T_i - \tau_i)e_\tau = \mathbf{c}_i(x)(s_i - 1)e_\tau$. Thus, $s_i e_\tau = e_\tau$ for all i , implying

$$we_\tau = e_\tau \quad \forall w \in W. \quad (3.27)$$

Therefore, $ee_\tau = e_\tau$. Next, recall that by (3.15) any $f(Y) \in \mathbb{C}[Y]^W$ commutes with all T_i , so $e_\tau f(Y) = f(Y)e_\tau$. Now we have

$$(e_\tau Y^\lambda e)(eL_f e) = e_\tau Y^\lambda L_f e = e_\tau Y^\lambda f(Y)e,$$

and similarly,

$$(eL_f e)(e_\tau Y^\lambda e) = f(Y)ee_\tau Y^\lambda e = f(Y)e_\tau Y^\lambda e = e_\tau f(Y)Y^\lambda e,$$

which implies (3.26).

Next, by [M2, (5.5.14), (5.5.15)] we have $e_\tau = e\mathbf{c}_+$ where $\mathbf{c}_+(x) = \prod_{\alpha \in R_+} \mathbf{c}_\alpha(-x)$ (up to a constant factor). Similarly, $e_{\tau, I} = e'\mathbf{c}'_+$ with $\mathbf{c}'_+(x) = \prod_{\alpha \in R'_+} \mathbf{c}_\alpha(-x)$. As a result, the action of the symmetrizer e_τ on $M' = e'M$ is calculated as follows:

$$e_\tau e' = e_\tau e_{\tau, I} / \mathbf{c}'_+ = e_\tau / \mathbf{c}'_+ = e\mathbf{c}_+ / \mathbf{c}'_+ = e\phi, \quad \phi(x) = \prod_{\alpha \in R_+ \setminus R'_+} \mathbf{c}_\alpha(-x). \quad (3.28)$$

The function ϕ is W' -invariant. If $f = \sum_i e'w_i f_i$ is an element of M' , then $\phi f = \sum_i e'w_i(\phi)^{w_i^{-1}} f_i$, where $\phi^{w_i^{-1}}(x) = \phi(w_i x)$. It follows that, up to a constant factor, $e_\tau = e\phi$ acts on M' by a matrix $\mathbf{v}\mathbf{u}$ as given in (3.25), and e acts as $\mathbf{v}\mathbf{w}$, where \mathbf{w} is a row of ones. We conclude that, up to a constant factor, $e_\tau Y^{k\lambda} e$ acts as $\mathbf{v}\mathbf{u} \mathcal{L}^k \mathbf{v}\mathbf{w} = (\mathbf{u} \mathcal{L}^k \mathbf{v})e$, and it is clear from the construction that $\mathbf{u} \mathcal{L}^k \mathbf{v}$ is W -invariant. Meanwhile, $ef(Y)e$ acts on M' as $L_f e$. Hence, $\mathbf{u} \mathcal{L}^k \mathbf{v}$ and L_f commute by (3.26).

This proves that each of $\hat{H}_k = \mathbf{u} \mathcal{L}^k \mathbf{v}$ is a W -invariant difference operator commuting with all L_f , $f \in \mathbb{C}[Y]^W$. By [LS, Theorem 3.15], \hat{H}_k must belong to the algebra of the Hamiltonians L_f , $f \in \mathbb{C}[Y]^W$. In particular, \hat{H}_k pairwise commute. \square

3.9. Let us calculate a quantum Lax pair for the trigonometric Ruijsenaars system. This corresponds to the case of a root system $R = A_{n-1}$ (or more precisely, to the \mathbf{GL}_n -case). This will provide an alternative derivation of some of the results of Nazarov and Sklyanin [NS]. Let us first describe the Cherednik operators in this case.

We take $V = \mathbb{C}^n$, with the orthonormal basis $\epsilon_1, \dots, \epsilon_n$ and the associated coordinates x_1, \dots, x_n . As in §2.6, the group $W = \mathfrak{S}_n$ acts on V by permuting the basis vectors, and we write $W' = \mathfrak{S}_{n-1}$ for the subgroup fixing ϵ_1 . As before, e, e' are the corresponding symmetrizers. The roots in R are $\alpha = \epsilon_i - \epsilon_j$ with $i \neq j$. The role of P^\vee in the \mathbf{GL}_n -case is played by the lattice $\Lambda = \sum_{i=1}^n \mathbb{Z}\epsilon_i$. There is only one coupling parameter τ , so $\tau_\alpha = \tau$ for all α . Fix $q = e^c$ and consider the algebra of difference operators $\mathcal{D}_q = \mathbb{C}(x) \rtimes t(\Lambda)$, generated by $\mathbb{C}(x)$ and the shift operators $t(\epsilon_k) = q^{\partial_k}$. For $i \neq j$ we set $R_{ij} := R(\epsilon_i - \epsilon_j)$ according to (3.14). Explicitly,

$$R_{ij} = a_{ij} + b_{ij}\mathfrak{s}_{ij}, \quad a_{ij} = a(x_i - x_j), \quad b_{ij} = b(x_i - x_j), \quad (3.29)$$

where

$$a(z) = \frac{\tau^{-1} - \tau e^z}{1 - e^z}, \quad b(z) = \frac{\tau - \tau^{-1}}{1 - e^z}. \quad (3.30)$$

The Cherednik operators are the following commuting elements of $\mathcal{D}_q * W$ [C2, BGHP], cf. [NS, (2.7)]:

$$Y_i = R_{i,i+1}R_{i,i+2} \dots R_{i,n} t(\epsilon_i) R_{i-1,i}^{-1} \dots R_{1,i}^{-1} \quad (i = 1, \dots, n). \quad (3.31)$$

The operators L_f , $f \in \mathbb{C}[Y_1^{\pm 1}, \dots, Y_n^{\pm 1}]^{\mathfrak{S}_n}$ are obtained according to §3.4. They can be given explicitly [R1, M1]; the simplest one is

$$L_f = \sum_{i=1}^n \left(\prod_{l \neq i} a_{il} \right) t(\epsilon_i), \quad f = Y_1 + \dots + Y_n. \quad (3.32)$$

A Lax matrix will be constructed from $Y_1 = Y^{\epsilon_1}$:

$$Y_1 = R_{12}R_{13} \dots R_{1n} t(\epsilon_1). \quad (3.33)$$

As explained in §3.7, Y_1 acts on $M' = e'M$. We have the following result, proved by Nazarov and Sklyanin (cf. a similar statement in the elliptic case in [KH1, Lemma 4.4]).

Lemma 3.4 ([NS], Proposition 2.4). *Upon restricting onto M' , Y_1 coincides with*

$$\left(A + \sum_{i \neq 1}^n B_i \mathfrak{s}_{1i} \right) t(\epsilon_1), \quad \text{where } A = \prod_{l \neq 1}^n a_{1l}, \quad B_i = b_{1i} \prod_{l \neq 1, i} a_{il}.$$

In [NS] this is proved by a cleverly organised induction. We give a simpler proof using symmetry arguments.

Proof. Since $t(\epsilon_1)$ commutes with e' , it is enough to show that

$$R_{12}R_{13} \dots R_{1n} e' = \left(A + \sum_{i \neq 1}^n B_i \mathfrak{s}_{1i} \right) e', \quad (3.34)$$

with A, B_i as stated. Using the definition of R_{1i} , we expand the product:

$$R_{12}R_{13} \dots R_{1n} = \sum_{2 \leq i_1 < \dots < i_k \leq n} g_{i_1 \dots i_k} \mathfrak{s}_{1i_1} \dots \mathfrak{s}_{1i_k}, \quad \text{with some } g_{i_1 \dots i_k} \in \mathbb{C}(x).$$

It is easy to see that $\mathfrak{s}_{1i_1} \dots \mathfrak{s}_{1i_k} e'$ equals $\mathfrak{s}_{1i_k} e'$ (or e' if the set $\{i_1, \dots, i_k\}$ is empty). It follows that in (3.34) one has $A = \prod_{l \neq 1}^n a_{1l}$. Similarly, B_2 is determined from

$$B_2 \mathfrak{s}_{12} = (b_{12} \mathfrak{s}_{12}) a_{13} \dots a_{1n} = b_{12} \left(\prod_{l \neq 1, 2} a_{2l} \right) \mathfrak{s}_{12}.$$

Finally, using Lemma 2.3(1) we find that $B_i = B_2^{\mathfrak{s}_{2i}} = b_{1i} \prod_{l \neq 1, i} a_{il}$, as needed. \square

Now a Lax matrix \mathcal{L} can be calculated using Lemma 2.3. This gives

$$\mathcal{L}_{ij} = \begin{cases} \left(\prod_{l \neq j} a_{jl} \right) t(\epsilon_j) & \text{for } i = j \\ \left(\prod_{l \neq i, j} a_{jl} \right) b_{ij} t(\epsilon_j) & \text{for } i \neq j. \end{cases}$$

This is the quantum Lax matrix obtained in [NS] (where it is denoted as Z), see also [Ha]. Here $t(\epsilon_j) = q^{\partial_j} = e^{\beta \hat{p}_j}$, assuming $q = e^{-i\hbar\beta}$. The classical Lax matrix is obtained by replacing \hat{p}_j with the classical momentum p_j :

$$L_{ij} = \begin{cases} \left(\prod_{l \neq j} a_{jl} \right) e^{\beta p_j} & \text{for } i = j \\ \left(\prod_{l \neq i, j} a_{jl} \right) b_{ij} e^{\beta p_j} & \text{for } i \neq j. \end{cases}$$

It is equivalent to the Lax matrix [R1, (3.19)] for the trigonometric Ruijsenaars–Schneider system.

3.10. We can now apply Proposition 3.3. In our case we have $R_+ = \{\epsilon_i - \epsilon_j \mid 1 \leq i < j \leq n\}$ and $R'_+ = \{\epsilon_i - \epsilon_j \mid 2 \leq i < j \leq n\}$. The function $\mathbf{c}_\alpha(-x)$ for $\alpha = \epsilon_i - \epsilon_j$ coincides with a_{ji} as given in (3.30). We can write down the function ϕ and the row/column vectors \mathbf{u}, \mathbf{v} (3.25):

$$\phi = \prod_{l \neq 1}^n a_{l1}, \quad \phi_i = \phi^{\sigma^{1i}} = \prod_{l \neq i}^n a_{li}, \quad \mathbf{u} = (\phi_1, \dots, \phi_n), \quad \mathbf{v} = (1, \dots, 1)^T.$$

According to Proposition 3.3, the difference operators $\mathbf{u}\mathcal{L}^j\mathbf{v}$, $j \in \mathbb{Z}$ belong to the commutative algebra of the Macdonald–Ruijsenaars operators L_f , $f \in \mathbb{C}[Y_1^{\pm 1}, \dots, Y_n^{\pm 1}]^{\mathfrak{S}_n}$.

We can compare this with one of the results of Nazarov and Sklyanin. Namely, [NS, Corollary 2.6] claims that the quantities UZ^jE belong to the algebra of the Macdonald–Ruijsenaars operators. Here Z and E coincide with our \mathcal{L} and \mathbf{v} , but the row vector $U = (U_1, \dots, U_n)$ is different from our \mathbf{u} (in particular, the expressions for U_i contain the shift operators $t(\epsilon_i)$, see [NS, (2.19)]). These two formulas may seem different, however, it can be checked that up to a constant factor, $U = \mathbf{u}\mathcal{L}$. Therefore, $UZ^jE = \mathbf{u}\mathcal{L}^{j+1}\mathbf{v}$, so both results agree.

3.11. A Lax partner \mathcal{A} for the Hamiltonian L_f (3.32) corresponding to $f = Y_1 + \dots + Y_n$ can be calculated by a similar method. We skip the details (see a similar calculation for the elliptic case in §6.3), and will only state the result:

$$\mathcal{A}_{ij} = \left(\prod_{l \neq i, j} a_{jl} \right) (b_{ij} t(\epsilon_j) - t(\epsilon_j) b_{ij}) \quad \text{for } i \neq j, \quad \mathcal{A}_{ii} = - \sum_{j \neq i} \mathcal{A}_{ij}.$$

To calculate its classical counterpart, we use that $t(\epsilon_j) = e^{-i\hbar\beta\partial_j}$ and

$$b_{ij} t(\epsilon_j) - t(\epsilon_j) b_{ij} = i\hbar\beta \frac{\partial b_{ij}}{\partial x_j} t(\epsilon_j) + o(\hbar).$$

Therefore, the classical limit A of $(i\hbar)^{-1}\mathcal{A}$ is given by

$$A_{ij} = \left(\prod_{l \neq i, j} a_{jl} \right) \beta \frac{\partial b_{ij}}{\partial x_j} e^{\beta p_j} \quad \text{for } i \neq j, \quad A_{ii} = - \sum_{j \neq i} A_{ij}.$$

4. LAX MATRIX FOR THE KOORNWINDER–VAN DIEJEN SYSTEM

4.1. Let us describe the Cherednik operators in the case $C^V C_n$, corresponding to the Koornwinder–van Diejen system [Ko, vD2]. We will follow [St] fairly closely, so the reader should consult that paper for further details. Let $V = \mathbb{C}^n$ with the standard orthonormal basis $\{\epsilon_i\}_{i=1}^n$ and the associated coordinates x_i . Let R be the root system of type C_n ,

$$R = \{\pm 2\epsilon_i \mid 1 \leq i \leq n\} \cup \{\pm\epsilon_i \pm \epsilon_j \mid 1 \leq i < j \leq n\}. \quad (4.1)$$

The Weyl group $W = \mathfrak{S}_n \times \{\pm 1\}^n$ of R consists of the transformations that permute the basis vectors ϵ_i and change their signs arbitrarily. As in §3.1, we write $\widehat{V} = V \oplus \mathbb{C}\delta$ for the space of affine-linear functions on V , with $\delta \equiv c$ on V . Let R_a be the affine root system associated with R (3.5). We choose a basis of simple roots in R_a ,

$$a_0 = \delta - 2\epsilon_1, \quad a_i = \epsilon_i - \epsilon_{i+1} \quad (i = 1, \dots, n-1), \quad a_n = 2\epsilon_n.$$

The affine Weyl group W_a of R is generated by $s_i = s_{a_i}$, $i = 0, \dots, n$. The action of the generators in coordinates x_i on V looks as follows:

$$\begin{aligned} s_0(x_1, \dots, x_n) &= (c - x_1, x_2, \dots, x_n), \\ s_i(x_1, \dots, x_n) &= (x_1, \dots, x_{i-1}, x_{i+1}, x_i, \dots, x_n) \quad (i = 1, \dots, n-1), \\ s_n(x_1, \dots, x_n) &= (x_1, \dots, x_{n-1}, -x_n). \end{aligned} \quad (4.2)$$

We have $W_a \cong W \ltimes \Lambda$, where $\Lambda = \sum_{i=1}^n \mathbb{Z}\epsilon_i$ is the coroot lattice of R , acting on V by (3.1). We consider the associated action (3.2), (3.3) of W_a on $\mathbb{C}(V)$, and form the algebra $\mathbb{C}(V) * W_a \cong \mathcal{D}_q * W$ of reflection-difference operators on V .

The affine Hecke algebra $\widehat{\mathfrak{H}}$ associated with W_a is generated by T_0, \dots, T_n subject to the following relations:

$$T_i T_{i+1} T_i T_{i+1} = T_{i+1} T_i T_{i+1} T_i \quad (i = 0, i = n-1), \quad (4.3)$$

$$T_i T_{i+1} T_i = T_{i+1} T_i T_{i+1} \quad (i = 1, \dots, n-2), \quad (4.4)$$

$$T_i T_j = T_j T_i, \quad |i - j| \geq 2, \quad (4.5)$$

$$(T_i - \tau_i)(T_i + \tau_i^{-1}) = 0 \quad (i = 0, \dots, n), \quad (4.6)$$

where τ_i are deformation parameters, with $\tau_1 = \dots = \tau_{n-1} = \tau$.

The basic representation $\beta : \widehat{\mathfrak{H}} \rightarrow \mathcal{D}_q * W$ is due to Noumi [No]. To describe it, we choose two additional parameters, τ_0^\vee, τ_n^\vee . It will be convenient to introduce parameters τ_α and functions \mathbf{c}_α for $\alpha \in R_a$ as follows (cf. [M2, (4.2.2), (4.3.9)]):

$$\tau_\alpha = \tau, \quad \mathbf{c}_\alpha = \frac{\tau^{-1} - \tau e^\alpha}{1 - e^\alpha} \quad \text{for } \alpha = k\delta \pm \epsilon_i \pm \epsilon_j \quad (k \in \mathbb{Z}, i \neq j), \quad (4.7)$$

$$\tau_\alpha = \tau_0, \quad \mathbf{c}_\alpha = \tau_0^{-1} \frac{(1 - \tau_0 \tau_0^\vee e^{\alpha/2})(1 + \tau_0 (\tau_0^\vee)^{-1} e^{\alpha/2})}{(1 - e^\alpha)} \quad \text{for } \alpha = (2k+1)\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}), \quad (4.8)$$

$$\tau_\alpha = \tau_n, \quad \mathbf{c}_\alpha = \tau_n^{-1} \frac{(1 - \tau_n \tau_n^\vee e^{\alpha/2})(1 + \tau_n (\tau_n^\vee)^{-1} e^{\alpha/2})}{(1 - e^\alpha)} \quad \text{for } \alpha = 2k\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}). \quad (4.9)$$

With this notation, we define β by setting

$$\beta : T_i \mapsto \tau_i + \mathbf{c}_{a_i}(s_i - 1), \quad i = 0, \dots, n,$$

and extend it to the whole of $\widehat{\mathfrak{H}}$ by multiplicativity, see [St, Theorem 3.2]. This defines a subalgebra of $\mathcal{D}_q * W$, depending on five parameters $\tau_0, \tau_0^\vee, \tau_n, \tau_n^\vee, \tau$.

By [St, (3.5)], inside $\widehat{\mathfrak{H}}$ we have a commutative subalgebra $\mathbb{C}[Y]$ generated by $Y_i^{\pm 1}$, where

$$Y_i = T_i \dots T_{n-1} T_n T_{n-1} \dots T_1 T_0 T_1^{-1} T_2^{-1} \dots T_{i-1}^{-1}, \quad i = 1, \dots, n.$$

The Hamiltonians of the Koornwinder–van Diejen system can be obtained as in §3.4, by taking symmetric combinations of $Y_i^{\pm 1}$. For $f(Y) = \sum_{i=1}^n (Y_i + Y_i^{-1})$ this reproduces the Koornwinder operator [Ko]; explicit formulas for the higher Hamiltonians can be found in [vD2].

4.2. Our goal to calculate a quantum Lax matrix corresponding to Y_1 . We will use the following notation for the reflections in W : \mathfrak{s}_{ij} is an elementary transposition of x_i and x_j , \mathfrak{s}_i is a sign reversal in the i th direction, and $\mathfrak{s}_{ij}^+ = \mathfrak{s}_{ij} \mathfrak{s}_i \mathfrak{s}_j$, which acts by $x_i \mapsto -x_j$, $x_j \mapsto -x_i$. Comparison with (4.2) gives $s_i = \mathfrak{s}_{i, i+1}$ for $i = 1, \dots, n-1$, $s_n = \mathfrak{s}_n$, $s_0 = \mathfrak{s}_1 t(\epsilon_1)$.

We have $Y_1 = Y^{\epsilon_1}$, so following §3.7 we consider the stabiliser W' of $\lambda = \epsilon_1$. This is the subgroup of signed permutations of x_2, \dots, x_n , and the Lax matrix will be of size $|W|/|W'| = 2n$. To calculate it, we need to determine the action of Y_1 on $M' = e'M$. By [St, (4.2)],

$$Y_1 = R(\epsilon_1 - \epsilon_2) R(\epsilon_1 - \epsilon_3) \dots R(\epsilon_1 - \epsilon_n) R(2\epsilon_1) R(\epsilon_1 + \epsilon_n) \dots R(\epsilon_1 + \epsilon_2) R(\delta + 2\epsilon_1) t(\epsilon_1). \quad (4.10)$$

Here $R(\alpha)$ are defined by (3.14) together with (4.7)–(4.9).

Let us introduce some shorthand notation. In addition to a_{ij}, b_{ij}, R_{ij} (3.29), (3.30), we define $a_{ij}^+ = a(x_i + x_j)$, $b_{ij}^+ = b(x_i + x_j)$, $a_{ij}^- = a(-x_i - x_j)$, $b_{ij}^- = b(-x_i - x_j)$, $R_{ij}^+ = a_{ij}^+ + b_{ij}^+ \mathfrak{s}_{ij}^+$.

Let us also introduce functions $u, v, \tilde{u}, \tilde{v}$ as follows:

$$u(z) = \tau_n^{-1} \frac{(1 - \tau_n \tau_n^\vee e^z)(1 + \tau_n (\tau_n^\vee)^{-1} e^z)}{1 - e^{2z}}, \quad v(z) = \tau_n - u(z), \quad (4.11)$$

$$\tilde{u}(z) = \tau_0^{-1} \frac{(1 - \tau_0 \tau_0^\vee q^{1/2} e^z)(1 + \tau_0 (\tau_0^\vee)^{-1} q^{1/2} e^z)}{1 - qe^{2z}}, \quad \tilde{v}(z) = \tau_0 - \tilde{u}(z). \quad (4.12)$$

Here $q = e^c$. Below we will use u_i, u_i^- to denote $u_i = u(x_i)$ and $u(-x_i)$, respectively, and similarly for v, \tilde{u}, \tilde{v} . With this notation, we have:

$$Y_1 = R_{12} R_{13} \dots R_{1n} R(2\epsilon_1) R_{1n}^+ \dots R_{12}^+ R(\delta + 2\epsilon_1) t(\epsilon_1), \\ R(2\epsilon_1) = u_1 + v_1 \mathfrak{s}_1, \quad R(\delta + 2\epsilon_1) t(\epsilon_1) = \tilde{u}_1 t(\epsilon_1) + \tilde{v}_1 \mathfrak{s}_1.$$

Lemma 4.1. *Let $e'' = \frac{1}{n!} \sum_{w \in \mathfrak{S}_{n-1}} w$, where \mathfrak{S}_{n-1} is the group of permutations on $\{2, \dots, n\}$, viewed as a subgroup of $W = \mathfrak{S}_n \times \mathbb{Z}_2^n$. The operators $\mathcal{R} = R_{12} \dots R_{1n}$ and $\mathcal{R}^+ = R_{1n}^+ \dots R_{12}^+$ preserve the subspace $M'' = e''M$. Their restriction onto M'' is calculated as follows:*

$$\mathcal{R}|_{M''} = U + \sum_{i \neq 1}^n V_i \mathfrak{s}_{1i}, \quad \mathcal{R}^+|_{M''} = U^+ + \sum_{i \neq 1}^n V_i^+ \mathfrak{s}_{1i}^+, \quad \text{where} \\ U = \prod_{l \neq 1}^n a_{1l}, \quad V_i = b_{1i} \prod_{l \neq 1, i} a_{il}, \quad U^+ = \prod_{l \neq 1}^n a_{1l}^+, \quad V_i^+ = b_{1i}^+ \prod_{l \neq 1, i} a_{li}.$$

Proof. For \mathcal{R} this has been shown in Lemma §3.10. The statement for \mathcal{R}^+ follows from the fact that $\mathcal{R}^+ = \mathcal{R}^\omega$, where $\omega \in W'$ is given by

$$\omega : V \rightarrow V, \quad (x_1, x_2, \dots, x_n) \mapsto (x_1, -x_2, \dots, -x_n). \quad (4.13)$$

□

Now let us restrict Y_1 further onto $M' = e'M \subset e''M$. From the above, Y_1 can be replaced with

$$(U + \sum_{i \neq 1}^n V_i \mathfrak{s}_{1i}) R(2\epsilon_1) (U^+ + \sum_{i \neq 1}^n V_i^+ \mathfrak{s}_{1i}^+) R(\delta + 2\epsilon_1) \quad (4.14)$$

Let us first work out the action of the product of the first three factors,

$$(U + \sum_{i \neq 1}^n V_i \mathfrak{s}_{1i}) (u_1 + v_1 \mathfrak{s}_1) (U^+ + \sum_{i \neq 1}^n V_i^+ \mathfrak{s}_{1i}^+). \quad (4.15)$$

Lemma 4.2. *The operator (4.15) preserves the subspace $M' = e'M$, and its restriction onto M' is given by $A + B\mathfrak{s}_1 + \sum_{i \neq 1} (C_i \mathfrak{s}_{1i} + D_i \mathfrak{s}_{1i}^+)$ with*

$$A = u_1 \prod_{l \neq 1}^n a_{1l} a_{1l}^+, \quad B = \tau^{2n-2} \tau_n - A - \sum_{i \neq 1} (C_i + D_i), \\ C_i = u_i b_{1i} a_{1i}^+ \prod_{l \neq 1, i} a_{il} a_{il}^+, \quad D_i = u_i^- b_{1i}^+ a_{1i} \prod_{l \neq 1, i} a_{li}^- a_{li}.$$

Proof. The operator (4.15) is composed of Y_1 and (the inverse of) $R(\delta + 2\epsilon_1)t(\epsilon_1)$. Both of these preserve M' , and so does (4.15). Next, for any element in W , its action on M' can be replaced with one of the elements $w = \text{id}, \mathfrak{s}_1, \mathfrak{s}_{1i}, \mathfrak{s}_{1i}^+$, since these elements represent the cosets in W/W' . Thus, restricting (4.15) to M' we obtain an expression of the form $A + B\mathfrak{s}_1 + \sum_{i \neq 1} (C_i \mathfrak{s}_{1i} + D_i \mathfrak{s}_{1i}^+)$. And since the operator (4.15) preserves M' , the resulting expression will be W' -invariant, in particular, $D_i = (C_i)^{\mathfrak{s}_i}$.

Let us now expand the product (4.15), moving the group elements to the right. We can pick either $w = \text{id}$ or \mathfrak{s}_{1i} from the first factor, $w = \text{id}$ or \mathfrak{s}_1 from the second, and $w = \text{id}$ or \mathfrak{s}_{1j}^+ from the third. It is easy to check that the product of three elements picked this way will represent the trivial coset $\text{id} \cdot W'$ if and only if $w = \text{id}$ is chosen from each factor. Therefore, $A = U u_1 U^+$. Also, the only way to obtain the representative \mathfrak{s}_{1i} is by picking $\mathfrak{s}_{1i}, \text{id}$, and id , respectively, from each

of the factors. Therefore, $C_i \mathfrak{s}_{1i} = V_i \mathfrak{s}_{1i} u_1 U^+$. This gives the required expressions for A and C_i , after which we find D_i as $D_i = (C_i)^{\mathfrak{s}_i}$.

To determine B , we use that $R_{1i} e = R_{1i}^+ e = \tau e$, $R(2\epsilon_1) e = \tau_n e$. As a result,

$$R_{12} R_{13} \dots R_{1n} R(2\epsilon_1) R_{1n}^+ \dots R_{12}^+ e = (\tau^{2n-2} \tau_n) e,$$

from which $A + B + \sum_{i \neq 1} (C_i + D_i) = \tau^{2n-2} \tau_n$. Lemma is proven. \square

4.3. Now we translate the obtained results into a matrix form. We choose $2n$ elements representing the (right and left) W' -cosets in W : $r_i = \mathfrak{s}_{1i}$, $r_{n+i} = \mathfrak{s}_{1i}^+$, where $1 \leq i \leq n$ and $r_1 = \mathfrak{s}_{11} := \text{id}$, $r_{n+1} = \mathfrak{s}_{11}^+ := \mathfrak{s}_1$. Suppose we have $Z = \sum_{i=1}^{2n} Z_i r_i$, with $Z_i \in \mathcal{D}_q(V)$. Assume that Z is W' -invariant, $wZ = Zw$ for all $w \in W'$, hence the multiplication by Z preserves $M' = e'M$. We write elements of M' as $\sum_{j=1}^{2n} e' r_j f_j$, with $f_j \in \mathbb{C}(x)$. Then we have

$$Z \left(e' \sum_{j=1}^{2n} r_j f_j \right) = e' Z \sum_{j=1}^{2n} r_j f_j = \sum_{i,j=1}^{2n} e' r_i r_j (Z_i)^{r_j r_i} f_j. \quad (4.16)$$

For each individual term in this sum we have: $e' r_i r_j (Z_i)^{r_j r_i} f_j = e' r_k (Z_i)^{r_j r_i} f_j$, where k is a uniquely defined index $k = k(i, j)$ such that $e' r_i r_j = e' r_k$. This means that this term represents the (k, j) -th entry of the matrix of Z . For the reader's convenience, here is an explicit description:

$$e' r_i r_j = e' r_k, \quad k = k(i, j) = \begin{cases} i & \text{for } i \notin \{1, n+1, j, j \pm n\} \\ j & \text{for } i = 1 \\ j+n & \text{for } i = n+1, 1 \leq j \leq n \\ j-n & \text{for } i = n+1, n+1 \leq j \leq 2n \\ 1 & \text{for } i = j \\ n+1 & \text{for } i = j \pm n. \end{cases}$$

Using this, we easily find all the entries of the matrix of Z . For example, for $1 < i \neq j \leq n$ we have $e' r_i r_j = e' r_i$, and so the (i, j) -th matrix entry is $(Z_i)^{r_j r_i}$.

Applying this procedure successively to $Z = A + B \mathfrak{s}_1 + \sum_{i \neq 1} (C_i \mathfrak{s}_{1i} + D_i \mathfrak{s}_{1i}^+)$ and to $Z = R(\delta + 2\epsilon_1) = \tilde{u}_1 t(\epsilon_1) + \tilde{v}_1 \mathfrak{s}_1$, we calculate the corresponding matrices. Let us state the result, leaving its (routine) verification to the reader.

Proposition 4.3. *Denote by \mathcal{P}, \mathcal{Q} the $2n \times 2n$ matrices representing the action on M' of $A + B \mathfrak{s}_1 + \sum_{i \neq 1} (C_i \mathfrak{s}_{1i} + D_i \mathfrak{s}_{1i}^+)$ and $\tilde{u}_1 t(\epsilon_1) + \tilde{v}_1 \mathfrak{s}_1$, respectively. Let us extend the set of vectors ϵ_i and variables x_i to the range $1 \leq i \leq 2n$ by setting $\epsilon_{i+n} = -\epsilon_i$ and $x_{n+i} = -x_i$ for $1 \leq i \leq n$. We also extend the definitions of a_{ij}, a_{ij}^+ , etc., accordingly. For instance, $a_{n+i, j}^+ = a(-x_i + x_j)$, $b_{i, n+j} = b(x_i + x_j)$, $u_{n+j} = u(-x_j)$ if $1 \leq i, j \leq n$. With this notation, we have:*

$$\begin{aligned} \mathcal{P}_{ij} &= u_j b_{ij} a_{ij}^+ \prod_{l=1}^{2n}{}' a_{jl} & (i-j \neq 0, \pm n), & \quad \mathcal{P}_{ii} &= u_i \prod_{l=1}^{2n}{}' a_{il}, \\ \mathcal{P}_{ij} &= \tau^{2n-2} \tau_n - \sum_{l \neq j}^{2n} \mathcal{P}_{il} & (i-j = \pm n), & \\ \mathcal{Q}_{ij} &= \tilde{u}_j t(\epsilon_j) & (i=j), & \quad \mathcal{Q}_{ij} = \tilde{v}_i & (i-j = \pm n), \quad \mathcal{Q}_{ij} = 0 & (i-j \neq 0, \pm n). \end{aligned}$$

Here the symbol \prod' in the formula for \mathcal{P}_{ij} indicates that we exclude those values of l where either $l-i$ or $l-j$ equals $0, \pm n$ (e.g., two values are excluded if $i=j$).

Corollary 4.4. *The quantum Lax matrix \mathcal{L} for the Koornwinder-van Diejen system is $\mathcal{L} = \mathcal{P} \mathcal{Q}$, with \mathcal{P}, \mathcal{Q} given above. The classical Lax matrix is found as $L = PQ$ where $P = \mathcal{P}$, while Q is obtained by setting $q = 1$ in the definitions of \tilde{u}, \tilde{v} and by replacing $t(\epsilon_j)$ with $e^{\beta p_j}$ (with $p_{n+i} = -p_i$ for $1 \leq i \leq n$).*

4.4. Finally, let us write down explicitly the analogue of Proposition 3.3 in the $C^\vee C_n$ case. We have

$$R_+ = \{2\epsilon_i \mid 1 \leq i \leq n\} \cup \{\epsilon_i \pm \epsilon_j \mid 1 \leq i < j \leq n\}, \quad R_+ \setminus R'_+ = \{2\epsilon_1\} \cup \{\epsilon_1 \pm \epsilon_i \mid 2 \leq i \leq n\}.$$

The functions \mathbf{c}_α are defined in (4.7)–(4.9). Substituting them into (3.25) gives

$$\phi = u_1^- \prod_{l \neq 1}^n a_{l1} a_{l1}^-, \quad \mathbf{u} = (\phi_1, \dots, \phi_{2n}), \quad \mathbf{v} = (1, \dots, 1)^T,$$

where

$$\phi_i = \phi^{s^{1i}} = u_i^- \prod_{l \neq i}^n a_{li} a_{li}^-, \quad \phi_{n+i} = \phi^{s^{1i}} = u_i \prod_{l \neq i}^n a_{li}^+ a_{li} \quad (1 \leq i \leq n).$$

Proposition 4.5. *The difference operators $\mathbf{u} \mathcal{L}^k \mathbf{v}$, $k \in \mathbb{Z}$ with the above \mathbf{u} , \mathbf{v} belong to the commutative algebra of the quantum integrals L_f , $f \in \mathbb{C}[Y_1^{\pm 1}, \dots, Y_n^{\pm 1}]^{\mathfrak{S}_n}$ of the Koornwinder–van Diejen system.*

The proof is identical to that of Proposition 3.3. By passing to the classical limit, the same result holds for the classical system. Of course, in the classical case one can also produce first integrals as $h_k = \text{tr}(L^k)$.

5. LAX PAIRS FOR THE ELLIPTIC CALOGERO–MOSER SYSTEMS

In this section we explain how our approach extends to the elliptic Calogero–Moser systems associated to any root system, including the BC_n case (Inozemtsev system). This will lead to Lax matrices with a spectral parameter, reproducing, in particular, the classical Lax pairs [Kr1, BMS, DHP]. For the Inozemtsev system we obtain a Lax pair of size $2n$, different from [I, DHP]. We will employ the theory of elliptic Dunkl operators, developed in [BFV, C5, EM, EFMV]. Note that in the elliptic case no quantum Lax pairs were known previously, although a quantum Lax operator in type A_n was considered in [Ha].

5.1. In the setting of Section §2.1, let W be a Weyl group with a root system $R \subset V = \mathbb{C}^n$ and a W -invariant set of parameters c_α , $\alpha \in R$. The Dunkl operators in the elliptic case depend on $t \neq 0$, the elliptic modulus τ , and additional *dynamical variables* represented by a vector $\lambda \in V$. They are as follows [BFV, EM, EFMV]:

$$y_\xi := t\partial_\xi + \sum_{\alpha \in R_+} c_\alpha \langle \alpha, \xi \rangle \sigma_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) s_\alpha, \quad \xi \in V. \quad (5.1)$$

Here $\alpha^\vee = 2\alpha/\langle \alpha, \alpha \rangle$, and

$$\sigma_\mu(z) = \frac{\theta_1(z - \mu)\theta_1'(0)}{\theta_1(z)\theta_1(-\mu)}, \quad (5.2)$$

where $\theta_1(z) = \theta_1(z|\tau)$ is the odd Jacobi theta function, associated with the elliptic curve $\mathbb{C}/\mathbb{Z} + \mathbb{Z}\tau$. Sometimes we will write $y_\xi(\lambda)$ to emphasize dependence on the dynamical variables. Note that as a function of λ , y_ξ has poles along the hyperplanes $\langle \alpha^\vee, \lambda \rangle = m + n\tau$ with $m, n \in \mathbb{Z}$. Below we will also need a classical limit of y_ξ , which is the following element of $\mathbb{C}(V)[p_1, \dots, p_n] * W$, cf. §2.3:

$$y_\xi^c := p_\xi + \sum_{\alpha \in R_+} c_\alpha \langle \alpha, \xi \rangle \sigma_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) s_\alpha, \quad \xi \in V. \quad (5.3)$$

Again, the two main properties of the Dunkl operators are their commutativity and equivariance [BFV]: for all $\xi, \eta \in V$ and $w \in W$,

$$y_\xi y_\eta = y_\eta y_\xi, \quad w y_\xi(\lambda) = y_{w\xi}(w\lambda) w. \quad (5.4)$$

Note that in the second relation the group action changes both ξ and λ . As before, the assignment $\xi \mapsto y_\xi$ extends to an algebra map

$$SV = \bigoplus_{i \geq 0} S^i V \rightarrow \mathcal{D}(V) * W, \quad q \mapsto q(y).$$

However, unlike in the rational case, this map is not W -equivariant. Despite that, there is a method for constructing commuting W -invariant quantum Hamiltonians from y_ξ , but it requires a certain regularization procedure [EFMV]. For the quadratic Hamiltonian this is straightforward

(this goes back to [BFV, Sec. 6]). Namely, let $\partial_i = \partial_{\xi_i}$ and $y_i = y_{\xi_i}$, where $\{\xi_i \mid i = 1 \dots n\}$ is an orthonormal basis in V . Let $\langle y, y \rangle = y_1^2 + \dots + y_n^2$. Then, by direct calculation,

$$\langle y, y \rangle = t^2 \Delta_V + t \sum_{\alpha \in R_+} c_\alpha \langle \alpha, \alpha \rangle \sigma'_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) s_\alpha + \sum_{\alpha \in R_+} c_\alpha^2 \langle \alpha, \alpha \rangle \sigma_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) \sigma_{\langle \alpha^\vee, \lambda \rangle}(-\langle \alpha, x \rangle).$$

Here $\sigma'_\mu(z) = \frac{d}{dz} \sigma_\mu(z)$. Using the identity $\sigma_\mu(z) \sigma_\mu(-z) = \wp(\mu) - \wp(z)$, we rewrite this as

$$\frac{1}{2} \langle y, y \rangle - \frac{1}{2} \sum_{\alpha \in R_+} c_\alpha^2 \langle \alpha, \alpha \rangle \wp(\langle \alpha^\vee, \lambda \rangle) = \widehat{H} + \widehat{A}, \quad (5.5)$$

where

$$\widehat{H} = \frac{1}{2} t^2 \Delta_V - \frac{1}{2} \sum_{\alpha \in R_+} c_\alpha (c_\alpha + t) \langle \alpha, \alpha \rangle \wp(\langle \alpha, x \rangle), \quad (5.6)$$

$$\widehat{A} = \frac{t}{2} \sum_{\alpha \in R_+} c_\alpha \langle \alpha, \alpha \rangle \left(\wp(\langle \alpha, x \rangle) + \sigma'_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) s_\alpha \right). \quad (5.7)$$

This is an elliptic analog of (2.11). Note that $\lim_{\mu \rightarrow 0} \sigma'_\mu(z) = -\wp(z) - 2\eta_1$, where $\eta_1 = \zeta(\frac{1}{2})$, see [BFV, Prop. 19(v)]. Thus, when λ approaches the hyperplane $\langle \alpha^\vee, \lambda \rangle = 0$, the term

$$c_\alpha \langle \alpha, \alpha \rangle \left(\wp(\langle \alpha, x \rangle) + \sigma'_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) s_\alpha \right)$$

in \widehat{A} can be replaced by its limit, i.e., by $c_\alpha \langle \alpha, \alpha \rangle (\wp(\langle \alpha, x \rangle)(1 - s_\alpha) - 2\eta_1 s_\alpha)$. This shows that the operator $\langle y, y \rangle$ after subtracting a λ -dependent term becomes regular in the neighbourhood of $\lambda = 0$.

Now, we obviously have $[y_\xi, \widehat{H} + \widehat{A}] = 0$, therefore we can construct a quantum Lax pair \mathcal{L}, \mathcal{A} of size $|W|$ in exactly the same way as in §2.2, but now the matrix \mathcal{A} will depend on the dynamical variables. Since \widehat{A} vanishes in the classical limit $t = -i\hbar \rightarrow 0$, the constructed Lax pair admits a classical limit.

To construct Lax pairs of smaller sizes, we use the same approach as in §2.4. We start by choosing ξ with nontrivial stabiliser W' . Denote by R' the root system of W' . Then $\langle \alpha, \xi \rangle = 0$ for $\alpha \in R'$, so

$$y_\xi = t \partial_\xi + \sum_{\alpha \in R_+ \setminus R'_+} c_\alpha \langle \alpha, \xi \rangle \sigma(\langle \alpha, x \rangle; \langle \alpha^\vee, \lambda \rangle) s_\alpha. \quad (5.8)$$

Choose λ (close to $\lambda = 0$) with the same stabiliser as ξ : this is possible because the singular terms with $\alpha \in R'$ are no longer present in (5.8). We also specialize λ in (5.5) (which is possible because the right-hand side is regular near $\lambda = 0$). It easily follows from (5.4) that under such a specialisation y_ξ and both sides of (5.5) become W' -invariant, that is, $w y_\xi = y_\xi w$ and $w(\widehat{H} + \widehat{A}) = (\widehat{H} + \widehat{A})w$ for all $w \in W'$. Since \widehat{H} is W -invariant, we obtain that $w \widehat{A} = \widehat{A}w$. As a result, both y_ξ and \widehat{A} preserve the subspace $M' = e' M$. Therefore, the same construction as in §2.4 applies, producing a Lax pair of size $|W|/|W'|$.

5.2. Let us illustrate the method in the case $W = \mathfrak{S}_n$. This is a modification of §2.6. We have n commuting Dunkl operators, depending on c, t and the dynamical parameters $\lambda = (\lambda_1, \dots, \lambda_n)$:

$$y_i = t \partial_i + c \sum_{j \neq i} \sigma_{\lambda_i - \lambda_j}(x_i - x_j) \mathfrak{s}_{ij}, \quad i = 1 \dots n.$$

Choose y_1 , so $\xi = (1, 0, \dots, 0)$ so W' and e' are the same as in §2.6, and

$$R'_+ = \{\epsilon_i - \epsilon_j \mid 2 \leq i < j \leq n\}.$$

Let us specialise λ in y_1 to $(\mu, 0, \dots, 0)$, where μ is an arbitrary parameter:

$$y_1 = t \partial_1 + c \sum_{j \neq 1} \sigma_\mu(x_1 - x_j) \mathfrak{s}_{1j}.$$

We also specialise λ in \widehat{A} , obtaining

$$\widehat{A} = ct \sum_{j \neq 1}^n (\wp(x_1 - x_j) + \sigma'_\mu(x_1 - x_j) \mathfrak{s}_{1j}) + ct \sum_{2 \leq i < j \leq n} (\wp(x_i - x_j)(1 - \mathfrak{s}_{ij}) - 2\eta_1 \mathfrak{s}_{ij}).$$

Since $\mathfrak{s}_{ij} = \text{id}$ on M' , the restriction of \widehat{A} onto M' can be replaced by

$$\widehat{A} = ct \sum_{j \neq 1}^n (\wp(x_1 - x_j) + \sigma'_\mu(x_1 - x_j) \mathfrak{s}_{1j}) - ct\eta_1(n-1)(n-2).$$

By removing an unimportant constant term, we may change \widehat{A} to

$$\widehat{A} = ct \sum_{j \neq 1}^n (\wp(x_1 - x_j) + \sigma'_\mu(x_1 - x_j) \mathfrak{s}_{1j}).$$

The quantum Lax pair \mathcal{L}, \mathcal{A} is now calculated in exactly the same way as in §2.6. This gives, after setting $t = -i\hbar$, $c = ig$, the following matrices:

$$\mathcal{L}_{kl} = \begin{cases} ig\sigma_\mu(x_k - x_l) & \text{for } k \neq l \\ \hat{p}_k & \text{for } k = l, \end{cases} \quad \mathcal{A}_{kl} = \begin{cases} g\hbar\sigma'_\mu(x_k - x_l) & \text{for } k \neq l \\ g\hbar \sum_{j \neq k}^n \wp(x_j - x_k) & \text{for } k = l. \end{cases} \quad (5.9)$$

This is an elliptic generalisation of (1.5)–(1.6) with

$$\widehat{H} = \frac{1}{2} \sum_{k=1}^n \hat{p}_k^2 + g(g - \hbar) \sum_{k < l}^n \wp(x_k - x_l).$$

In the classical limit $\hbar \rightarrow 0$ it gives the well-known Krichever's Lax pair with a spectral parameter [Kr1].

5.3. Let us describe the BC_n -case related to the Inozemtsev system [I]. This system depends on five coupling constants c, g_0, g_1, g_2, g_3 . We have n commuting Dunkl operators [EFMV]:

$$y_i = t\partial_i + v_{\lambda_i}(x_i)\mathfrak{s}_i + c \sum_{j \neq i} (\sigma_{\lambda_i - \lambda_j}(x_i - x_j)\mathfrak{s}_{ij} + \sigma_{\lambda_i + \lambda_j}(x_i + x_j)\mathfrak{s}_{ij}^+), \quad i = 1 \dots n.$$

Here $\lambda = (\lambda_1, \dots, \lambda_n)$ are the dynamical parameters, and

$$v_\mu(z) = v_\mu(z; g_0, g_1, g_2, g_3) = \sum_{r=0}^3 g_r \sigma_{2\mu}^r(z), \quad \sigma_\mu^r(z) := \frac{\theta_{r+1}(z - \mu)\theta_1'(0)}{\theta_{r+1}(z)\theta_1(-\mu)}, \quad (5.10)$$

where $\theta_r(z) = \theta_r(z|\tau)$, $r = 0 \dots 3$ are the Jacobi theta functions, with $\theta_4(z) := \theta_0(z)$. Clearly, $\sigma_\mu^0(z)$ coincides with (5.2). For an account of the properties of $\sigma_\mu^r(z)$, see [KH1, Appendix]. Note the following identities:

$$v_{-\mu}(-z) = -v_\mu(z), \quad v_\mu(z)v_\mu(-z) = \sum_{r=0}^3 ((g_r^\vee)^2 \wp(\mu + \omega_r) - g_r^2 \wp(z + \omega_r)), \quad (5.11)$$

where $(\omega_0, \omega_1, \omega_2, \omega_3) = (0, \frac{1}{2}, \frac{1+\tau}{2}, \frac{\tau}{2})$ are half-periods, and

$$\begin{pmatrix} g_0^\vee \\ g_1^\vee \\ g_2^\vee \\ g_3^\vee \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & -1 & -1 \\ 1 & -1 & 1 & -1 \\ 1 & -1 & -1 & 1 \end{pmatrix} \begin{pmatrix} g_0 \\ g_1 \\ g_2 \\ g_3 \end{pmatrix}. \quad (5.12)$$

Another property of $v_\mu(z)$ which will be needed later is the following symmetry between μ and z :

$$v_\mu(z; g_0, g_1, g_2, g_3) = v_{-z}(-\mu; g_0^\vee, g_1^\vee, g_2^\vee, g_3^\vee) = -v_z(\mu; g_0^\vee, g_1^\vee, g_2^\vee, g_3^\vee). \quad (5.13)$$

This can be checked by comparing transaltion properties and residues in the z -variable.

In the formulas below we will use the abbreviations $x_{ij} := x_i - x_j$, $x_{ij}^+ := x_i + x_j$, and similarly for the λ -variables. One calculates $\langle y, y \rangle = y_1^2 + \dots + y_n^2$ to get:

$$\begin{aligned} \langle y, y \rangle &= t^2 \Delta + 2c^2 \sum_{i < j}^n \left(\sigma_{\lambda_{ij}}(x_{ij}) \sigma_{\lambda_{ij}}(-x_{ij}) + \sigma_{\lambda_{ij}^+}(x_{ij}^+) \sigma_{\lambda_{ij}^+}(-x_{ij}^+) \right) \\ &\quad + 2ct \sum_{i < j}^n \left(\sigma_{\lambda_{ij}}(x_{ij}) \mathfrak{s}_{ij} + \sigma_{\lambda_{ij}^+}(x_{ij}^+) \mathfrak{s}_{ij}^+ \right) + \sum_{i=1}^n v_{\lambda_i}(x_i) v_{\lambda_i}(-x_i) + t \sum_{i=1}^n v'_{\lambda_i}(x_i) \mathfrak{s}_i. \end{aligned}$$

We can rewrite this as

$$\langle y, y \rangle - 2c^2 \sum_{i < j}^n (\wp(\lambda_{ij}) + \wp(\lambda_{ij}^+)) - \sum_{i=1}^n \sum_{r=0}^3 (g_r^\vee)^2 \wp(\lambda_i + \omega_r) = \widehat{H} + \widehat{A}, \quad (5.14)$$

where

$$\begin{aligned} \widehat{H} &= t^2 \Delta_V - 2c(c+t) \sum_{i < j}^n (\wp(x_{ij}) + \wp(x_{ij}^+)) - \sum_{i=1}^n \sum_{r=0}^3 g_r (g_r + t) \wp(x_i + \omega_r), \\ \widehat{A} &= 2ct \sum_{i < j}^n \left(\wp(x_{ij}) + \wp(x_{ij}^+) + \sigma'_{\lambda_{ij}}(x_{ij}) \mathfrak{s}_{ij} + \sigma'_{\lambda_{ij}^+}(x_{ij}^+) \mathfrak{s}_{ij}^+ \right) + t \sum_{i=1}^n \sum_{r=0}^3 g_r \wp(x_i + \omega_r) + t \sum_{i=1}^n v'_{\lambda_i}(x_i) \mathfrak{s}_i. \end{aligned}$$

We choose $\xi = (1, 0, \dots, 0)$ so W' and e' are the same as in §4.3, and the elements of $M' = e'M$ have the form $\sum_{i=1}^{2n} e' r_i f_i$, where $r_i = \mathfrak{s}_{1i}$, $r_{n+i} = \mathfrak{s}_{1i}^+$ ($1 \leq i \leq n$). The Lax matrix, therefore, will be of size $2n$. We will construct it from the operator y_1 , in which we set $\lambda = (\mu, 0, \dots, 0)$:

$$y_1 = t\partial_1 + v_\mu(x_1) \mathfrak{s}_1 + c \sum_{j \neq 1}^n (\sigma_\mu(x_{1j}) \mathfrak{s}_{1j} + \sigma_\mu(x_{1j}^+) \mathfrak{s}_{1j}^+).$$

Similarly, we specialise λ in \widehat{A} and obtain (after removing a constant):

$$\widehat{A}|_{M'} = 2ct \sum_{j \neq 1}^n (\wp(x_{1j}) + \wp(x_{1j}^+) + \sigma'_\mu(x_{1j}) \mathfrak{s}_{1j} + \sigma'_\mu(x_{1j}^+) \mathfrak{s}_{1j}^+) + t \sum_{r=0}^3 g_r \wp(x_1 + \omega_r) + t v'_\mu(x_1) \mathfrak{s}_1.$$

The quantum Lax pair \mathcal{L}, \mathcal{A} is now calculated in exactly the same way as in §4.3. To write down the answer in compact form, let us extend the range of the variables x_i to $1 \leq i \leq 2n$ by setting $x_{n+i} = -x_i$; we also set $\partial_{n+i} = -\partial_i$. Then we obtain the following $2n \times 2n$ matrices:

$$\mathcal{L}_{ij} = \begin{cases} c\sigma_\mu(x_i - x_j) & \text{for } i - j \neq 0, \pm n \\ v_\mu(x_i) & \text{for } i - j = \pm n \\ t\partial_i & \text{for } i = j, \end{cases}$$

$$\mathcal{A}_{ij} = \begin{cases} 2ct\sigma'_\mu(x_i - x_j) & \text{for } i - j \neq 0, \pm n \\ tv'_\mu(x_i) & \text{for } i - j = \pm n \\ 2ct \sum_{l: l-i \neq 0, \pm n}^{2n} \wp(x_i - x_l) + t \sum_{r=0}^3 g_r \wp(x_i + \omega_r) & \text{for } i = j. \end{cases}$$

The classical Lax pair is obtained in the limit $t = -i\hbar \rightarrow 0$, resulting in

$$L_{ij} = \begin{cases} c\sigma_\mu(x_i - x_j) & \text{for } i - j \neq 0, \pm n \\ v_\mu(x_i) & \text{for } i - j = \pm n \\ p_i & \text{for } i = j, \end{cases} \quad A = -t^{-1} \mathcal{A}. \quad (5.15)$$

Here we keep the same convention, $p_{n+i} = -p_i$. Note that the previously known Lax pairs for the classical Inozentsev system were of a larger size ($3n$ as in [I], or $2n + 1$ as in [DHP]). Probably, they can be brought to the above form by a suitable reduction.

5.4. According to [C5, EFMV], the elliptic Calogero–Moser Hamiltonian (5.6) is completely integrable: there is a commutative algebra of quantum Hamiltonians $L_q \in \mathcal{D}(V)$, $q \in (SV)^W$, each with the leading term $q(\partial)$. These Hamiltonians are W -invariant partial differential operators with *elliptic coefficients*, i.e. they are invariant under translations $t(v)$, $v \in P^\vee \oplus \tau P^\vee$. In [EFMV], a general procedure was given for constructing these higher Hamiltonians from the elliptic Dunkl operators. It consists of the following three steps, see [EFMV, Theorem 3.1]. First, one substitutes the elliptic Dunkl operators $y_\xi(\lambda)$ as momenta into suitable *classical* rational Calogero–Moser Hamiltonians, with the dynamical parameters λ as the position variables. Next, one goes to the limit $\lambda = 0$ (it is shown in [EFMV] that this limit exists), obtaining a reflection-difference operator. Finally, one restricts this reflection-differential operator to W -invariant functions (so it becomes a differential operator). There is also a parallel construction in the classical case [EFMV, Theorem 3.4].

For our purposes, we need a modification of that procedure, where the substitution is made into the classical *elliptic* Calogero–Moser Hamiltonians (cf. Remark 3.8 in [EFMV]). Namely, consider the following classical Hamiltonian $H^{\vee,c}(\xi, \lambda)$:

$$H^{\vee,c} = \frac{1}{2} \langle \xi, \xi \rangle - \frac{1}{2} \sum_{\alpha \in R_+} c_\alpha^2 \langle \alpha, \alpha \rangle \wp(\langle \alpha^\vee, \lambda \rangle). \quad (5.16)$$

By the above results of [EFMV], there is a family of Poisson-commuting Hamiltonians $L_q^{\vee,c}(\xi, \lambda)$, $q \in (SV)^W$, with the above $H^{\vee,c}$ corresponding to $q = \frac{1}{2} \langle \xi, \xi \rangle$. Note that each of these Hamiltonians is elliptic in λ with respect to the lattice $P \oplus \tau P$. Then we have the following result, whose proof was suggested to the author by P. Etingof.

Proposition 5.1. *Let $y_\xi(\lambda)$, $y_\xi^c(\lambda)$ denote the quantum and classical elliptic Dunkl operators associated to a root system R according to (5.1), (5.3). Let $L_q^{\vee,c}(\xi, \lambda)$, $q \in (SV)^W$ be the “dual” classical higher Hamiltonians associated with (5.16). Then:*

- (i) $L_q^{\vee,c}(y_\xi(\lambda), \lambda)$, viewed as an element of $\mathcal{D}(V) * W$ depending on λ , is regular for λ near $\lambda = 0$;
- (ii) $L_q^{\vee,c}(y_\xi^c(\lambda), \lambda)$, viewed as an element of $\mathbb{C}(V \times V) * W$ depending on λ , is regular for all $\lambda \in V$;
- (iii) $L_q^{\vee,c}(y_\xi^c(\lambda), \lambda)$ is constant in λ . Moreover, expanding $L_q^{\vee,c}(y_\xi^c(\lambda), \lambda)$ as $\sum_{w \in W} a_w w$ with $a_w \in \mathbb{C}(V \times V)$, we have $a_w = 0$ for $w \neq \text{id}$.

Proof. *Part (i).* In the case when $L_q^{\vee,c}$ is replaced by the Hamiltonian of the *rational* Calogero–Moser system, the regularity statement is Theorem 3.1 in [EFMV], and we can use the same method. More precisely, two different proofs of the regularity are given in [EFMV]. The first proof does not extend easily to the elliptic case because it requires establishing (iii) in advance, which we cannot do. However, the second proof as in [EFMV, 5.3] extends to the elliptic case verbatim.

Part (ii). The operators $y_\xi^c(\lambda)$ are the $\hbar = 0$ limit of $y_\xi(\lambda)$, thus, the expression $L_q^{\vee,c}(y_\xi^c(\lambda), \lambda)$ remains regular near $\lambda = 0$ by (i). Other possible singularities are along the hyperplanes $\langle \alpha^\vee, \lambda \rangle = m + n\tau$ with $m, n \in \mathbb{Z}$. To rule them out, let us see how $y_\xi^c(\lambda)$ changes under translations $\lambda \mapsto \lambda + u + \tau v$ with $u, v \in P$. Note that $\sigma_{\mu+1}(z) = \sigma_\mu(z)$ and $\sigma_{\mu+\tau}(z) = e^{2\pi i z} \sigma_\mu(z)$, by properties of $\theta_1(z|\tau)$. This gives

$$y_\xi^c(\lambda + u + \tau v) = p_\xi + \sum_{\alpha > 0} c_\alpha \langle \alpha, \xi \rangle e^{2\pi i \langle \alpha, x \rangle \langle \alpha^\vee, v \rangle} \sigma_{\langle \alpha^\vee, \lambda \rangle}(\langle \alpha, x \rangle) s_\alpha = e^{2\pi i \langle v, x \rangle} y_\xi^c(\lambda) e^{-2\pi i \langle v, x \rangle}.$$

Each Hamiltonian $L_q^{\vee,c}$ is elliptic in λ , so the expression $L_q^{\vee,c}(y_\xi^c(\lambda), \lambda)$ has the same translation properties:

$$L_q^{\vee,c}(y_\xi^c(\lambda), \lambda) \xrightarrow{\lambda \mapsto \lambda + u + \tau v} e^{2\pi i \langle v, x \rangle} L_q^{\vee,c}(y_\xi^c(\lambda), \lambda) e^{-2\pi i \langle v, x \rangle}. \quad (5.17)$$

Since we know that the right-hand side is regular along each of the hyperplane $\langle \alpha^\vee, \lambda \rangle = 0$, the left-hand side must be regular along the shifted hyperplanes. As a result, it is regular everywhere.

Part (iii). Let us expand $L_q^{\vee,c}(y_\xi^c(\lambda), \lambda)$ as $\sum_{w \in W} a_w w$. Each coefficient a_w is a function of x, p, λ , and by (ii) it is globally holomorphic in λ . From (5.17) we have that, as a function of λ , a_w is quasi-periodic with respect to the lattice $P \oplus \tau P$. However, a holomorphic quasi-periodic function must be a constant, which proves that each a_w is constant in λ . Invoking (5.17) once

again, we obtain $\widehat{a}_w w = e^{2\pi i \langle v, x \rangle} a_w w e^{-2\pi i \langle v, x \rangle}$ for all $v \in P$, from which it follows that $a_w = 0$ for $w \neq \text{id}$. \square

Remark 5.2. For $q = \frac{1}{2} \langle \xi, \xi \rangle$, all the statements of the proposition follow directly from the calculations in (5.5)–(5.7). Note that when the quantum Dunkl operators are substituted, \widehat{A} will contain terms of the form $\sigma'_\mu(z)$, which are not everywhere regular in μ . This shows that the global regularity property (ii) does not hold in the quantum setting (contrary to the claim made in [EFMV, Remark 3.8]).

Remark 5.3. For the BC_n -case with coupling constants c, g_0, g_1, g_2, g_3 , the proposition remains true, with the same proof as above. Note that the dual classical Hamiltonians in this case have coupling constants c and $g_0^\vee, g_1^\vee, g_2^\vee, g_3^\vee$ as defined in (5.12) (cf. Example 3.9 in [EFMV] and (5.14) above). A small modification is required for the proof of regularity at $\xi_i = \omega_r$ with $r = 1, 2, 3$. For this, one needs to employ shifts $\xi \mapsto \xi + \omega_r(\epsilon_1 + \dots + \epsilon_n)$. For such translations there is a formula similar to (5.17), but now it also involves a permutation of the parameters g_r . The rest of the proof does not change.

Remark 5.4. Lax pairs for the trigonometric Calogero–Moser system can be obtained by replacing $\sigma_\mu(z)$ and $\wp(z)$ by their trigonometric versions, $\sigma_\mu(z) = \cot z - \cot \mu$ and $\wp(z) = \sin^{-2} z$. In the BC_n case, the function $v_\mu(z)$ should also be replaced by its trigonometric version, $v_\mu(z) = g_0(\cot z - \cot 2\mu) + g_1(-\tan z - \cot 2\mu)$, cf. [FeP, Pu]. It is customary in the trigonometric case to set the spectral parameter to a specific value. For instance, a trigonometric version of the Lax pair (5.9) would have off-diagonal entries $\mathcal{L}_{kl} = \text{ig} \sigma_\mu(x_k - x_l) = \text{ig} \cot(x_k - x_l) - \cot \mu$, and setting $\mu = \pi/2$ would make it into $\mathcal{L}_{kl} = \text{ig} \cot(x_k - x_l)$.

Remark 5.5. By the same method, the proposition can be proved for a more general class of *crystallographic elliptic Calogero–Moser systems* constructed in [EFMV], for which the group W is not necessarily a Weyl group.

5.5. We can now construct a Lax partner \mathcal{A} for any of the higher quantum Hamiltonians. Namely, consider $L_q^{\vee, c}(y_\xi(\lambda), \lambda)$ with $q \in (SV)^W$. From the W -invariance of $L_q^{\vee, c}$ and W -equivariance of $y_\xi(\lambda)$, we have

$$w_x L_q^{\vee, c}(y_\xi(\lambda), \lambda) w_x^{-1} = L_q^{\vee, c}(y_{w\xi}(w\lambda), \lambda) = L_q^{\vee, c}(y_\xi(w\lambda), w^{-1}\lambda) \quad \text{for all } w \in W, \quad (5.18)$$

where we use the subscript to indicate that w_x acts in the x -variable. In the limit $\lambda = 0$ this becomes W -invariant, so we have

$$L_q^{\vee, c}(y_\xi(0), 0)e = \widehat{H}e, \quad \widehat{H} \in \mathcal{D}(V)^W,$$

where \widehat{H} is one of the commuting Hamiltonians of the Calogero–Moser system (5.6). By Proposition 5.1, the classical limit of \widehat{H} can be obtained as

$$H = L_q^{\vee, c}(y_\xi^c(\lambda), \lambda). \quad (5.19)$$

Now write

$$L_q^{\vee, c}(y_\xi(\lambda), \lambda) = \widehat{H} + \widehat{A}, \quad (5.20)$$

for a suitable $\widehat{A} \in \mathcal{D}(V) * W$ (depending on λ).

Obviously, $[y_\xi(\lambda), \widehat{H} + \widehat{A}] = 0$, so a quantum Lax pair of size $|W|$ can be constructed as before. Let us now specialise both ξ and λ to have the same stabiliser W' . In this case, $y_\xi(\lambda)$ will be W' -invariant by (5.4). Also, $L_q^{\vee, c}(y_\xi(\lambda), \lambda)$ and, therefore, \widehat{A} are W' -invariant by (5.18). As a result, $\mathcal{L}, \mathcal{H}, \mathcal{A}$ can be restricted onto $M' = e'M$, giving a Lax pair of size $|W/W'|$.

It remains to explain why the constructed Lax pairs have a classical limit. We know that the classical limit of \widehat{H} is the classical Hamiltonian H . Now, comparing (5.19) and (5.20), we see that the classical limit of \widehat{A} is zero. Therefore, the classical limit of $(i\hbar)^{-1}\widehat{A}$ is well-defined and this produces the classical Lax partner A in the same way as before.

5.6. Let us mention some consequences of the above for the classical systems.

Proposition 5.6. *Let $\xi = b_i$ be a fundamental (co)weight for a root system R , and W' denote the stabiliser of ξ in the Weyl group W of R . The classical elliptic Calogero–Moser system for a root system R admits a Lax matrix L of size $|W/W'|$ with a spectral parameter. Each of the commuting Hamiltonian flows of the system induces an isospectral deformation of L . The functions $h_k = \text{tr} L^k$, $k \in \mathbb{N}$, form an involutive family, that is, $\{\text{tr} L^a, \text{tr} L^b\} = 0$ for all a, b .*

Proof. The Lax matrix L is constructed from $y_\xi(\lambda)$ by taking $\xi = b_i$ and $\lambda = \mu b_i$, so it has μ as a spectral parameter. The isospectrality of L is a direct corollary of the existence of a Lax partner for each of the Hamiltonians. Let H_1, \dots, H_n be the full set of the commuting Hamiltonians. From the isospectrality of L we know that each $h_k = \text{tr} L^k$ remains constant under any of the commuting flows. Therefore, $\{h_k, H_i\} = 0$ for all $i = 1, \dots, n$. An abelian Poisson subalgebra in $\mathbb{C}(V \times V)$ cannot have more than n functionally independent elements, therefore, each h_k is a function of H_1, \dots, H_n . It follows that $\{h_k, h_l\} = 0$, as needed. \square

As a corollary, we can derive the following result.

Corollary 5.7. *Let L be the classical Lax matrix (5.15) for the Inosemtev system. Then the functions $h_k = \text{tr} L^{2k}$, $k = 1, \dots, n$ form a complete set of first integrals in involution.*

Proof. From the structure of L it is clear that each h_k is polynomial in momenta, with the leading terms $p_1^{2k} + \dots + p_n^{2k}$. Thus, h_1, \dots, h_n are functionally independent. Their involutivity was shown in the proposition above. \square

6. ELLIPTIC DIFFERENCE CASE

Here we extend our approach to the elliptic Ruijsenaars–Schneider system and its versions for other root systems. The corresponding generalisation of the Cherednik and Macdonald operators to the elliptic case was found by Komori and Hikami in [KH2], by developing the approach of [C6]. We will refer to these systems as *generalised Ruijsenaars systems*. The main tool is elliptic R -matrices, first introduced in type A in [SU]. We adjust some of the notation of [KH2] to make it closer to ours. The $C^\vee C_n$ case is related to the elliptic van Diejen system [vD1] and is treated separately in Subsections § 6.6–§ 6.10.

6.1. The setting is the same as in § 3.1: $R \subset V$ is a reduced, irreducible root system with Weyl group W , $R_a \subset \widehat{V} = V \oplus \mathbb{C}\delta$ is the associated affine root system, $W_a = W \ltimes t(Q^\vee)$ (resp. $\widehat{W} = W \ltimes t(P^\vee)$) is the affine (resp. extended affine) Weyl group. We choose a basis a_0, \dots, a_n of R_a , writing s_i for the corresponding simple reflections. Recall that $\widehat{W} = W_a \rtimes \Omega$ with $\Omega \cong P^\vee/Q^\vee$, and we have the length function $l(w)$ on \widehat{W} .

As in § 3.1, we choose $q = e^c$ and consider the algebra $\mathbb{C}(V) * W_a \cong \mathcal{D}_q * W$ of reflection-difference operators on V . Fix a set of W -invariant coupling constants m_α , $\alpha \in R$ (so in particular $m_{-\alpha} = m_\alpha$). For $\tilde{\alpha} = \alpha + k\delta \in R_a$, define R -matrices $R(\tilde{\alpha})$ to be the following elements of $\mathbb{C}(V) * W_a$:

$$R(\tilde{\alpha}) = \sigma_{m_\alpha}(\tilde{\alpha}) - \sigma_{\langle \alpha^\vee, \xi \rangle}(\tilde{\alpha}) s_{\tilde{\alpha}}, \quad (6.1)$$

where $\xi \in V$ are the dynamical variables, and $\sigma_\mu(z)$ is the function (5.2). According to [KH2, (4.5)], we have

$$R_\alpha R_{-\alpha} = \wp(m_\alpha) - \wp(\langle \alpha^\vee, \xi \rangle). \quad (6.2)$$

Definition 6.1. Define a set $\{R_w | w \in \widehat{W}\}$ by taking a reduced decomposition $w = s_{i_1} \dots s_{i_l} \pi$, $\pi \in \Omega$ and setting

$$R_w = R(\alpha^1) \dots R(\alpha^l), \quad \text{where } \alpha^1 = a_{i_1}, \alpha^2 = s_{i_1}(a_{i_2}), \dots, \alpha^l = s_{i_1} \dots s_{i_{l-1}}(a_{i_l}). \quad (6.3)$$

In particular, we have $R_{s_i} = R(a_i)$, $i = 0 \dots n$, and $R_\pi = 1$ for $\pi \in \Omega$. Elliptic Cherednik operators are defined as $Y^b = R_{t(b)} t(b)$, $b \in P^\vee$.

Theorem 6.2 ([KH2], Theorems 3.2 & 4.3). (i) *The elements R_w do not depend on the choice of a reduced decomposition for w ;*

(ii) *Y^b , $b \in P_+^\vee$ are pairwise commuting elements of $\mathcal{D}_q * W$.*

The proof is based on the fact that $R(\tilde{\alpha})$ satisfy the affine Yang–Baxter relations as in [KH2, (3.1a)–(3.1d)]; this idea goes back to Cherednik [C6]. The elliptic Macdonald–Ruijsenaars operators are obtained from the operators Y^b in the following way. Introduce the vector

$$\rho_m = \frac{1}{2} \sum_{\alpha \in R_+} m_\alpha \alpha. \quad (6.4)$$

It satisfies the following equations, cf. [KH2, (4.9)]:

$$\langle a_i^\vee, \rho_m \rangle = m_{\alpha_i}, \quad i = 1, \dots, n. \quad (6.5)$$

Theorem 6.3 ([KH2], Theorem 4.5). *Given $b \in P_+^\vee$, let $L^b \in \mathcal{D}_q$ be the unique difference operator such that $Y^b e = L^b e$. If $\xi = -\rho_m$ then each L^b is W -invariant, and the difference operators L^b , $b \in P_+^\vee$ form a commutative family.*

Remark 6.4. In [KH2] the operators Y^b with antidominant $b \in P_-^\vee$ are considered, but in fact they are the same as ours with $b \in P_+^\vee$. This is because the translations $t(v)$ are defined in [KH2, (2.10)] with the opposite sign compared to our conventions.

The operators L^b are complicated in general, but admit an explicit description when b is minuscule or quasi-minuscule.

Theorem 6.5 ([KH2], Theorems 6.1 & 6.5). *(i) Let b be a minuscule coweight, so that $\langle \alpha, b \rangle$ is either 0 or 1 for any $\alpha \in R_+$. Then we have*

$$L^b = \sum_{\pi \in Wb} A_\pi t(\pi), \quad A_\pi = \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle > 0}} \sigma_{m_\alpha}(\alpha). \quad (6.6)$$

(ii) Let b be a quasi-minuscule coweight of the form $b = \varphi^\vee$, with $\varphi \in R_+$ the highest root. In this case, $\langle \alpha, b \rangle \in \{0, 1\}$ for any $\alpha \in R_+ \setminus \{b\}$. Then

$$L^b = \sum_{\pi \in Wb} (A_\pi t(\pi) - B_\pi), \quad A_\pi = \sigma_{m_\varphi}(\pi^\vee + \delta) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle > 0}} \sigma_{m_\alpha}(\alpha), \quad (6.7)$$

$$B_\pi = \sigma_{\langle \varphi^\vee, -\rho_m \rangle}(\pi^\vee + \delta) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle > 0}} \sigma_{m_\alpha}(\alpha). \quad (6.8)$$

In these formulas the roots are viewed as affine-linear functions, so, for example, $\sigma_{m_\alpha}(\alpha + \delta) = \sigma_{m_\alpha}(\langle \alpha, x \rangle + c)$.

Remark 6.6. To connect the operators Y^b to the trigonometric Cherednik operators from §3, we rescale $x \mapsto x/2\pi i$ and let $\tau \rightarrow +i\infty$. In this limit (6.1) becomes

$$R(\tilde{\alpha}) = \frac{\sinh((\tilde{\alpha} - m_\alpha)/2)}{\sinh(\tilde{\alpha}/2) \sinh(-m_\alpha/2)} - \frac{\sinh(\tilde{\alpha} - \langle \alpha^\vee, \xi \rangle/2)}{\sinh(\tilde{\alpha}/2) \sinh(-\langle \alpha^\vee, \xi \rangle/2)} s_{\tilde{\alpha}}. \quad (6.9)$$

Let us set $\tau_{\tilde{\alpha}}^2 = e^{-m_\alpha}$, $\eta_{\tilde{\alpha}}^2 = e^{\langle \alpha^\vee, \xi \rangle}$ and rescale $R(\tilde{\alpha})$ by multiplying it by $\sinh(-m_\alpha/2)$. This gives

$$R(\tilde{\alpha}) = \frac{\tau_{\tilde{\alpha}}^{-1} - \tau_{\tilde{\alpha}} e^{\tilde{\alpha}}}{1 - e^{\tilde{\alpha}}} + \frac{(\tau_{\tilde{\alpha}} - \tau_{\tilde{\alpha}}^{-1})(\eta_{\tilde{\alpha}} - \eta_{\tilde{\alpha}}^{-1} e^{\tilde{\alpha}})}{(\eta_{\tilde{\alpha}} - \eta_{\tilde{\alpha}}^{-1})(1 - e^{\tilde{\alpha}})} s_{\tilde{\alpha}}.$$

If we take $\xi \rightarrow \infty$ deep inside the positive Weyl cone, then $\eta_\alpha \rightarrow \infty$ for $\alpha > 0$, which turns the above formula into (3.14) (assuming $\alpha > 0$). It is known that for $w = t(b)$ with $b \in P_+^\vee$ all the R -matrices that appear in the decomposition (6.3) will be of the form $R(\alpha + k\delta)$ with $\alpha \in R_+$, $k \geq 0$. Thus, the elliptic operators Y^b , $b \in P_+^\vee$ in the trigonometric limit coincide with the operators from §3, up to an overall factor.

Remark 6.7. In the GL_n case, the operators L^b from Theorem 6.5 are equivalent to the commuting Hamiltonians of the elliptic Ruijsenaars system [R1]. For other root systems, the explicit operators L^b can be viewed as elliptic generalisations of the Macdonald difference operators [M1].

6.2. For later purposes we also need a “dual” version of the Cherednik operators for the affine root system $(R_a)^\vee$. To any affine root $\tilde{\alpha} = \alpha + k\delta$ we associate a coroot $\tilde{\alpha}^\vee$ by the formula

$$(\alpha + k\delta)^\vee = \frac{2}{\langle \alpha, \alpha \rangle} (\alpha + k\delta).$$

The coroots $\tilde{\alpha}^\vee$ with $\tilde{\alpha} \in R_a$ form a dual affine root system $(R_a)^\vee \subset \widehat{V}$. The hyperplanes $\tilde{\alpha} = 0$ and $\tilde{\alpha}^\vee = 0$ are the same, so $s_{\tilde{\alpha}} = s_{\tilde{\alpha}^\vee}$ and both R_a and $(R_a)^\vee$ share the same affine Weyl group W_a . The group $\widehat{W} = W \rtimes P^\vee = W_a \rtimes \Omega$ permutes both coroots and roots, so we view it as an extended affine Weyl group for both systems. We can take $a_0^\vee, \dots, a_n^\vee$ as a basis of $(R_a)^\vee$; note that Ω acts on this basis by permutations. For $\tilde{\alpha} = \alpha + k\delta \in R_a$ we define

$$R(\tilde{\alpha}^\vee) = \sigma_{m_\alpha}(\tilde{\alpha}^\vee) - \sigma_{\langle \alpha, \zeta \rangle}(\tilde{\alpha}^\vee) s_{\tilde{\alpha}}, \quad (6.10)$$

where the dynamical variable is $\zeta \in V$. We also define a set of elements $\{R_w^\vee | w \in \widehat{W}\}$ by “dualising” (6.3):

$$R_w^\vee = R(\alpha^1) \dots R(\alpha^l), \quad \text{where } \alpha^1 = a_{i_1}^\vee, \alpha^2 = s_{i_1}(a_{i_2}^\vee), \dots, \alpha^l = s_{i_1} \dots s_{i_{l-1}}(a_{i_l}^\vee).$$

In particular, $R_{s_i}^\vee = R(a_i^\vee)$, $i = 0 \dots n$, and $R_\pi^\vee = 1$ for $\pi \in \Omega$. Then the same arguments as in [KH2] (see also Proposition 6.17 below) prove the following results.

Theorem 6.8. *The elements R_w^\vee do not depend on the choice of a reduced decomposition for $w \in \widehat{W}$. The elements $Y^{b,\vee} := R_{t(b)}^\vee t(b)$, $b \in P_+^\vee$ pairwise commute.*

Theorem 6.9. *Assume $\zeta = -\rho_m^\vee$ with $\rho_m^\vee = \frac{1}{2} \sum_{\alpha \in R_+} m_\alpha \alpha^\vee$. Given $b \in P_+^\vee$, define $L^{b,\vee} \in \mathcal{D}_q$ as the unique difference operator such that $Y^{b,\vee} e = L^{b,\vee} e$. Then each $L^{b,\vee}$ is W -invariant, and $L^{b,\vee}$ with $b \in P_+^\vee$ form a commutative family of difference operators.*

Theorem 6.10. (i) *Let b be a minuscule coweight. Then we have*

$$L^{b,\vee} = \sum_{\pi \in Wb} A_\pi^\vee t(\pi), \quad A_\pi^\vee = \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle > 0}} \sigma_{m_\alpha}(\alpha^\vee). \quad (6.11)$$

(ii) *Let b be a quasi-minuscule coweight of the form $b = \varphi^\vee$, with $\varphi \in R_+$ the highest root. Then*

$$L^{b,\vee} = \sum_{\pi \in Wb} (A_\pi^\vee t(\pi) - B_\pi^\vee), \quad A_\pi^\vee = \sigma_{m_\varphi}(\pi + 2\delta / \langle \pi, \pi \rangle) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle > 0}} \sigma_{m_\alpha}(\alpha^\vee), \quad (6.12)$$

$$B_\pi^\vee = \sigma_{\langle \varphi, -\rho_m^\vee \rangle}(\pi + 2\delta / \langle \pi, \pi \rangle) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle > 0}} \sigma_{m_\alpha}(\alpha^\vee). \quad (6.13)$$

6.3. Before proceeding to the general construction of quantum Lax pairs, it will be instructive to discuss the \mathfrak{GL}_n -case. Our setting will be similar to §3.9: we take $V = \mathbb{C}^n$, with an orthonormal basis $\epsilon_1, \dots, \epsilon_n$ and the associated coordinates x_1, \dots, x_n , and with the standard action of $W = \mathfrak{S}_n$ on V . We set $\Lambda = \sum_{i=1}^n \mathbb{Z}\epsilon_i$, and consider the algebra of difference operators $\mathcal{D}_q = \mathbb{C}(x) \rtimes t(\Lambda)$, generated by $\mathbb{C}(x)$ and the shift operators $t(\epsilon_i) = q^{\partial_i}$, where $q = e^c$ with fixed c . As before, we will view reflection-difference operators in $\mathcal{D}_q * W$ acting on the module M (3.21), and identify $M \cong \mathbb{C}W \otimes \mathbb{C}(x)$. As a result, we represent elements of $\mathcal{D}_q * W$ as operator-valued matrices of size $|W|$ (3.22).

We have one coupling constant $m_\alpha = \mu$ for all $\alpha \in R$. For $\alpha = \epsilon_i - \epsilon_j$, $i \neq j$, the R -matrices (3.14) take the form [SU]:

$$R_{ij} = \sigma_\mu(x_i - x_j) - \sigma_{\xi_i - \xi_j}(x_i - x_j) \mathfrak{s}_{ij}. \quad (6.14)$$

They satisfy the Yang–Baxter relations, $R_{ij} R_{ik} R_{jk} = R_{jk} R_{ik} R_{ij}$. The elliptic Cherednik operators $Y_i := Y^{\epsilon_i}$ can be calculated from Definition 6.1 (cf. (6.2)):

$$Y_i = R_{i,i+1} R_{i,i+2} \dots R_{i,n} t(\epsilon_i) R_{i1} \dots R_{i,i-1} \quad (i = 1, \dots, n). \quad (6.15)$$

The elliptic Ruijsenaars operator $\widehat{H} = L^{\epsilon_1}$ is obtained from $Y_1 = Y^{\epsilon_1}$ by Theorems 6.3, 6.5:

$$\widehat{H} = \sum_{i=1}^n \prod_{j \neq i} \sigma_\mu(x_i - x_j) t(\epsilon_i). \quad (6.16)$$

Up to a gauge transformation, this is the quantum Hamiltonian from [R1].

A quantum Lax pair will be constructed using Y_1 and Y_2 , which are:

$$Y_1 = R_{12} \dots R_{1n} t(\epsilon_1), \quad Y_2 = R_{23} \dots R_{2n} t(\epsilon_2) R_{21}. \quad (6.17)$$

Lemma 6.11. *Write $W' = \mathfrak{S}_{n-1}$ for the subgroup fixing ϵ_1 , and e' for the corresponding symmetrizer. If $\xi_i - \xi_{i+1} = -\mu$ for all $1 < i < n$, then Y_1 and Y_2 preserve the subspace $M' = e'M$.*

Proof. This can be proved similarly to [KH2, Theorem 4.5], but for the reader's convenience we give a self-contained proof. From the assumptions on ξ , for any $1 < i < n$ we have

$$R_{i,i+1} = (1 + \mathfrak{s}_{i,i+1})\sigma_\mu(x_i - x_{i+1}), \quad R_{i+1,i} = \sigma_\mu(x_{i+1} - x_i)(1 - \mathfrak{s}_{i,i+1}), \quad R_{i+1,i}R_{i,i+1} = 0.$$

Next, the Yang–Baxter relations imply that

$$R_{i+1,i}Y_1 = R_{12} \dots R_{1,i-1}R_{1,i+1}R_{1,i}R_{1,i+2} \dots R_{1n} t(\epsilon_1)R_{i+1,i}.$$

Assuming $i > 1$, and multiplying by e' from the right, we obtain

$$(1 - \mathfrak{s}_{i,i+1})Y_1e' = 0 \quad \text{for } i > 1, \quad (6.18)$$

so Y_1 acts on $M' = e'M$ as a consequence. For Y_2 , we first notice that $R_{32}Y_2 = 0$ since $R_{32}R_{23} = 0$, and so $(1 - \mathfrak{s}_{23})Y_2 = 0$ as a consequence. Also, we have $(1 - \mathfrak{s}_{i,i+1})Y_2e' = 0$ for $i > 2$; this follows from the Yang–Baxter relations in the same way as (6.18). Putting this together, we conclude that $(1 - w)Y_2e' = 0$ for all $w \in W'$, hence $Y_2(M') \subset M'$. \square

Lemma 6.12. *Assume that $\xi_i - \xi_{i+1} = -\mu$ for all $1 < i < n$. We have:*

$$Y_1|_{M'} = (A + \sum_{i \neq 1}^n B_i \mathfrak{s}_{1i}) t(\epsilon_1), \quad Y_2|_{M'} = E + \sum_{i \neq 1}^n F_i \mathfrak{s}_{1i}, \quad \text{where}$$

$$A = \prod_{l \neq 1}^n \sigma_\mu(x_1 - x_l), \quad B_i = -\sigma_{\xi_1 - \xi_2}(x_1 - x_i) \prod_{l \neq 1, i} \sigma_\mu(x_i - x_l),$$

$$E = \sum_{i > 1} \sigma_\mu(x_i - x_1 + c) \prod_{l \neq 1, i} \sigma_\mu(x_i - x_l) t(\epsilon_i), \quad F_i = \sigma_{\xi_1 - \xi_2}(x_1 - x_i - c) \prod_{l \neq 1, i} \sigma_\mu(x_i - x_l) t(\epsilon_i).$$

Proof. The formula for Y_1 is proved in the same way as in Lemma §3.10. For Y_2 we can argue similarly. Namely, after expanding Y_2 (6.17), we obtain a sum of terms of the form

$$\mathfrak{s}_{2i_1} \dots \mathfrak{s}_{2i_k} t(\epsilon_2) w_0, \quad \text{with } 3 \leq i_1 < \dots < i_k \leq n \text{ and } w_0 \in \{\text{id}, \mathfrak{s}_{12}\}.$$

Multiplication by e' reduces this to $\mathfrak{s}_{1i}e'$ with $1 \leq i \leq n$, which equals $\mathfrak{s}_{12}e'$ only if the set $\{i_1, \dots, i_k\}$ is empty and $w_0 = \mathfrak{s}_{12}$. Therefore, the coefficient F_2 is found from

$$F_2 \mathfrak{s}_{12} = \prod_{l \neq 1, 2} \sigma_\mu(x_2 - x_l) t(\epsilon_2) \sigma_{\xi_1 - \xi_2}(x_1 - x_2) \mathfrak{s}_{12}.$$

This gives F_2 as stated in the lemma. The other coefficients F_i are determined by the symmetry, $F_i = (F_2)^{\mathfrak{s}_{2i}}$. It remains to determine the coefficient E . Arguing as above, we have several terms that reduce to ge' with $g \in \mathcal{D}_q$, but only one of them will contain $t(\epsilon_2)$. Namely, this happens if $\{i_1, \dots, i_k\} = \emptyset$ and $w_0 = \text{id}$, so the corresponding term is

$$\prod_{l \neq 1, 2} \sigma_\mu(x_2 - x_l) t(\epsilon_2) \sigma_\mu(x_2 - x_1).$$

If we combine this with the fact that E is W' -invariant, we will arrive at the expression given in the lemma. \square

To construct a Lax pair, we set $\widehat{A} = Y_1 + Y_2 - \widehat{H}$, where \widehat{H} is the Hamiltonian (6.16). Then from the above lemma, we have:

$$\widehat{A}|_{M'} = \sum_{i=1}^n Z_i \mathfrak{s}_{1i}, \quad \text{with } Z_1 = \sum_{i \neq 1} (\sigma_\mu(x_{i1} + c) - \sigma_\mu(x_{i1})) \prod_{l \neq 1, i} \sigma_\mu(x_{il}) t(\epsilon_i), \quad (6.19)$$

$$Z_i = (\sigma_\eta(x_{1i} - c) - \sigma_\eta(x_{1i})) \prod_{l \neq 1, i} \sigma_\mu(x_{il}) t(\epsilon_i) \quad (i \neq 1). \quad (6.20)$$

Here $\eta := \xi_1 - \xi_2$ and $x_{il} := x_i - x_l$.

Using Lemma 2.3, we can now calculate the matrices \mathcal{L} , \mathcal{A} of size n that represent the action of Y_1 and \widehat{A} on M' . The result is:

$$\mathcal{L}_{ij} = \begin{cases} \left(\prod_{l \neq j} \sigma_\mu(x_{jl}) \right) t(\epsilon_j) & \text{for } i = j \\ -\sigma_\eta(x_{ij}) \left(\prod_{l \neq i, j} \sigma_\mu(x_{jl}) \right) t(\epsilon_j) & \text{for } i \neq j, \end{cases}$$

$$\mathcal{A}_{ij} = \begin{cases} \sum_{k \neq j} (\sigma_\mu(x_{kj} + c) - \sigma_\mu(x_{kj})) \left(\prod_{l \neq j, k} \sigma_\mu(x_{kl}) \right) t(\epsilon_k) & \text{for } i = j \\ (\sigma_\eta(x_{ij} - c) - \sigma_\eta(x_{ij})) \left(\prod_{l \neq i, j} \sigma_\mu(x_{jl}) \right) t(\epsilon_j) & \text{for } i \neq j. \end{cases}$$

Since $[Y_1, Y_1 + Y_2] = 0$, the above \mathcal{L} , \mathcal{A} and $\mathcal{H} = \widehat{H}\mathbb{1}$ satisfy the quantum Lax equation (1.6). Note that \mathcal{L} is equivalent to the Lax matrix found by Hasegawa using a different approach, cf. [Ha, eq. (38)].

The classical limit corresponds to $c = -i\hbar\beta \rightarrow 0$. In \mathcal{L} we simply replace $t(\epsilon_j) = e^{c\partial_j}$ with $e^{\beta p_j}$:

$$L_{ij} = \begin{cases} \left(\prod_{l \neq j} \sigma_\mu(x_{jl}) \right) e^{\beta p_j} & \text{for } i = j \\ -\sigma_\eta(x_{ij}) \left(\prod_{l \neq i, j} \sigma_\mu(x_{jl}) \right) e^{\beta p_j} & \text{for } i \neq j. \end{cases}$$

The Lax partner A is then found as the classical limit of $(i\hbar)^{-1}\mathcal{A}$, which gives

$$A_{ij} = \begin{cases} -\sum_{k \neq j} \beta \sigma'_\mu(x_{kj}) \left(\prod_{l \neq j, k} \sigma_\mu(x_{kl}) \right) e^{\beta p_k} & \text{for } i = j \\ \beta \sigma'_\eta(x_{ij}) \left(\prod_{l \neq i, j} \sigma_\mu(x_{jl}) \right) e^{\beta p_j} & \text{for } i \neq j. \end{cases}$$

These matrices are equivalent to those known from [R1, BCa, KrZ].

6.4. Let us now discuss a method for constructing Lax pairs in general. As before, we will view reflection-difference operators in $\mathcal{D}_q * W$ acting on the module M (3.21), and identify $M \cong \mathbb{C}W \otimes \mathbb{C}(V)$. As a result, we represent elements of $\mathcal{D}_q * W$ as operator-valued matrices of size $|W|$ (3.22). Take an elliptic Cherednik operator Y^b , $b \in P_+^\vee$ and the corresponding quantum Hamiltonians $L^b \in \mathcal{D}_q^W$, constructed in Theorem 6.3. Our task is to find a combination of Cherednik operators which has the same classical limit as L^b . In the \mathbf{GL}_n -case above such a combination was $Y_1 + Y_2$, but this does not seem very helpful in regards to the general case. There is, nevertheless, a natural analogue of Proposition 5.1, but it requires a renormalisation of the operators $R(\alpha)$ and Y^b . Namely, we define the *unitary* affine R -matrices as follows:

$$\widehat{R}(\tilde{\alpha}) = \sigma_{m_\alpha}(\langle \alpha^\vee, \xi \rangle)^{-1} R(\tilde{\alpha}) = \frac{\sigma_{m_\alpha}(\tilde{\alpha})}{\sigma_{m_\alpha}(\langle \alpha^\vee, \xi \rangle)} - \frac{\sigma_{\langle \alpha^\vee, \xi \rangle}(\tilde{\alpha})}{\sigma_{m_\alpha}(\langle \alpha^\vee, \xi \rangle)} s_{\tilde{\alpha}}. \quad (6.21)$$

This can be rewritten as

$$\widehat{R}(\tilde{\alpha}) = \frac{\theta(\tilde{\alpha} - m_\alpha)\theta(\langle \alpha^\vee, \xi \rangle)}{\theta(\tilde{\alpha})\theta(\langle \alpha^\vee, \xi \rangle - m_\alpha)} - \frac{\theta(\tilde{\alpha} - \langle \alpha^\vee, \xi \rangle)\theta(-m_\alpha)}{\theta(\tilde{\alpha})\theta(\langle \alpha^\vee, \xi \rangle - m_\alpha)} s_{\tilde{\alpha}}. \quad (6.22)$$

From (6.2) it follows that

$$\widehat{R}(\tilde{\alpha})\widehat{R}(-\tilde{\alpha}) = 1 \quad \text{for all } \tilde{\alpha} \in R_a. \quad (6.23)$$

We also define \widehat{R}_w , $w \in W_a$ in the same way as in Definition 6.1, but using the unitary R -matrices instead, and we set $\widehat{Y}^b = \widehat{R}_{t(b)} t(b)$ for $b \in P_+^\vee$. The elements $\widehat{R}_w, \widehat{Y}^b$ differ from R_w, Y^b by a ξ -depending factor, and Theorem 6.2 remains valid for them. The classical Cherednik operators are $Y_c^\lambda := \eta_0(Y^\lambda)$ and $\widehat{Y}_c^\lambda := \eta_0(\widehat{Y}^\lambda)$, where the classical-limit map η_0 is defined in § 3.6. We also write $L_c^{b, \vee}$ for the classical limit of the Hamiltonians $L^{b, \vee}$ from Theorems 6.9, 6.10; these are elements of $A_0 = \mathbb{C}(V)[P^\vee]$, i.e. linear combinations of the terms $g(x)e^{\beta p_\lambda}$ with $\lambda \in P^\vee$, see § 3.6 for the notation. Denote by $L_c^{b, \vee}(\xi, \widehat{Y})$ the result of replacing each term $g(x)e^{\beta p_\lambda}$ by $g(\xi)\widehat{Y}^\lambda$. Similarly, we write $L_c^{b, \vee}(\xi, \widehat{Y}_c)$ for the result of substituting the classical operators \widehat{Y}_c^λ . We then have the following analogue of Proposition 5.1.

Proposition 6.13. *Let $L_c^{b, \vee}$, $b \in P_+^\vee$ be one of the dual classical Hamiltonians. Then, assuming that b is (quasi-)minuscule, we have:*

- (i) $L_c^{b, \vee}(\xi, \widehat{Y})$, viewed as an element of $\mathcal{D}(V) * W$ depending on ξ , is regular for ξ near $\xi = 0$;
- (ii) $L_c^{b, \vee}(\xi, \widehat{Y}_c)$, viewed as an element of $\mathbb{C}(V \times V) * W$ depending on ξ , is regular for all $\xi \in V$;

(iii) $L_c^{b,\vee}(\xi, \widehat{Y}_c)$ is constant in ξ . Moreover, expanding $L_c^{b,\vee}(\xi, \widehat{Y}_c)$ as $\sum_{w \in W} a_w w$ with $a_w \in \mathbb{C}(V \times V)$, we have $a_w = 0$ for $w \neq \text{id}$.

(iv) We have $L_c^{b,\vee}(\xi, \widehat{Y}_c) = L_c^b + \text{const}$.

We will prove the proposition in §6.5. We expect that a similar result is true for any of the classical Hamiltonians $L_c^{b,\vee}$, not only for those associated with (quasi-)minuscule coweights.

A quantum Lax pair can now be constructed as in (5.19)–(5.20), by writing

$$L_c^{b,\vee}(\xi, \widehat{Y}) = L^b + \widehat{A}, \quad (6.24)$$

for a suitable $\widehat{A} \in \mathcal{D}_q * W$. Pick a Cherednik operator Y^λ , $\lambda \in P_+^\vee$. Then Y^λ and $L^b + \widehat{A}$ commute, and L^b is W -invariant, which leads to a quantum Lax pair of size $|W|$. Since $L_c^{b,\vee}(\xi, \widehat{Y}_c) = L_c^b + \text{const}$ by Proposition 6.13(iv), the classical limit of \widehat{A} is a constant. We can modify \widehat{A} by subtracting this constant; as a result, the constructed Lax pair will admit a classical limit. Finally, to reduce it to a Lax pair of a smaller size, we use the following lemma whose proof is postponed to §6.5.

Lemma 6.14. *Let $\lambda = b_k$, $1 \leq k \leq n$ be one of the fundamental coweights, W' be its stabiliser and e' be the corresponding symmetrizer. Furthermore, assume that $\xi = -\rho_m + \eta b_k$, with arbitrary η . Then for such λ, ξ the action of Y^λ and $L_c^{b,\vee}(\xi, \widehat{Y})$ on the module M preserve the subspace $M' = e'M$.*

This means that the quantum Lax pair constructed from (6.24) can be restricted to M' , giving matrices of size $|W|/|W'|$. Note that the Lax pair depends on a spectral parameter η . The following theorem summarizes the obtained results.

Theorem 6.15. *For any fundamental coweight $\lambda = b_k \in P_+^\vee$, with the stabiliser W' , and any (quasi-)minuscule coweight b , there exists a quantum Lax pair \mathcal{L}, \mathcal{A} of size $|W|/|W'|$ satisfying the Lax equation (1.6) with $\mathcal{H} = \widehat{H}\mathbb{1}$, $\widehat{H} = L^b$. This Lax pair depends on a spectral parameter and admits a classical limit.*

6.5. In this subsection we prove Proposition 6.13 and Lemma 6.14. For this we need to analyse the operators \widehat{R}_w and \widehat{Y}^b in more detail. First, let us introduce some notation. The R -matrices $R(\tilde{\alpha})$, $\widehat{R}(\tilde{\alpha})$ depend on $\xi \in V$, so it will be convenient to introduce $\mathbb{V} = V^\vee \times V$ and $\mathbb{C}(\mathbb{V})$, where the first factor $V^\vee \cong V$ represents the ξ -variable. We also introduce $\mathbb{W} = W \times \widehat{W}$, to allow W acting on ξ ; the subgroup $\text{id} \times \widehat{W}$ will be identified with \widehat{W} . This makes $\mathbb{C}(\mathbb{V})$ into a $\mathbb{C}\mathbb{W}$ -module, so we form the product $\mathbb{C}(\mathbb{V}) \rtimes \mathbb{C}\mathbb{W}$. With each $w \in \widehat{W}$ we associate an element $w^\vee \otimes w \in \mathbb{W}$, where $w^\vee \in W$ is the linear part of w ; in particular, for $w = t(b)$ we have $w^\vee = 1$, so $w^\vee \otimes w = 1 \otimes t(b)$.

With this notation, the R -matrices can be viewed as elements of $\mathbb{C}(\mathbb{V}) \rtimes \mathbb{C}\mathbb{W}$. An important property is their equivariance in the following sense:

$$(w^\vee \otimes w)R(\tilde{\alpha})(w^\vee \otimes w)^{-1} = R(w\tilde{\alpha}), \quad \tilde{\alpha} \in R_a, \quad w \in \widehat{W}, \quad (6.25)$$

and the same for $\widehat{R}(\tilde{\alpha})$. In particular, for $w = t(b)$, $b \in P^\vee$ this implies (cf. [KH2, (4.6)])

$$t(b)R(\alpha + k\delta)t(-b) = R(\alpha + k'\delta), \quad k' = k + \langle \alpha, b \rangle. \quad (6.26)$$

Another crucial property is that $R(\tilde{\alpha})$ and $\widehat{R}(\tilde{\alpha})$ satisfy the affine Yang–Baxter relations [KH2, (3.1)(a)–(d)]. This property can be reformulated by setting

$$\mathbb{T}_i = \widehat{R}(a_i)(s_i^\vee \otimes s_i), \quad \mathbb{T}_\pi = \pi^\vee \otimes \pi \quad (6.27)$$

for $i = 0, \dots, n$ and $\pi \in \Omega$. Then one can check that the relations [KH2, (3.1)(a)–(d)] imply that \mathbb{T}_i and \mathbb{T}_π satisfy the relations of the braid group (3.7)–(3.9). Moreover, the unitarity (6.23) implies that $\mathbb{T}_i^2 = 1$ for all i . Therefore, we have the following result.

Proposition 6.16. *The assignment $s_i \mapsto \mathbb{T}_i$ ($i = 0, \dots, n$), $\pi \mapsto \mathbb{T}_\pi$ ($\pi \in \Omega$) extends to a group homomorphism $w \mapsto \mathbb{T}_w$, $w \in \widehat{W}$. In particular, the elements $\mathbb{T}_{t(b)}$, $b \in P^\vee$ pairwise commute.*

The elements \mathbb{T}_w can be rewritten in terms of the affine R -matrices, giving the following result.

Proposition 6.17. *For any reduced decomposition of w into $w = s_{i_1} \dots s_{i_l} \pi$ with $\pi \in \Omega$ we have $\mathbb{T}_w = \widehat{R}_w(w^\vee \otimes w)$, where*

$$\widehat{R}_w = \widehat{R}(\alpha^1) \dots \widehat{R}(\alpha^l), \quad \alpha^1 = a_{i_1}, \quad \alpha^2 = s_{i_1}(a_{i_2}), \quad \dots, \quad \alpha^l = s_{i_1} \dots s_{i_{l-1}}(a_{i_l}). \quad (6.28)$$

As a result, \widehat{R}_w does not depend on the choice of a decomposition. For $b \in P^\vee$ we have $\mathbb{T}_{t(b)} = \widehat{R}_{t(b)}(1 \otimes t(b))$, so it can be identified with \widehat{Y}^b as defined in § 6.4.

The roots α^k appearing in (6.42) can be characterised geometrically. For this it will be convenient to work over \mathbb{R} , assuming $V \cong \mathbb{R}^n$ and setting $c = 1$ in the definition of $t(\lambda)$ (3.1); then $\delta \equiv 1$ on V . Let C_a be the Weyl alcove,

$$C_a = \{x \in V \mid a_i(x) > 0 \text{ for } i = 0, \dots, n\}.$$

This is a fundamental domain for W_a , and each $\pi \in \Omega$ maps C_a to itself. The set of positive roots R_a^+ consists, by definition, of all those $\tilde{\alpha}$ which take positive values on C_a . Then for any $w \in \widehat{W}$, the set $\{\alpha^1, \dots, \alpha^l\}$ as defined in (6.42) consists of all $\tilde{\alpha} \in R_a^+$ for which the hyperplane $\tilde{\alpha} = 0$ separates C_a and $w(C_a)$, see [M2, 2.2] (note that our w corresponds to w^{-1} in [M2]). Furthermore, the sequence of the hyperplanes $\alpha^1 = 0, \dots, \alpha^l = 0$ can be obtained by taking a straight line between two generic points $x \in C_a$, $y \in wC_a$ and by listing the reflection hyperplanes which this line intersects as you go from x to y . This determines each α^k up to a sign which can be further fixed by prescribing that α^k decreases as you move from x to y (this is because α^k is > 0 on C_a and < 0 on wC_a).

Lemma 6.18. *For any $\alpha \in R$, we have $\widehat{Y}^{\alpha^\vee} = 1$ if $\langle \alpha^\vee, \xi \rangle = 0$.*

Proof. We may assume that $\alpha > 0$. Consider $w = t(\alpha^\vee) \in \widehat{W}$; it acts on V by $w(x) = x - \alpha^\vee$. We need to show that for $\langle \alpha^\vee, \xi \rangle = 0$ we have $\widehat{R}_w = t(-\alpha^\vee)$. Take a generic $x \in C_a$ and consider a line from x to $y = x - \alpha^\vee \in wC_a$. We have $0 < \langle \alpha, x \rangle < 1$ and $\langle \alpha, y \rangle = \langle \alpha, x \rangle - 2$, therefore the line intersects the hyperplanes $\alpha = 0$ and $\alpha + \delta = 0$. This tells us that $\alpha^i = \alpha$ and $\alpha^j = \alpha + \delta$ for some $1 \leq i < j \leq l$. Note that $s_\alpha s_{\alpha+\delta}$ is the translation by α^\vee , which is $t(-\alpha^\vee)$ under our conventions.

Let $z = s_\alpha(x) = s_{\alpha+\delta}(y)$; this point divides the interval between x, y in two parts. The interval between x, z is symmetric under s_α , therefore the collection of the hyperplanes it intersects with will be symmetric as well, giving $\alpha^{i+r} = \pm s_\alpha(\alpha^{i-r})$ for $0 < r < i$. Moreover, the rate of change of each α^k in direction of α^\vee should be positive; this gives $\alpha^{i+r} = -s_\alpha(\alpha^{i-r})$. Now, from the formula (6.22) we get $\widehat{R}(\alpha) = -s_\alpha$ when $\langle \alpha^\vee, \xi \rangle = 0$. Therefore, assuming $\langle \alpha^\vee, \xi \rangle = 0$, we get

$$\widehat{R}(\alpha^1) \dots \widehat{R}(\alpha^{2i-1}) = \widehat{R}(\alpha^1) \dots \widehat{R}(\alpha^{i-1}) (-s_\alpha) \widehat{R}(-s_\alpha(\alpha^{i-1})) \dots \widehat{R}(-s_\alpha(\alpha^1)).$$

The middle factor s_α can be replaced by $s_\alpha^\vee \otimes s_\alpha \in \mathbb{W}$, since s_α^\vee acts trivially if $\langle \alpha^\vee, \xi \rangle = 0$. Using the equivariance (6.25), we can move $s_\alpha^\vee \otimes s_\alpha$ to the left and obtain

$$\widehat{R}(\alpha^1) \dots \widehat{R}(\alpha^{2i-1}) = -\widehat{R}(\alpha^1) \dots \widehat{R}(\alpha^{i-1}) \widehat{R}(-\alpha^{i-1}) \dots \widehat{R}(-\alpha^1) (s_\alpha^\vee \otimes s_\alpha) = -s_\alpha,$$

by unitarity. By the same arguments applied to the interval from z to y , we get $\alpha^{j+r} = -s_{\alpha+\delta}(\alpha^{j-r})$ for all $1 \leq r \leq l - j$ and so the product $\widehat{R}(\alpha^{2i}) \dots \widehat{R}(\alpha^l)$ reduces to $-s_{\alpha+\delta}$. Putting this together gives us $\widehat{R}_w = s_\alpha s_{\alpha+\delta} = t(-\alpha^\vee)$, as needed. \square

Recall that we have previously defined $R_{t(b)}$ and Y^b in accordance with definition 6.1; let us now compare them to $\widehat{R}_{t(b)}$ and \widehat{Y}^b .

Lemma 6.19. *For any $b \in P^\vee$, we have $R_{t(b)} = G_b \widehat{R}_{t(b)}$ and $Y^b = G_b \widehat{Y}^b$, where*

$$G_b(\xi) = \prod_{\substack{\alpha \in R \\ \langle \alpha, b \rangle > 0}} (\sigma_{m_\alpha}(\langle \alpha^\vee, \xi \rangle))^{\langle \alpha, b \rangle}.$$

Proof. This follows directly from the geometric description of the sequence $\alpha^1, \dots, \alpha^l$ given above. The exponent $\langle \alpha, b \rangle$ in the formula counts how many roots of the form $\alpha + k\delta$, $k \in \mathbb{Z}$ will appear in that sequence. \square

Remark 6.20. It follows that the elements Y^b pairwise commute for all $b \in P^\vee$. However, the formula $Y^\lambda Y^\mu = Y^{\lambda+\mu}$ does not hold in general: it is only valid up to a ξ -depending factor.

We also need to know how the classical operators Y_c^b, \widehat{Y}_c^b behave under the shifts $\xi \mapsto \xi + u + \tau v$ with $u, v \in P$.

Lemma 6.21. *Under $\xi \mapsto \xi + u + \tau v$, $u, v \in P$, the operator Y_c^b changes to $e^{2\pi i \langle v, x \rangle} Y_c^b e^{-2\pi i \langle v, x \rangle}$.*

Proof. We have $\sigma_{\mu+1} = \sigma_\mu(z)$ and $\sigma_{\mu+\tau}(z) = e^{2\pi iz} \sigma_\mu(z)$. Hence,

$$R(\tilde{\alpha}) \xrightarrow{\xi \mapsto \xi + \tau v} \sigma_{m_\alpha}(\tilde{\alpha}) - e^{2\pi i \langle \alpha^\vee, v \rangle} \tilde{\alpha} \sigma_{\langle \alpha^\vee, \xi \rangle}(\tilde{\alpha}) s_{\tilde{\alpha}} = e^{2\pi i \langle v, x \rangle} R(\tilde{\alpha}) e^{-2\pi i \langle v, x \rangle}.$$

Also, $t(b) = e^{2\pi i c \langle v, b \rangle} e^{2\pi i \langle v, x \rangle} t(b) e^{-2\pi i \langle v, x \rangle}$. Therefore, for $Y^b = R_{t(b)} t(b)$ we obtain

$$Y^b \xrightarrow{\xi \mapsto \xi + \tau v} e^{2\pi i c \langle v, b \rangle} e^{2\pi i \langle v, x \rangle} Y^b e^{-2\pi i \langle v, x \rangle}.$$

In the classical limit we have $c = 0$, so the lemma follows. \square

Lemma 6.22. *For a function $g(\xi)$ and Y^b , $b \in P^\vee$, we have $\mathbb{T}_w g(\xi) \widehat{Y}^b = g(w^{-1}\xi) \widehat{Y}^{wb} \mathbb{T}_w$ for any $w \in W$.*

Proof. It suffices to check that $\mathbb{T}_i g(\xi) = g(s_i \xi) \mathbb{T}_i$ and $\mathbb{T}_i \widehat{Y}^b = \widehat{Y}^{s_i b} \mathbb{T}_i$ for $i = 1, \dots, n$. The first claim is immediate from (6.21), (6.27). Also, since $\widehat{Y}^b = \mathbb{T}_{t(b)}$, we have $\mathbb{T}_i \widehat{Y}^b = \mathbb{T}_{s_i t(b)} = \mathbb{T}_{t(s_i b) s_i} = \widehat{Y}^{s_i b} \mathbb{T}_i$, as needed. \square

For the rest of this subsection, $b \in P^\vee$ is assumed to be dominant and (quasi-)minuscule.

Lemma 6.23. *The only possible singularities of $L_c^{b, \vee}(\xi, \widehat{Y})$ in the ξ -variable are along the hyperplanes $\langle \alpha^\vee, \xi \rangle = k + l\tau$, with $\alpha \in R$, $k, l \in \mathbb{Z}$.*

Proof. The operators \widehat{Y}^π are defined in terms of $\widehat{R}(\tilde{\alpha})$, which have singularities along hyperplanes $\theta(\langle a_i^\vee, \xi_0 \rangle - m_\alpha) = 0$. However, these are removable due to the coefficients in $L_c^{b, \vee}(\xi, \gamma)$, as can be readily seen from the formulas in Theorem 6.10. Indeed, in the classical limit the coefficients A_π^\vee coincide with $G_b(\xi)$ from Lemma 6.19. Therefore, $A_\pi^\vee \widehat{Y}^\pi = Y^\pi$, and the latter has singularities only where $\theta(\langle \alpha^\vee, \xi \rangle) = 0$. \square

Lemma 6.24. *For $u, v \in P$, we have*

$$L_c^{b, \vee}(\xi, \widehat{Y}_c) \xrightarrow{\xi \mapsto \xi + \tau v} e^{2\pi i \langle v, x \rangle} L_c^{b, \vee}(\xi, \widehat{Y}_c) e^{-2\pi i \langle v, x \rangle}.$$

Proof. If b is minuscule, then from the proof of the previous lemma we see that $L_c^{b, \vee}(\xi, \widehat{Y}_c)$ is a sum of the terms Y_c^π , so the result follows from Lemma 6.21. If b is quasi-minuscule, then we have in addition the sum of B_π^\vee over $\pi \in Wb$, which is a function of ξ only. One checks directly from the formula that each B_π^\vee is elliptic in ξ , in both the quantum and classical settings. This implies that the lemma is true in this case as well. \square

Corollary 6.25. (i) *For any $b \in P^\vee$ we have $\mathbb{T}_i L_c^{b, \vee}(\xi, \widehat{Y}) = L_c^{b, \vee}(\xi, \widehat{Y}) \mathbb{T}_i$ for $i = 1, \dots, n$.*

(ii) *If $\lambda = b_k$ then $\mathbb{T}_i Y^\lambda = Y^\lambda \mathbb{T}_i$ for $i \neq k$.*

Proof. The first statement follows from Lemma 6.22, since the classical Hamiltonian $L_c^{b, \vee}(\xi, \gamma)$ is W -invariant. By the same lemma, $\mathbb{T}_i \widehat{Y}^\lambda = \widehat{Y}^\lambda \mathbb{T}_i$ if $s_i(\lambda) = \lambda$. The second statement now follows from Lemma 6.19, since it is easy to check that the factor G_b in this case is invariant under s_i . \square

Proof of Lemma 6.14. Using the above corollary together with (6.27), we have $\widehat{R}(a_i)(s_i^\vee \otimes s_i) Y^\lambda = Y^\lambda \widehat{R}(a_i)(s_i^\vee \otimes s_i)$ for $i \neq k$, from which $\widehat{R}(-a_i) Y^\lambda = (s_i^\vee \otimes s_i) Y^\lambda (s_i^\vee \otimes s_i) \widehat{R}(-a_i)$. Multiplying this by $\sigma_{m_{a_i}}(-\langle a_i^\vee, \xi \rangle)$, we obtain (cf. [KH2, (4.7)]):

$$R(-a_i) Y^\lambda = (s_i^\vee \otimes s_i) Y^\lambda (s_i^\vee \otimes s_i) R(-a_i) \quad \text{for } i \neq k.$$

If $\xi = -\rho_m + \eta b_k$ then $\langle a_i^\vee, \xi \rangle = -m_{a_i}$ for $i \neq k$. In this case directly from the definitions, $R(-a_i) = \sigma_{m_{a_i}}(-a_i)(1 - s_i)$. So if we use this in the previous relation, and multiply it by e' , we get $(1 - s_i) Y^\lambda e' = 0$. Since this holds for all $i \neq k$, we conclude that Y^λ preserves the subspace $M' = e' M$. The statement about $L_c^{b, \vee}(\xi, \widehat{Y})$ is proved in exactly the same way. \square

Proof of Proposition 6.13. *Part (i).* We pick $\alpha \in R$ and want to show that $L_c^{b, \vee}(\xi, \widehat{Y})$ is regular along $\langle \alpha, \xi \rangle = 0$. Let us first take b to be minuscule. The orbit Wb breaks into pairs π, π' with $\pi' = s_\alpha(\pi)$, plus a number of s_α -invariant π 's. If $s_\alpha(\pi) = \pi$, then the coefficient A_π^\vee does not contain the factor $\sigma_{m_\alpha}(\langle \alpha^\vee, \xi \rangle)$ and so is regular. The other case leads to

$$A \widehat{Y}^\pi + A' \widehat{Y}^{\pi'}, \quad A = A_\pi^\vee, \quad A' = A_{\pi'}^\vee, \quad \pi' = \pi - \alpha^\vee.$$

The coefficients A, A' have first order poles along $\langle \alpha^\vee, \xi \rangle = 0$ and satisfy $A' = A^{s_\alpha}$ due to the W -symmetry of $L^{b, \vee}$. As a result, $A + A'$ is regular along the hyperplane $\langle \alpha^\vee, \xi \rangle = 0$. Now, by

Lemma 6.18 we have $\widehat{Y}^{\pi'} = \widehat{Y}^\pi \widehat{Y}^{\alpha^\vee} = \widehat{Y}^\pi$ at the hyperplane $\langle \alpha^\vee, \xi \rangle = 0$. It follows that at this hyperplane $A\widehat{Y}^\pi + A'\widehat{Y}^{\pi'} = (A + A')\widehat{Y}^\pi$, and this expression is regular.

Now consider the case when $b = \varphi^\vee$ is quasi-minuscule. In this case, there is one additional possibility when $\pi' = s_\alpha(\pi)$ with $\pi' = \pi - 2\alpha^\vee$; this happens only when $\alpha^\vee \in Wb$. In this case, we are led to consider

$$A\widehat{Y}^{\alpha^\vee} + A'\widehat{Y}^{-\alpha^\vee} + B, \quad A = (A_\pi^\vee)_c, \quad A' = (A_{\pi'}^\vee)_c, \quad B = - \sum_{\pi \in Wb} (B_\pi^\vee)_c. \quad (6.29)$$

Note that in the classical limit we have $\delta = 0$, so the formulas (6.12)–(6.13) tell us that the coefficients A, A', B will have second order poles along $\langle \alpha^\vee, \xi \rangle = 0$ (the singularity in B comes from $B_\alpha^\vee + B_{-\alpha^\vee}^\vee$). We have the following properties of A, A', B , first two of which follow from the W -symmetry of $L^{b,\vee}$, and the last one can be checked by inspecting the formulas (6.12)–(6.13) (in the case $\delta = 0$):

$$\begin{aligned} A' &= A^{s_\alpha}, \quad B = B^{s_\alpha}; \\ A + A' \text{ and } B &\text{ have zero residue at } \langle \alpha^\vee, \xi \rangle = 0; \\ A + A' + B &\text{ has at most simple pole along } \langle \alpha^\vee, \xi \rangle = 0. \end{aligned} \quad (6.30)$$

In addition, it follows from Lemma 6.18 that

$$\widehat{Y}^{\alpha^\vee} = 1 + \varepsilon \langle \alpha^\vee, \xi \rangle + o(\langle \alpha^\vee, \xi \rangle) \quad \text{near } \langle \alpha^\vee, \xi \rangle = 0,$$

for some $\varepsilon \in \mathbb{C}(\mathbb{V}) \rtimes \mathbb{C}\widehat{W}$, which is regular along $\langle \alpha^\vee, \xi \rangle = 0$. It follows that $\widehat{Y}^{-\alpha^\vee} = (\widehat{Y}^{\alpha^\vee})^{-1}$ satisfies

$$\widehat{Y}^{-\alpha^\vee} = 1 - \varepsilon \langle \alpha^\vee, \xi \rangle + o(\langle \alpha^\vee, \xi \rangle) \quad \text{near } \langle \alpha^\vee, \xi \rangle = 0.$$

Putting this together, one sees that (6.29) is regular at $\langle \alpha^\vee, \xi \rangle = 0$, which proves part (i).

Parts (ii), (iii). These are proved in the same way as in Proposition 5.1, using Lemmas 6.23, 6.24.

Part (iv). It is sufficient to prove that for $\xi = -\rho_m$, we have $L_c^{b,\vee}(\xi, \widehat{Y})e = (L_c^b + \text{const})e$. Indeed, the statement (iv) is then obtained by passing to the classical limit and by using part (iii). We already know that $L_c^{b,\vee}(\xi, \widehat{Y}) = \sum_{\pi \in Wb} Y^\pi + c_0$, where c_0 is a function of ξ . Therefore, it suffices to prove that for $\xi = -\rho_m$ we have $Y^\pi e = 0$ if π is not dominant. Recall that

$$Y^\pi = R(\alpha^1) \dots R(\alpha^l) t(\pi),$$

where the sequence $\alpha^1, \dots, \alpha^l$ is obtained by going from $x \in C_a$ to $x - \pi$ and listing all reflection hyperplanes transversed in that process; in addition, each α^i should be positive at x and negative at $x - \pi$. Using the translation properties (6.26), we can rewrite this as

$$Y^\pi = t(\pi) R(\beta^1) \dots R(\beta^l),$$

where the sequence of affine roots β^i is obtained similarly by going from $x + \pi$ to $x \in C_a$, listing all transversed hyperplanes. Since x lies inside the Weyl alcove C_a , the last root β^l should correspond to one of the faces of the alcove; also, we know that $\beta^l(x) < 0$. This tells us that $\beta^l = -a_i$ with $0 \leq i \leq n$. Moreover, if we choose x inside C_a sufficiently close to 0, then we necessarily have $\beta^l = -a_i$ with $i \neq 0$ (otherwise π would be dominant). Now, for $\xi = -\rho_m$ we have $\langle a_i^\vee, \xi \rangle = -m_{a_i}$, and so $R(-a_i) = \sigma_{m_{a_i}}(-a_i)(1 - s_i)$ as a result. Therefore, $R(\beta^l)e = R(-a_i)e = 0$, implying $Y^\pi e = 0$. \square

6.6. Let us proceed to the case of the affine root system $C^\vee C_n$. In the setting of [KH2] this corresponds to the case of a reduced root system $R = BC_n$, but for us, as in Section §4, R will be a root system (4.1) of type C_n , and R_a will denote the associated affine root system (3.5), with the following basis of simple roots:

$$a_0 = \delta - 2\epsilon_1, \quad a_i = \epsilon_i - \epsilon_{i+1} \quad (i = 1, \dots, n-1), \quad a_n = 2\epsilon_n,$$

where $\delta \equiv c$ on V . We have the Weyl group $W = \mathfrak{S}_n \times \{\pm 1\}^n$ and the group $\widehat{W} = W_a$ generated by $s_i = s_{a_i}$, acting on V in accordance with (4.2). Note that $\widehat{W} \cong W \times \Lambda$ with $\Lambda = \sum_{i=1}^n \mathbb{Z}\epsilon_i$. The dominant cone is $\Lambda_+ = \{(\lambda_1, \dots, \lambda_n) \in \Lambda \mid \lambda_1 \geq \dots \geq \lambda_n \geq 0\}$.

The R -matrices are elements of the algebra $\mathbb{C}(V) * \widehat{W} \cong \mathcal{D}_q * W$ of reflection-difference operators on V . They depend on dynamical parameters $\xi \in V$, coupling constants $\mu, \nu, \bar{\nu}, g = (g_i), \bar{g} = (\bar{g}_i)$ ($i = 0 \dots 3$), and are as follows [KH2]:

$$R(\tilde{\alpha}) = \sigma_\mu(\tilde{\alpha}) - \sigma_{\langle \alpha^\vee, \xi \rangle}(\tilde{\alpha}) s_{\tilde{\alpha}} \quad \text{for } \tilde{\alpha} = k\delta \pm \epsilon_i \pm \epsilon_j \quad (k \in \mathbb{Z}, i \neq j), \quad (6.31)$$

$$R(\tilde{\alpha}) = v_{\nu, g}(\tilde{\alpha}/2) - v_{\langle \alpha^\vee, \xi \rangle, g}(\tilde{\alpha}/2) s_{\tilde{\alpha}} \quad \text{for } \tilde{\alpha} = 2k\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}), \quad (6.32)$$

$$R(\tilde{\alpha}) = v_{\bar{\nu}, \bar{g}}(\tilde{\alpha}/2) - v_{\langle \alpha^\vee, \xi \rangle, \bar{g}}(\tilde{\alpha}/2) s_{\tilde{\alpha}} \quad \text{for } \tilde{\alpha} = (2k+1)\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}). \quad (6.33)$$

In these formulas, $v_{\nu, g}(z) = v_\nu(z; g_0, g_1, g_2, g_3)$ is the function (5.10).

According to [KH2, Theorems 4.1, 4.2], the elements $R(\tilde{\alpha})$ satisfy the affine Yang–Baxter relations. We can now define the elements R_w and $Y^b = R_{t(b)}t(b)$, $b \in \Lambda$ in the same way as in 6.1 (note that the group Ω is trivial in this case), and Theorem 6.2 remains valid in this setting.

Let us write down the expressions for $Y^{\pm \epsilon_i}$. We have a reduced decomposition (cf. [St, (3.6)])

$$t(\epsilon_i) = s_i \dots s_{n-1} s_n s_{n-1} \dots s_1 s_0 s_1 \dots s_{i-1}, \quad 1 \leq i \leq n.$$

Applying Definition 6.1, we calculate

$$\begin{aligned} Y^{\epsilon_i} &= R(\epsilon_i - \epsilon_{i+1}) R(\epsilon_i - \epsilon_{i+2}) \dots R(\epsilon_i - \epsilon_n) R(2\epsilon_i) \\ &\quad \times R(\epsilon_i + \epsilon_n) \dots R(\epsilon_i + \epsilon_{i+1}) R(\epsilon_i + \epsilon_{i-1}) \dots R(\epsilon_i + \epsilon_1) \\ &\quad \times R(\delta + 2\epsilon_i) t(\epsilon_i) R(\epsilon_i - \epsilon_1) \dots R(\epsilon_i - \epsilon_{i-1}). \end{aligned} \quad (6.34)$$

Remark 6.26. In [KH2], a slightly different affine root system is considered, with the roots in (6.32) replaced by $k\delta \pm \epsilon_i$, but the corresponding R -matrices and the resulting elements R_w are the same.

For $i = 1, \dots, n$ define m_i by $m_1 = \dots = m_{n-1} = \mu$, $m_n = \nu$. Let $\xi = \xi_0$ be a solution to the system of equations

$$\langle a_i^\vee, \xi_0 \rangle = -m_i, \quad i = 1, \dots, n. \quad (6.35)$$

Explicitly, we have $\xi_0 = (\xi_1, \dots, \xi_n)$ with $\xi_i = -\nu - (n-i)\mu$.

Theorem 6.27 ([KH2, Theorems 4.5 & 6.5]). *Let $\xi = \xi_0$ as above.*

(1) *Given $b \in \Lambda_+$, let $L^b \in \mathcal{D}_q$ be the unique difference operator such that $Y^b e = L^b e$. Then each L^b is W -invariant, and the difference operators L^b , $b \in \Lambda_+$ form a commutative family.*

(2) *Let $b = \epsilon_1$. Then*

$$L^b = \sum_{\pi \in Wb} (A_\pi t(\pi) - B_\pi), \quad A_\pi = v_{\nu, g}(\pi) v_{\bar{\nu}, \bar{g}}(\pi + \delta/2) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle = 1}} \sigma_\mu(\alpha), \quad (6.36)$$

$$B_\pi = v_{\nu, g}(\pi) v_{-\nu - (n-1)\mu, \bar{g}}(\pi + \delta/2) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle = 1}} \sigma_\mu(\alpha). \quad (6.37)$$

The operator (6.36)–(6.37) contains 11 parameters $\mu, \nu, \bar{\nu}, g_i, \bar{g}_i$, but multiplying all g_i (or all \bar{g}_i) by a constant simply rescales the Hamiltonian. Thus, effectively we have 9 coupling parameters. This Hamiltonian was first introduced by van Diejen [vD1], in a different form and under an additional constraint on $\nu, \bar{\nu}$. For general coupling parameters it was introduced by Komori and Hikami in [KH1], where higher quantum Hamiltonians were also constructed. See [KH1, (4.21)] for an alternative presentation of L^b which links it to [vD1]. The classical Hamiltonian L_c^b looks as follows:

$$L_c^b = \sum_{\pi \in Wb} (A_\pi e^{\beta p_\pi} - B_\pi), \quad A_\pi = v_{\nu, g}(\pi) v_{\bar{\nu}, \bar{g}}(\pi) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle = 1}} \sigma_\mu(\alpha), \quad (6.38)$$

$$B_\pi = v_{\nu, g}(\pi) v_{-\nu - (n-1)\mu, \bar{g}}(\pi) \prod_{\substack{\alpha \in R \\ \langle \pi, \alpha \rangle = 1}} \sigma_\mu(\alpha). \quad (6.39)$$

6.7. Before proceeding to constructing a quantum Lax pair, we need to renormalise the R -matrices. The R -matrices (6.31) have the property (6.2), so we can define $\widehat{R}(\tilde{\alpha})$ by (6.21), with $m_\alpha = \mu$. For the R -matrices (6.32)–(6.33) the procedure is more subtle. First, considering (6.32), we obtain:

$$R(\tilde{\alpha}) R(-\tilde{\alpha}) = v_{\nu, g}(\tilde{\alpha}/2) v_{\nu, g}(-\tilde{\alpha}/2) + v_{\langle \alpha^\vee, \xi \rangle, g}(\tilde{\alpha}/2) v_{-\langle \alpha^\vee, \xi \rangle, g}(\tilde{\alpha}/2). \quad (6.40)$$

Using (5.11), this can be rewritten as (cf. [KH2, (4.5)])

$$R(\tilde{\alpha})R(-\tilde{\alpha}) = \sum_{r=0}^3 (g_r^\vee)^2 (\wp(\nu + \omega_r) - \wp(\langle \alpha^\vee, \xi \rangle + \omega_r)),$$

which is independent of x .

Definition 6.28. Let $v_{\nu,g}(z)$ be the function (5.10) with parameters $\nu, g = (g_0, g_1, g_2, g_3)$. The dual parameters $\nu^\vee, g^\vee = (g_0^\vee, g_1^\vee, g_2^\vee, g_3^\vee)$ are defined by (5.12) and the condition $v_{\nu,g}(\nu^\vee) = 0$.

Remark 6.29. The function $v_{\nu,g}(z)$ can be parametrised by its zeros. Following [KH1], let $v_{\nu,g}(z) = A \prod_{r=0}^3 \sigma_{\nu_r}^r(z)$. Then we have $2\nu = \sum_{r=0}^3 \nu_r$ and g_r can be expressed in terms of A and ν_r , see [KH1, Lemma 4.5]. If such a parametrisation is used, then ν^\vee can be taken simply as ν_0 (or any of $\nu_r + \omega_r, r = 0 \dots 3$). However, the parametrisation of $v_{\nu,g}$ by its zeros is inconvenient for writing the R -matrices.

Lemma 6.30. Let $\tilde{\alpha} \in R_a$ and $R(\tilde{\alpha})$ be as in (6.32). Then $\widehat{R}(\tilde{\alpha}) := v_{\nu^\vee, g^\vee}(\langle \alpha^\vee, \xi \rangle)^{-1} R(\tilde{\alpha})$ satisfies $\widehat{R}(\tilde{\alpha})\widehat{R}(-\tilde{\alpha}) = 1$. We have $\widehat{R}(\tilde{\alpha}) = -e^{\pi i(\tilde{\alpha} - 2\nu^\vee)\beta_r} s_{\tilde{\alpha}}$ when $\langle \alpha^\vee, \xi \rangle = \omega_r, r = 0 \dots 3$. Here ω_r are the half-periods, and $(\beta_0, \beta_1, \beta_2, \beta_3) = (0, 0, 1, 1)$. Similarly, for the case (6.33) we define $\widehat{R}(\tilde{\alpha}) := v_{\bar{\nu}^\vee, \bar{g}^\vee}(\langle \alpha^\vee, \xi \rangle)^{-1} R(\tilde{\alpha})$. Then $\widehat{R}(\tilde{\alpha}) = -e^{\pi i(\tilde{\alpha} - 2\bar{\nu}^\vee)\beta_r} s_{\tilde{\alpha}}$ when $\langle \alpha^\vee, \xi \rangle = \omega_r$.

Proof. Since the expression (6.40) is independent of x , we may assume $\tilde{\alpha}/2 = \nu^\vee$. Then

$$R(\tilde{\alpha})R(-\tilde{\alpha}) = v_{\langle \alpha^\vee, \xi \rangle, g}(\nu^\vee) v_{-\langle \alpha^\vee, \xi \rangle, g}(\nu^\vee) = v_{\nu^\vee, g^\vee}(\langle \alpha^\vee, \xi \rangle) v_{\nu^\vee, g^\vee}(-\langle \alpha^\vee, \xi \rangle),$$

by (5.13). Hence, $\widehat{R}(\tilde{\alpha})\widehat{R}(-\tilde{\alpha}) = 1$.

From the definition of $R(\tilde{\alpha})$ we see that it has a first order pole at the hyperplane $\langle \alpha^\vee, \xi \rangle = 0$, with the residue equal to $-1/2(\sum_{r=0}^3 g_r) s_{\tilde{\alpha}}$. On the other hand, $v_{\nu^\vee, g^\vee}(\langle \alpha^\vee, \xi \rangle)$ also has a first order pole at $\langle \alpha^\vee, \xi \rangle = 0$, with the residue equal to g_0^\vee , which is $1/2(\sum_{r=0}^3 g_r)$ from (5.12). This implies that $\widehat{R}(\tilde{\alpha})$ tends to $-s_{\tilde{\alpha}}$ as $\langle \alpha^\vee, \xi \rangle$ approaches zero. This proves the claim for $\omega_r = 0$; for other half-periods ω_r proof is similar. \square

Motivated by the above, for given ν, g and $\bar{\nu}, \bar{g}$ we choose (and fix) $\nu^\vee, \bar{\nu}^\vee$ and define *unitary R -matrices* by

$$\widehat{R}(\tilde{\alpha}) = \begin{cases} \sigma_\mu(\langle \alpha^\vee, \xi \rangle)^{-1} R(\tilde{\alpha}) & \text{for } \tilde{\alpha} = k\delta \pm \epsilon_i \pm \epsilon_j \quad (k \in \mathbb{Z}, i \neq j), \\ \widehat{R}(\tilde{\alpha}) = v_{\nu^\vee, g^\vee}(\langle \alpha^\vee, \xi \rangle)^{-1} R(\tilde{\alpha}) & \text{for } \tilde{\alpha} = 2k\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}), \\ \widehat{R}(\tilde{\alpha}) = v_{\bar{\nu}^\vee, \bar{g}^\vee}(\langle \alpha^\vee, \xi \rangle)^{-1} R(\tilde{\alpha}) & \text{for } \tilde{\alpha} = (2k+1)\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}). \end{cases}$$

Then we have $\widehat{R}(\tilde{\alpha})\widehat{R}(-\tilde{\alpha}) = 1$ for all $\tilde{\alpha} \in R_a$.

6.8. Now let us recall the notation of §6.5, by which we view $R(\tilde{\alpha})$ and $\widehat{R}(\tilde{\alpha})$ as elements of $\mathbb{C}(\mathbb{V}) \rtimes \mathbb{C}W$, where $\mathbb{V} = V^\vee \times V$ incorporates the dynamical variables. Both $R(\tilde{\alpha})$ and $\widehat{R}(\tilde{\alpha})$ satisfy the affine Yang–Baxter relations [KH2, (3.1)(a)-(c)]. This allows us to introduce

$$\mathbb{T}_i = \widehat{R}(a_i)(s_i^\vee \otimes s_i), \quad i = 0, \dots, n. \quad (6.41)$$

Then one checks that the relations [KH2, (3.1)(a)-(c)] imply that \mathbb{T}_i satisfy the relations (4.3)–(4.5), while the unitarity of $\widehat{R}(\tilde{\alpha})$ implies that $\mathbb{T}_i^2 = 1$ for all i . Therefore, we have the following result, analogous to Propositions 6.16, 6.17.

Proposition 6.31. (1) The assignment $s_i \mapsto \mathbb{T}_i$ ($i = 0, \dots, n$) extends to a group homomorphism $w \mapsto \mathbb{T}_w, w \in \widehat{W}$. In particular, the elements $\mathbb{T}_{t(b)}, b \in \Lambda$ pairwise commute.

(2) For any reduced decomposition of w into $w = s_{i_1} \dots s_{i_l}$ we have $\mathbb{T}_w = \widehat{R}_w(w^\vee \otimes w)$, where

$$\widehat{R}_w = \widehat{R}(\alpha^1) \dots \widehat{R}(\alpha^l), \quad \alpha^1 = a_{i_1}, \alpha^2 = s_{i_1}(a_{i_2}), \dots, \alpha^l = s_{i_1} \dots s_{i_{l-1}}(a_{i_l}). \quad (6.42)$$

As a result, \widehat{R}_w does not depend on the choice of a decomposition. For $b \in \Lambda$ we have $\mathbb{T}_{t(b)} = \widehat{R}_{t(b)}(1 \otimes t(b)) = \widehat{Y}^b$. This implies the commutativity of \widehat{Y}^b and, therefore, of $Y^b := R_{t(b)}t(b)$.

We also have an analogue of Proposition 6.13. Let $b = \epsilon_1$ and L_c^b is the classical van Diejen Hamiltonian (6.38)–(6.39). By $L_c^{b,\vee}$ we denote the classical operator with the dual coupling parameters $\nu^\vee, \bar{\nu}^\vee, g^\vee, \bar{g}^\vee$ (and with $\mu^\vee = \mu$), see Definition 6.28.

Proposition 6.32. *Let $L_c^{b,\vee}(\xi, \widehat{Y})$ and $L_c^{b,\vee}(\xi, \widehat{Y}_c)$ be the result of substituting the dynamical variables and Cherednik operators into the dual classical Hamiltonian. Then we have:*

- (i) $L_c^{b,\vee}(\xi, \widehat{Y})$, viewed as an element of $\mathcal{D}(V) * W$ depending on ξ , is regular for ξ near $\xi = 0$;
- (ii) $L_c^{b,\vee}(\xi, \widehat{Y}_c)$, viewed as an element of $\mathbb{C}(V \times V) * W$ depending on ξ , is regular for all $\xi \in V$;
- (iii) $L_c^{b,\vee}(\xi, \widehat{Y}_c)$ is constant in ξ . Moreover, expanding $L_c^{b,\vee}(\xi, \widehat{Y}_c)$ as $\sum_{w \in W} a_w w$ with $a_w \in \mathbb{C}(V \times V)$, we have $a_w = 0$ for $w \neq \text{id}$.
- (iv) We have $L_c^{b,\vee}(\xi, \widehat{Y}_c) = L_c^b + \text{const}$.

This is proved in the same way as Proposition 6.13. One additional complication arises when considering what happens for $\xi_i = \omega_r$ (this is further explained in a more general situation in the next subsection). Here is a suitable generalisation of Lemma 6.18, which is proved by a similar argument.

Lemma 6.33. *For $\alpha = 2\epsilon_l$ we have $\widehat{Y}^{\alpha^\vee} = e^{-\pi c \beta_r - \lambda_r}$ if $\langle \alpha^\vee, \xi \rangle = \omega_r$, where*

$$\lambda_r = 2\pi i \beta_r (\nu^\vee + \bar{\nu}^\vee + (n-1)\mu), \quad r = 0 \dots 3, \quad (6.43)$$

and β_r are the same as in Lemma 6.30. In particular, for the classical operator $\widehat{Y}_c^{\alpha^\vee}$ we have $\widehat{Y}_c^{\alpha^\vee} = e^{-\lambda_r}$ if $\langle \alpha^\vee, \xi \rangle = \omega_r$.

The same arguments as in Section §6.4 lead to a construction of a Lax pair.

Theorem 6.34. *For any fundamental coweight λ with the stabiliser W' , there exists a quantum Lax pair \mathcal{L}, \mathcal{A} of size $|W|/|W'|$ satisfying the Lax equation (1.6) with $\mathcal{H} = \widehat{H}\mathbb{1}$, $\widehat{H} = L^{\epsilon_1}$. This Lax pair depends on a spectral parameter and admits a classical limit.*

The smallest Lax pair of size $2n$ is obtained for $\lambda = \epsilon_1$. We calculate the corresponding Lax matrix in §6.10.

6.9. Let us generalise Proposition 6.32 to any of the higher Hamiltonians of the van Diejen system. A direct proof is problematic since we do not know an explicit formula for these Hamiltonians. Instead, we will use a result of Rains, who in [Ra] developed a geometric approach to elliptic DAHAs. To formulate his result we will need some notation. Let us introduce

$$v_{\tilde{\alpha}} = \begin{cases} \sigma_\mu(\tilde{\alpha}) & \text{for } \tilde{\alpha} = k\delta \pm \epsilon_i \pm \epsilon_j \quad (k \in \mathbb{Z}, i \neq j), \\ v_{\nu, g}(\tilde{\alpha}/2) & \text{for } \tilde{\alpha} = 2k\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}), \\ v_{\bar{\nu}, \bar{g}}(\tilde{\alpha}/2) & \text{for } \tilde{\alpha} = (2k+1)\delta \pm 2\epsilon_i \quad (k \in \mathbb{Z}), \end{cases}$$

as well as \widehat{G}_π , $\pi \in \Lambda$:

$$\widehat{G}_\pi = \prod_{\substack{\alpha \in R: \langle \pi, \alpha \rangle > 0 \\ 0 \leq k \leq \langle \pi, \alpha \rangle}} v_{\alpha + k\delta} \quad (\widehat{G}_0 := 1).$$

The coupling parameters $\mu, \nu, \bar{\nu}, g_i, \bar{g}_i$ will be assumed generic. Next, let

$$\Pi = \{\lambda = (\lambda_1, \dots, \lambda_n) \mid \lambda_i \in \{-1, 0, 1\}\}.$$

For any $\alpha \in R$, define an α^\vee -string in Π as $\{\pi + \mathbb{Z}\alpha^\vee\} \cap \Pi$ where $\pi \in \Pi$; the number of the lattice points on a string will be called its *length*. It is easy to see that α^\vee -strings in Π are of one of the following types:

$$\text{length-one: } \{\pi\}, \quad \text{with } s_\alpha \pi = \pi, \quad (6.44)$$

$$\text{length-two: } \{\pi, \pi'\}, \quad \text{with } s_\alpha \pi = \pi', \quad \pi' - \pi = \pm \alpha^\vee, \quad (6.45)$$

$$\text{length-three: } \{\pi, \pi \pm \alpha^\vee\}, \quad \text{with } s_\alpha \pi = \pi. \quad (6.46)$$

Definition 6.35. For given generic $\mu, \nu_r, \bar{\nu}_r$, let \mathcal{V} be the vector space of all difference operators $L \in \mathcal{D}_q(V)$ of the form

$$L = \sum_{\pi \in \Pi} a_\pi t(\pi), \quad a_\pi \in \mathbb{C}(V), \quad (6.47)$$

satisfying the following conditions: (1) L is W -invariant; (2) a_π / \widehat{G}_π is elliptic w.r.t. the lattice $\Lambda + \tau\Lambda$; (3) a_π and a_π / \widehat{G}_π have at most simple poles along the hyperplanes $\tilde{\alpha} = l + m\tau$ with $\tilde{\alpha} \in R_\alpha$ and $l, m \in \mathbb{Z}$, and no other singularities.

In addition, we are going to impose certain “residue conditions” on the coefficients a_π for each α^\vee -string in Π . In what follows, we call a function $f \in \mathbb{C}(V)$ α -regular if it has no singularities along hyperplanes $\langle \alpha, x \rangle = \text{const}$. For the length-one strings (6.44) the residue condition is simply that

$$a_\pi \text{ is } \alpha\text{-regular.} \quad (6.48)$$

For the length-two strings (6.45) the conditions are that

$$\theta(\alpha)a_\pi \text{ and } \theta(\alpha)a_{\pi'} \text{ are } \alpha\text{-regular.} \quad (6.49)$$

Here, as before, we view the roots as affine-linear functions, e.g., $\theta(\alpha) = \theta(\langle \alpha, x \rangle)$ for $\alpha \in R$. For the length-three strings (6.46), the conditions are more involved:

$$\theta(\alpha)\theta(\alpha + \delta)a_{\pi+\alpha^\vee}, \theta(\alpha + \delta)\theta(\alpha - \delta)a_\pi \text{ and } \theta(\alpha)\theta(\alpha - \delta)a_{\pi-\alpha^\vee} \text{ are } \alpha\text{-regular;} \quad (6.50)$$

$$a_{\pi+\alpha^\vee} + a_\pi + a_{\pi-\alpha^\vee} \text{ is regular for } \langle \alpha, x \rangle = 0, \pm c. \quad (6.51)$$

Additionally, for $\alpha = 2\epsilon_l$ we require that

$$e^{-\lambda_r}a_{\pi+\alpha^\vee} + a_\pi + e^{\lambda_r}a_{\pi-\alpha^\vee} \text{ is regular for } \langle \alpha, x \rangle = 2\omega_r, 2\omega_r \pm c, \quad (6.52)$$

where $\lambda_r = 2\pi i\beta_r(\nu + \bar{\nu} + (n-1)\mu)$, cf. (6.43). See [RaR] where similar residue conditions were considered in the rank-one case.

The following result can be extracted from [Ra].

Theorem 6.36 (cf. Theorem 7.22 of [Ra]). *For each $b_k = \epsilon_1 + \dots + \epsilon_k$ ($0 \leq k \leq n$), there exists an operator $L_k \in \mathcal{V}$ with the leading terms $\sum_{\pi \in W_{b_k}} \widehat{G}_\pi t(\pi)$, whose coefficients satisfy the residue conditions (6.48)–(6.52) for any α^\vee -string, $\alpha \in R$. The operators L_k pairwise commute and admit a classical limit.*

Remark 6.37. The Hamiltonian L_0 is trivial, $L_0 = 1$. It is expected that L_k are the Hamiltonians of the van Diejen system, but the construction below makes no use of this or the commutativity of L_k . Note that for the first van Diejen Hamiltonian L^{ϵ_1} (6.36)–(6.37) the residue conditions can be easily checked from the explicit formula.

Remark 6.38. In the classical limit, the residue conditions for α^\vee -strings of length one or two remain the same. For a length-three string, (6.50)–(6.51) are replaced with their $c = 0$ limit:

$$a_{\pi+\alpha^\vee}\theta^2(\alpha), a_\pi\theta^2(\alpha) \text{ and } a_{\pi-\alpha^\vee}\theta^2(\alpha) \text{ are } \alpha\text{-regular;} \quad (6.53)$$

$$a_{\pi+\alpha^\vee} + a_\pi + a_{\pi-\alpha^\vee} \text{ is regular for } \langle \alpha, x \rangle = 0. \quad (6.54)$$

Similarly, for $\alpha = 2\epsilon_l$ and $r = 1, 2, 3$ the conditions (6.52) are replaced with:

$$e^{-\lambda_r}a_{\pi+\alpha^\vee} + a_\pi + e^{\lambda_r}a_{\pi-\alpha^\vee} \text{ is regular for } \langle \alpha, x \rangle = 2\omega_r. \quad (6.55)$$

This tells us that the second-order poles in this sum must cancel. Also, from $\langle \alpha, \pi \rangle = 0$ it can be checked that \widehat{G}_π is periodic with respect to translations by α^\vee and $\tau\alpha^\vee$, and so must be a_π , by the definition of \mathcal{V} . Together with the s_α -invariance of a_π , this gives that a_π has zero residue at $\langle \alpha, x \rangle = 2\omega_r$. As a result, $e^{-\lambda_r}a_{\pi+\alpha^\vee} + e^{\lambda_r}a_{\pi-\alpha^\vee}$ also have zero residue at $\langle \alpha, x \rangle = 2\omega_r$ (cf. the properties (6.30)).

We now have the following analogue of Proposition (6.32).

Proposition 6.39. *Consider the classical limits of the operators L_k from Theorem 6.36, and with the dual parameters $\nu^\vee, g^\vee, \bar{\nu}^\vee, \bar{g}^\vee$. Denote these classical Hamiltonians as L_k^\vee , $k = 0, \dots, n$, with $L_0^\vee = 1$. Let $L_k^\vee(\xi, \widehat{Y})$ and $L_k^\vee(\xi, \widehat{Y}_c)$ denote the result of substituting the dynamical variables and Cherednik operators into L_k^\vee . Then we have:*

- (i) $L_k^\vee(\xi, \widehat{Y})$, viewed as an element of $\mathcal{D}(V) * W$ depending on ξ , is regular for ξ near $\xi = 0$;
- (ii) $L_k^\vee(\xi, \widehat{Y}_c)$, viewed as an element of $\mathbb{C}(V \times V) * W$ depending on ξ , is regular for all $\xi \in V$;
- (iii) $L_k^\vee(\xi, \widehat{Y}_c)$ is constant in ξ . Moreover, expanding $L_k^\vee(\xi, \widehat{Y}_c)$ as $\sum_{w \in W} a_w w$, we have $a_w = 0$ for $w \neq \text{id}$.

(iv) We have $L_k^\vee(\xi, \widehat{Y}_c) = L_c^{b_k} + \sum_{0 \leq l < k} a_{kl} L_c^{b_l}$, with some $a_{kl} \in \mathbb{C}$, where $L_c^{b_l}$ are the higher (classical) van Diejen Hamiltonians for $b_l = \epsilon_1 + \dots + \epsilon_l$.

This is proved similarly to Proposition 6.13. Namely, the regularity of $L_k^\vee(\xi, \widehat{Y}_c)$ for ξ close to 0, as well as for $\xi_i = \omega_r$, follows from the residue conditions on the coefficients of L_k^\vee , as specified in Remark 6.38 (used together with Lemmas 6.18, 6.33). The global regularity then follows from an analogue of Lemma 6.24; to prove such an analogue we do not need an explicit formula for L_k^\vee , but make use of the Definition 6.35 instead. We leave the details to the reader. \square

As a consequence, we can perform a construction of a Lax pair for each of the higher van Diejen Hamiltonians, and so Theorem 6.34 remains valid for any of L^b with $b = b_k$, $k = 1, \dots, n$.

6.10. Let us calculate a quantum Lax matrix corresponding to $Y_1 = Y^{\epsilon_1}$. The calculation and the notation will be very similar to those in § 4.2. The stabiliser W' of $\lambda = \epsilon_1$ is the subgroup of signed permutations of x_2, \dots, x_n . The dynamical variable ξ needs to satisfy the conditions $\langle a_i^\vee, \xi \rangle = -m_i$ for $i = 2, \dots, n$, cf. (6.35), that is,

$$\xi_1 = \eta, \quad \xi_i = -\nu - (n-i)\mu \quad (i = 2, \dots, n), \quad (6.56)$$

with η being a spectral parameter. Let us abbreviate $R(\epsilon_i - \epsilon_j)$ and $R(\epsilon_i + \epsilon_j)$ to R_{ij} and R_{ij}^+ , respectively. Using (6.34), we get

$$Y_1 = R_{12}R_{13} \dots R_{1n} R(2\epsilon_1)R_{1n}^+ \dots R_{12}^+ R(\delta + 2\epsilon_1) t(\epsilon_1), \quad \text{with}$$

$$R_{ij} = \sigma_\mu(x_{ij}) - \sigma_{\xi_{ij}}(x_{ij})\mathfrak{s}_{ij}, \quad R_{ij}^+ = \sigma_\mu(x_{ij}^+) - \sigma_{\xi_{ij}^+}(x_{ij}^+)\mathfrak{s}_{ij}^+,$$

$$R(2\epsilon_1) = v_{\nu,g}(x_1) - v_{\xi_1,g}(x_1)\mathfrak{s}_1, \quad R(\delta + 2\epsilon_1) t(\epsilon_1) = v_{\bar{\nu},\bar{g}}(x_1 + c/2) t(\epsilon_1) - v_{\xi_1,\bar{g}}(x_1 + c/2)\mathfrak{s}_1.$$

Introducing $\mathcal{R} = R_{12} \dots R_{1n}$ and $\mathcal{R}^+ = R_{1n}^+ \dots R_{12}^+$, we have the following.

Lemma 6.40 (cf. [KH1], Lemma 4.4). *\mathcal{R} and \mathcal{R}^+ preserve the subspace $M'' = e''M$, where $e'' = \frac{1}{n!} \sum_{w \in \mathfrak{S}_{n-1}} w$. Their restriction onto M'' is calculated as follows:*

$$\begin{aligned} \mathcal{R}|_{M''} &= U - \sum_{i \neq 1}^n V_i \mathfrak{s}_{1i}, & U &= \prod_{l \neq 1}^n \sigma_\mu(x_{1l}), & V_i &= \sigma_{\xi_{12}}(x_{1i}) \prod_{l \neq 1, i} \sigma_\mu(x_{il}), \\ \mathcal{R}^+|_{M''} &= U^+ - \sum_{i \neq 1}^n V_i^+ \mathfrak{s}_{1i}^+, & U^+ &= \prod_{l \neq 1}^n \sigma_\mu(x_{1l}^+), & V_i^+ &= \sigma_{\xi_{1n}^+}(x_{1i}^+) \prod_{l \neq 1, i} \sigma_\mu(x_{li}). \end{aligned}$$

Proof. Note that the dynamical variables satisfy $\xi_{i+1} - \xi_i = -\mu$ for $1 < i < n$, so the statement about \mathcal{R} is known already from Lemmas 6.11, 6.12. The result for \mathcal{R}^+ follows by observing that it can be obtained as $(\omega^\vee \otimes \omega)\mathcal{R}(\omega^\vee \otimes \omega)^{-1}$, where $\omega \in W'$ is the transformation (4.13). Note that under this transformation, the dynamical variables change to $(\xi_1, -\xi_n, \dots, -\xi_2)$, so $\xi_{12} = \xi_1 - \xi_2$ becomes $\xi_{1n}^+ = \xi_1 + \xi_n$. \square

Next, we restrict Y_1 further onto $M' = e'M$, as in Section § 4.2. From the above,

$$Y_1|_{M'} = (U - \sum_{i \neq 1}^n V_i \mathfrak{s}_{1i}) R(2\epsilon_1) (U^+ - \sum_{i \neq 1}^n V_i^+ \mathfrak{s}_{1i}^+) R(\delta + 2\epsilon_1) t(\epsilon_1)|_{M'}. \quad (6.57)$$

The main step is to work out the restriction onto M' of the product

$$(U - \sum_{i \neq 1}^n V_i \mathfrak{s}_{1i})(v_{\nu,g}(x_1) - v_{\xi_1,g}(x_1)\mathfrak{s}_1)(U^+ - \sum_{i \neq 1}^n V_i^+ \mathfrak{s}_{1i}^+). \quad (6.58)$$

Lemma 6.41. *The operator (6.58) preserves the subspace $M' = e'M$, and its restriction onto M' is given by $A + B\mathfrak{s}_1 - \sum_{i \neq 1} (C_i \mathfrak{s}_{1i} + D_i \mathfrak{s}_{1i}^+)$ with*

$$A = v_{\nu,g}(x_1) \prod_{l \neq 1}^n \sigma_\mu(x_{1l}) \sigma_\mu(x_{1l}^+), \quad B = \alpha v_{\eta,g}(x_1) + \beta v_{\nu,g}(-x_1), \quad (6.59)$$

$$C_i = v_{\nu,g}(x_i) \sigma_{\xi_{12}}(x_{1i}) \sigma_\mu(x_{1i}^+) \prod_{l \neq 1, i} \sigma_\mu(x_{il}) \sigma_\mu(x_{il}^+), \quad (6.60)$$

$$D_i = (C_i)^{\mathfrak{s}_i} = v_{\nu,g}(-x_i) \sigma_{\xi_{12}}(x_{1i}^+) \sigma_\mu(x_{1i}) \prod_{l \neq 1, i} \sigma_\mu(-x_{li}^+) \sigma_\mu(x_{li}), \quad (6.61)$$

where

$$\alpha = \left\{ -1 + \sum_{i=1}^{n-1} \frac{\sigma_{\xi_{12}}(-i/n)\sigma_{\eta+\nu}(i/n)}{\sigma_{\mu}(i/n)^2} \right\} \prod_{l=1}^{n-1} \sigma_{\mu}(l/n)^2, \quad (6.62)$$

$$\beta = \sum_{i \neq 1} \sigma_{\xi_{12}}(x_{1i})\sigma_{\eta+\nu}(x_{1i}^+) \prod_{l \neq 1, i} \sigma_{\mu}(x_{il})\sigma_{\mu}(x_{l1}). \quad (6.63)$$

In these formulas $\xi_1 = \eta$ is a spectral parameter, $\xi_{12} = \eta + \nu + (n-2)\mu$.

Proof. Let us write $v_{\nu}(z)$ instead of $v_{\nu, g}(z)$ to simplify the notation. The coefficients A , C_i and D_i are calculated in the same way as in Proposition 4.2. For calculating B we expand the product (6.58) and collect the terms that reduce to \mathfrak{s}_1 when restricted on M' . It is easy to check that this happens only for the following choices of the terms in each of the factors: (1) U , $v_{\xi_1}(x_1)\mathfrak{s}_1$, U^+ ; (2) $V_i\mathfrak{s}_{1i}$, $v_{\nu}(x_1)$, $V_i^+\mathfrak{s}_{1i}^+$; (3) $V_i\mathfrak{s}_{1i}$, $v_{\xi_1}(x_1)\mathfrak{s}_1$, $V_i^+\mathfrak{s}_{1i}^+$. As a result, B is calculated from

$$B = -v_{\xi_1}(x_1)U(U^+)^{\mathfrak{s}_1} + \sum_{i \neq 1} v_{\nu}(x_i)V_i(V_i^+)^{\mathfrak{s}_{1i}} - \sum_{i \neq 1} v_{\xi_1}(x_i)V_i(V_i^+)^{\mathfrak{s}_{1i}\mathfrak{s}_1},$$

resulting in

$$B = -v_{\xi_1}(x_1) \prod_{l \neq 1} \sigma_{\mu}(x_{1l})\sigma_{\mu}(x_{l1}) \\ + \sum_{i \neq 1} \sigma_{\xi_{12}}(x_{1i}) \left\{ v_{\nu}(x_i)\sigma_{\xi_{1n}^+}(x_{1i}^+) - v_{\xi_1}(x_i)\sigma_{\xi_{1n}^+}(x_{1i}) \right\} \prod_{l \neq 1, i} \sigma_{\mu}(x_{il})\sigma_{\mu}(x_{l1}).$$

Here $\xi_{12} = \eta + \nu + (n-2)\mu$, $\xi_{1n}^+ = \eta - \nu$ by (6.56). We now use the following identity [KH1, (2.8a)] which follows from the addition formulas for $\sigma_{\mu}(z)$:

$$v_{\nu}(x_i)\sigma_{\eta-\nu}(x_{1i}^+) - v_{\eta}(x_i)\sigma_{\eta-\nu}(x_{1i}) = v_{\nu}(-x_1)\sigma_{\eta+\nu}(x_{1i}^+) + v_{\eta}(x_1)\sigma_{\eta+\nu}(x_{1i}).$$

Using this in the previous formula leads to $B = \alpha v_{\eta}(x_1) + \beta v_{\nu}(-x_1)$, with

$$\alpha = - \prod_{l \neq 1} \sigma_{\mu}(x_{1l})\sigma_{\mu}(x_{l1}) + \sum_{i \neq 1} \sigma_{\xi_{12}}(x_{1i})\sigma_{\eta+\nu}(x_{1i}) \prod_{l \neq 1, i} \sigma_{\mu}(x_{il})\sigma_{\mu}(x_{l1}), \quad (6.64)$$

and with β as in (6.63). To see why α is, in fact, constant in x , we use symmetry arguments. Indeed, the operator given in the lemma must commute with the action of any $w \in W'$. Therefore, the coefficient B must be invariant under signed permutations of x_2, \dots, x_n . Now view B as a function of the parameters η, ν, μ . Then the residue of B at $\nu = 0$ is given by

$$\beta|_{\nu=0} = \sum_{i \neq 1} \sigma_{\eta+(n-2)\mu}(x_{1i})\sigma_{\eta}(x_{1i}^+) \prod_{l \neq 1, i} \sigma_{\mu}(x_{il})\sigma_{\mu}(x_{l1}).$$

Thus, this also must be W' -invariant for all η, μ . Obviously, replacing in this expression η by $\eta + \nu$ gives us back β , hence β is W' -invariant, and so must be α .

Now, from the formula for α it is easy to see that it is elliptic function of x_i , regular at $x_1 + x_i = 0$ and $x_i + x_l = 0$ with $i, l > 1$. By W' -symmetry, it follows that α is also regular at hyperplanes $x_1 - x_i = 0$ and $x_i - x_l = 0$. As a result, α is globally regular, so is a constant (depending on η, μ, ν). It can now be evaluated by setting $x_l = l/n$, $1 \leq l \leq n$ in (6.64), which leads to the expression (6.62). \square

Remark 6.42. We see from the proof that β is W' -invariant. This is easy to confirm for $n = 1$ when $B = -v_{\xi_1}(x_1)$, and for $n = 2$, in which case $\xi_{12} = \eta + \nu$ so we have

$$\alpha = -\wp(\mu) + \wp(\eta + \nu), \quad \beta = \sigma_{\eta+\nu}(x_{12})\sigma_{\eta+\nu}(x_{12}^+).$$

However, for $n > 2$ such a symmetry is not obvious from the formulas. Note that by this symmetry β does not have poles at $x_i - x_l = 0$ for $i, l > 1$.

To write down the Lax matrix, we use the same notation as in Proposition 4.3, namely, extend the set of vectors ϵ_i and variables x_i to the range $1 \leq i \leq 2n$ by setting $\epsilon_{i+n} = -\epsilon_i$ and $x_{n+i} = -x_i$

for $1 \leq i \leq n$. Denote by \mathcal{P}, \mathcal{Q} the following $2n \times 2n$ matrices:

$$\begin{aligned} \mathcal{P}_{ij} &= -v_{\nu,g}(x_j) \sigma_{\xi_{12}}(x_{ij}) \sigma_{\mu}(x_{ij}^{\dagger}) \prod_{l=1}^{2n}{}' \sigma_{\mu}(x_{jl}) \quad (i-j \neq 0, \pm n), & \mathcal{P}_{ii} &= v_{\nu,g}(x_i) \prod_{l=1}^{2n}{}' \sigma_{\mu}(x_{il}), \\ \mathcal{P}_{i,n+i} &= \alpha v_{\eta,g}(x_i) + \beta^{s_{1i}} v_{\nu,g}(-x_i), & \mathcal{P}_{n+i,i} &= \alpha v_{\eta,g}(-x_i) + \beta^{s_{1i}} v_{\nu,g}(x_i) \quad (1 \leq i \leq n), \\ \mathcal{Q}_{ii} &= v_{\bar{\nu},\bar{g}}(x_i + c/2) t(\epsilon_i), & \mathcal{Q}_{ij} &= -v_{\xi_{1,\bar{g}}}(x_i + c/2) \quad (i-j = \pm n), & \mathcal{Q}_{ij} &= 0 \quad (i-j \neq 0, \pm n). \end{aligned}$$

Here α, β are given by (6.62), (6.63), and the symbol \prod' in the formula for \mathcal{P}_{ij} indicates that we exclude those values of l where either $l-i$ or $l-j$ equals $0, \pm n$ (e.g., two values are excluded if $i=j$). Explicitly, we have

$$\begin{aligned} \beta^{s_{1i}} &= \sum_{j:j \neq i}^n \sigma_{\xi_{12}}(x_{ij}) \sigma_{\eta+\nu}(x_{ij}^{\dagger}) \prod_{l \neq i,j} \sigma_{\mu}(x_{jl}) \sigma_{\mu}(x_{li}), \\ \beta^{s_{1i}^{\dagger}} &= \sum_{j:j \neq i}^n \sigma_{\xi_{12}}(-x_{ij}^{\dagger}) \sigma_{\eta+\nu}(x_{ji}) \prod_{l \neq i,j} \sigma_{\mu}(x_{jl}) \sigma_{\mu}(x_{li}^{\dagger}). \end{aligned}$$

Proposition 6.43. *The quantum Lax matrix \mathcal{L} for the elliptic van Diejen system is $\mathcal{L} = \mathcal{P}\mathcal{Q}$. It satisfies the quantum Lax equation (1.6) for every quantum Hamiltonian $\widehat{H} = L^b$ of the van Diejen system and suitable \mathcal{A} . The classical Lax matrix is $L = PQ$ where $P = \mathcal{P}$, while Q is obtained from \mathcal{Q} by setting $c = 0$ and replacing $t(\epsilon_i)$ with $e^{\beta p_i}$ (with $p_{n+i} = -p_i$ for $1 \leq i \leq n$). The matrix L deforms isospectrally under each of the Hamiltonian flows of the classical van Diejen system. As a result, the functions $\text{tr} L^k$ are in involution.*

Remark 6.44. Let us remark on how one can calculate a Lax partner for the above \mathcal{L} . Let us look back at the calculation for the GL_n -case in Subsection §6.3. According to Proposition 6.13, the Lax partner \mathcal{A} in that case can be found by considering

$$\sum_{i=1}^n \prod_{j \neq i}^n \sigma_{\mu}(\xi_i - \xi_j) \widehat{Y}^{\epsilon_i} = \sum_{i=1}^n Y^{\epsilon_i}.$$

The reason why in Subsection §6.3 we used $Y^{\epsilon_1} + Y^{\epsilon_2}$ is that all other terms Y^{ϵ_j} vanish after specialising ξ and restricting onto M' . By a similar reasoning, in constructing a Lax partner for the van Diejen system one can use $Y^{\epsilon_1} + Y^{\epsilon_2} + Y^{-\epsilon_1}$.

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